Electric Potential

Flashback to PHY113...

- Work done by a conservative force is independent of the path of the object.
- Gravity and elastic forces are examples.
- This leads to the concept of potential energy and helps us avoid tackling problems using only forces.

 $W = \int_{x_i}^{x_f} F_s dx = -\Delta U$

Electrostatic forces are also conservative.

Electric Potential Energy



$$\vec{\mathbf{F}} = q_0 \vec{\mathbf{E}}$$
$$\Delta U = -\int_{path} \vec{\mathbf{F}} \cdot d\vec{\mathbf{s}} = -q_0 \int_{path} \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}}$$

...but **F** is a conservative force

...so the path we take does not matter

$$\Delta U = -q_0 \int_A^B \vec{\mathbf{E}} d\vec{\mathbf{s}}$$

Electric Potential

- Work, ΔU , is dependent on the magnitude of the test charge, q_0 .
- We'd like to have a quantity independent of the test charge and only an attribute of the electric field.

$$V = \frac{U}{q_0} \qquad \Delta V = \frac{\Delta U}{q_0} = -\int_A^B \vec{\mathbf{E}} \cdot d\vec{\mathbf{s}} \qquad V_P = -\int_\infty^P \vec{\mathbf{E}} \cdot d\mathbf{s}$$

Electric Potential (J/C=V) Potential Difference between A and B

Electric Potential at P

Unit Pit Stop

Potential
$$V = J/C$$
Energy $J = N \cdot m$ Electric Field $N/C = \frac{N}{C} \left(\frac{V.C}{J}\right) \left(\frac{J}{N.m}\right) = V/m$

Electron-Volt
$$1eV = e(1V) = (1.6 \times 10^{-19} C)(1J/C) = 1.6 \times 10^{-19} J$$

Electric Potential in a Uniform Field



$$\Delta V = -\int_{A}^{B} \mathbf{E} d\mathbf{s} = -\int_{A}^{B} E \cos 0^{\circ} ds = -E \int_{A}^{B} ds$$
$$\Delta V = -E d \qquad \Delta U = -q_{0} E d$$

Electric field lines point to decreasing potential.

A positive charge will lose potential energy and gain kinetic energy when moving in the direction of the field.

Equipotential Surfaces



$$\Delta V = -\int_{A}^{B} \mathbf{E} \cdot d\mathbf{s} = -\mathbf{E} \cdot \int_{A}^{B} d\mathbf{s} = -\mathbf{E} \cdot \mathbf{s}$$

$$\Delta V = -(\mathbf{E} \cdot \mathbf{s})_{AC} - (\mathbf{E} \cdot \mathbf{s})_{CB}$$
$$\Delta V = -Es_{AC} \cos 0^{\circ} - Es_{CB} \cos 90^{\circ}$$
$$\Delta V = -Es \cos \theta$$

$$\Delta U = -q_0 (\mathbf{E} \cdot \mathbf{s})_{AC} - q_0 (\mathbf{E} \cdot \mathbf{s})_{CB}$$
$$\Delta U = -q_0 E s_{AC} \cos 0^\circ$$
$$\Delta U = -q_0 E s \cos \theta$$

$$\Delta V_{AD} = -\int_{A}^{D} \mathbf{E} \cdot d\mathbf{s} = -\mathbf{E} \cdot \int_{A}^{D} d\mathbf{s} = -\mathbf{E} \cdot \mathbf{s} = -(\mathbf{E} \cdot \mathbf{s})_{AC} - (\mathbf{E} \cdot \mathbf{s})_{CD} = -Es\cos\theta$$

No work is done moving a charged particle perpendicular to a field (along equipotential surfaces)

Equipotential Surfaces







Uniform Field

Point Charge

Electric Dipole

Electric Potential of a Point Charge



$$V_B - V_A = -\int_A^B \mathbf{E} \cdot d\mathbf{s}$$

$$\mathbf{E}.d\mathbf{s} = k_e \frac{q}{r^2} \hat{\mathbf{r}}.d\mathbf{s} = k_e \frac{q}{r^2} ds \cos\theta = k_e \frac{q}{r^2} dr$$

$$V_B - V_A = -\int E_r dr$$
$$V_B - V_A = -k_e q \int_{r_A}^{r_B} \frac{dr}{r^2} = k_e q \left[\frac{1}{r_B} - \frac{1}{r_A} \right]$$

$$V_{A=\infty}=0$$

$$V = k_e \frac{q}{r}$$

Two Point Charges



$$V_P = k_e \left(\frac{q_1}{r_1} + \frac{q_2}{r_2}\right)$$

A System of Point Charges



Concept Question

Two test charges are brought separately into the vicinity of a charge +Q. First, test charge +q is brought to a point a distance *r* from +Q. Then this charge is removed and test charge -q is brought to the same point. The electrostatic potential energy of which test charge is greater:



+*q* -*q* It is the same for both.

Getting From V to E

$$V = -\int_{A}^{B} \mathbf{E}.d\mathbf{s}$$
$$dV = -\mathbf{E}.d\mathbf{s}$$









n general:
$$E_x = -\frac{\partial V}{\partial x}$$

$$E_{y} = -\frac{\partial V}{\partial y}$$

$$E_z = -\frac{\partial V}{\partial z}$$

Electric Potential of a Dipole



$$V_{P} = \frac{2k_{e}qa}{x^{2} - a^{2}}$$
$$V \approx \frac{2k_{e}qa}{x^{2}} \quad (x \gg a)$$
$$E_{x} = -\frac{dV}{dx} = \frac{4k_{e}qa}{x^{3}}$$

Between the Charges

$$V_{P} = \frac{2k_{e}qx}{a^{2} - x^{2}} \quad E_{x} = -\frac{dV}{dx} = -2k_{e}q\left(\frac{a^{2} + x^{2}}{(a^{2} - x^{2})^{2}}\right)$$



Electric Potential Due to Continuous Charge Distributions



Start with an infinitesimal charge, *dq*.

$$dV = k_e \frac{dq}{r}$$

Then integrate over the whole distribution

$$V = k_e \int \frac{dq}{r}$$

Electric Potential Due to a Uniformly Charged Ring



$$V = \frac{k_e Q}{\sqrt{x^2 + a^2}}$$

$$E = \frac{k_e Q x}{\left(x^2 + a^2\right)^{3/2}}$$

Electric Potential Due to a Finite Line of Charge



 $k O \left(l \pm \sqrt{l^2 \pm a^2} \right)$

$$V = \frac{k_e Q}{l} \ln \left(\frac{l + \sqrt{l^2 + a^2}}{a} \right)$$

Electric Potential Due to a Uniformly Charged Sphere



$$E_r = \frac{k_e Q}{R^3} r$$

r < R

$$V_{D} - V_{C} = -\int_{R}^{r} E_{r} dr = \frac{k_{e}Q}{2R^{3}} \left(R^{2} - r^{2}\right)$$

$$V_D = \frac{k_e Q}{2R} \left(3 - \frac{r^2}{R^2} \right)$$

$$E_r = k_e \frac{Q}{r^2}$$

$$V_B = -\int_{\infty}^{r} E_r dr = -k_e Q \int_{\infty}^{r} \frac{dr}{r^2}$$

$$V_B = k_e \frac{Q}{r}$$

r > R

$$V_C = k_e \frac{Q}{R}$$

Potential Due to a Charged Conductor



- Charges always reside at the outer surface of the conductor.
- The field lines are always perpendicular to surface.
- Then **E**.*d***s**=0 on the surface at any point.
- Which means, $V_B V_A = 0$ along the surface.
- The surface is an equipotential surface.
- Finally, since E=0 inside the conductor, the potential V is constant and equal to the surface value.



Connected Charged Conducting Spheres



Cavity Within a Conductor



$$V_B - V_A = -\int_A^B \mathbf{E}.d\mathbf{s} = 0$$

We can always find a path where E.*d*s is non-zero.

But, since $\Delta V=0$ for all paths, E must be zero everywhere in the cavity.

A cavity without any charges enclosed by a conducting wall is field free.

Summary

- Charges can have different electric potential energy at different points in an electric field.
- Electric potential is the electric potential energy per unit charge.
- All points inside a conductor are at the same potential.