# CHAPTER-22
# ELECTROSTATICS

## 1. INTRODUCTION

The branch of physics which deals with electric effect of static charge is called electrostatics.

## 2. ELECTRIC CHARGE

Charge of a material body or particle is the property (acquired or natural) due to which it produces and experiences electrical and magnetic effects. Some of naturally occurring charged particles are electrons, protons, $\alpha$-particles etc.

Charge is a derived physical quantity & is measured in Coulomb in S.I. unit. In practice we use mC($10^{-3}$C), $\mu$C ($10^{-6}$C), nC($10^{-9}$C) etc.
C.G.S. unit of charge = electrostatic unit = esu.
1 coulomb = $3 \times 10^9$ esu of charge
Dimensional formula of charge = $[M^oL^oT^1I^1]$

### 2.1 Properties of Charge

**(i) Charge is a scalar quantity :** It adds algebraically and represents excess or deficiency of electrons.

**(ii) Charge is of two types : (i) Positive charge and (ii) Negative charge** Charging a body implies transfer of charge (electrons) from one body to another. Positively charged body means loss of electrons i.e. deficiency of electrons. Negatively charged body means excess of electrons. This also shows that **mass of a negatively charged body > mass of a positively charged identical body.**

**(iii) Charge is conserved :** In an isolated system, total charge (sum of positive and negative) remains constant whatever change takes place in that system.

**(iv) Charge is quantized :** Charge on any body always exists in integral multiples of a fundamental unit of electric charge. This unit is equal to the magnitude of charge on electron ($1e = 1.6 \times 10^{-19}$ coulomb). So charge on anybody is $Q = \pm ne$, where n is an integer and e is the charge of the electron. **Millikan's oil drop** experiment proved the quantization of charge or atomicity of charge

**Note :** Recently, the existence of particles of charge $\pm\frac{1}{3}e$ and $\pm\frac{2}{3}e$ has been postulated. These particles are called quarks but still this is not considered as the quantum of charge because these are unstable (They have very short span of life).

**(v)** Like point charges repel each other while unlike point charges attract each other.

**(vi)** Charge is always associated with mass, i.e., charge can not exist without mass though mass can exist without charge. The particle such as photon or neutrino which have no (rest) mass can never have a charge.

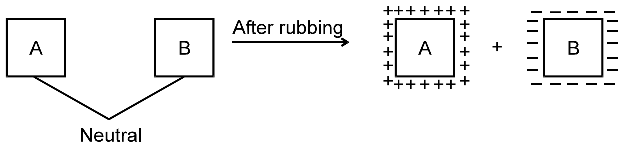
**(vii) Charge is relativistically invariant:** This means that charge is independent of frame of reference i.e. charge on a body does not change whatever be its speed. This property is worth mentioning as in contrast to charge, the mass of a body depends on its speed and increases with increase in speed.

**(viii)** A charge at rest produces only electric field around itself, a charge having uniform motion produces electric as well as magnetic field around itself while a charge having accelerated motion emits electromagnetic radiations.

## 2.2 Charging of a body

A body can be charged by means of (a) friction, (b) conduction, (c) induction, (d) thermionic ionization or thermionic emission (e) photoelectric effect and (f) field emission.

**(a) Charging by Friction :** When a neutral body is rubbed against other neutral body then some electrons are transferred from one body to other. The body which can hold electrons tightly, draws some electrons and the body which can not hold electrons tightly, loses some electrons. The body which draws electrons becomes negatively charged and the body which loses electrons becomes positively charged.




For example : Suppose a glass rod is rubbed with a silk cloth. As the silk can hold electrons more tightly and a glass rod can hold electrons less tightly (due to their chemical properties), some electrons will leave the glass rod and get transferred to the silk. So, in the glass rod their will be deficiency of electrons, therefore it will become positively charged. And in the silk, there will be some extra electrons, so it will become negatively charged

**(b) Charging by conduction (flow):** There are three types of materials in nature

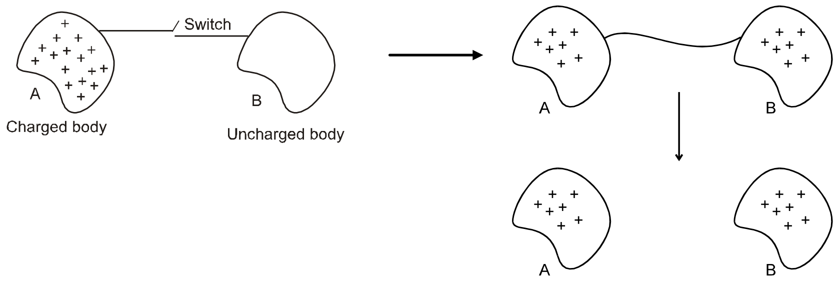
**(i) Conductor :** Conductors are the material in which the outer most electrons are very loosely bound, so they are free to move (flow). So in a conductor, there are large number of free electrons.

Ex. Metals like Cu, Ag, Fe, Al............

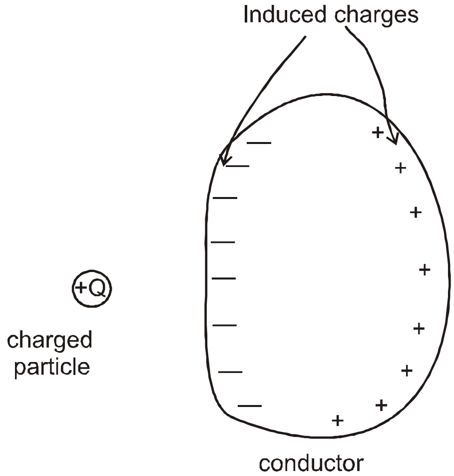
**(ii) Insulator or Dielectric or Nonconductor :** Non-conductors are the materials in which outer most electrons are very tightly bound, so that they cannot move (flow). Hence in a non-conductor there are no free electrons. Ex. plastic, rubber, wood etc.

**(iii) Semi conductor :** Semiconductors are the materials which have free electrons but very less in number.

Now lets see how the charging is done by conduction. In this method, we take a charged conductor 'A' and an uncharged conductor 'B'. When both are connected, some charge will flow from the charged body to the uncharged body. If both the conductors are identical & kept at large distance and connected to each other, then charge will be divided equally in both the conductors otherwise they will flow till their electric potential becomes same. Its detailed study will be done in last section of this chapter.



**(c) Charging by Induction :** To understand this, let's have introduction to induction.



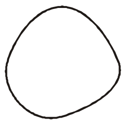

We have studied that there are lot of free electrons in the conductors. When a charged particle +Q is brought near a neutral conductor, due to attraction of +Q charge, many electrons (–ve charges) come closer and accumulate on the closer surface.

On the other hand, a positive charge (deficiency of electrons) appears on the other surface. The flow of charge continues till the resultant force on free electrons of the conductor becomes zero. This phenomena is called induction and charges produced are called induced charges.

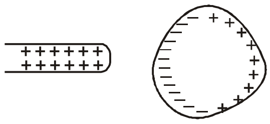
**A body can be charged by induction in the following two ways :**

**Method I :**

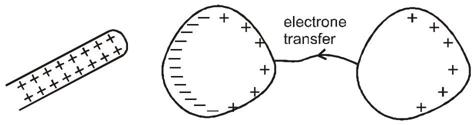
**Step 1 :** Take an isolated neutral conductor..



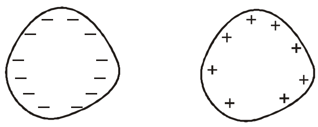
**Step 2:** Bring a charged rod near it. Due to the charged rod, charges will induce on the conductor.



**Step 3 :** Connect another neutral conductor with it. Due to attraction of the rod, some free electrons will move from the right conductor to the left conductor and due to deficiency of electrons positive charges will appear on right conductor and on the left conductor, there will be excess of electrons due to transfer from right conductor.



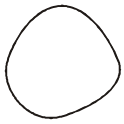

**Step 4 :** Now disconnect the connecting wire and remove the rod.



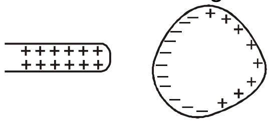
The first conductor will be negatively charged and the second conductor will be positively charged.

**Method II**

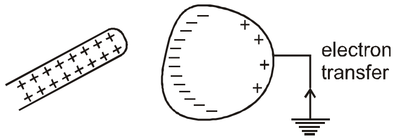
**Step 1:** Take an isolated neutral conductor.



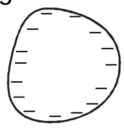
**Step 2 :** Bring a charged rod near it. Due to the charged rod, charges will induce on the conductor.



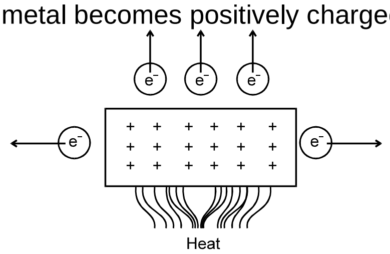
**Step 3 :** Connect the conductor to the earth (this process is called grounding or earthing). Due to attraction of the rod, some free electrons will move from earth to the conductor, so in the conductor there will be excess of electrons due to transfer from the earth, so net charge on conductor will be negative.



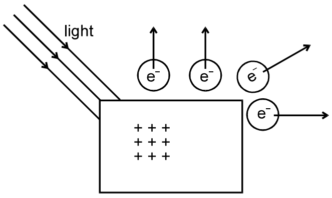

**Step 4.** Now disconnect the connecting wire. Conductor becomes negatively charged.



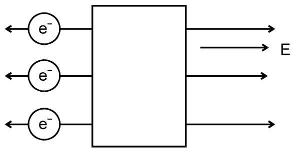
**(d)** **Thermionic emission :** When the metal is heated at a high temperature then some electrons of metals are ejected and the metal becomes positively charged.




**(e)** **Photoelectric effect :** When light of sufficiently high frequency is incident on metal surface then some electrons gain energy from light and come out of the metal surface and remaining metal becomes positively charged.

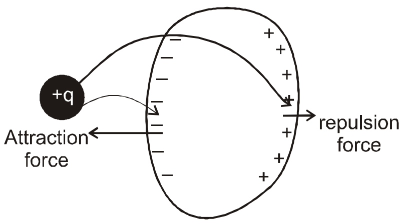



**(f)** **Field emission :** When electric field of large magnitude is applied near the metal surface then some electrons come out from the metal surface and hence the metal gets positively charged.




**Example 1.** If a charged body is placed near a neutral conductor, will it attract the conductor or repel it ?

**Solution :**



If a charged body (+ve) is placed left side near a neutral conductor, (–ve) charge will induce at left surface and (+ve) charge will induce at right surface. Due to positively charged body –ve induced charge will feel attraction and the +ve induced charge will feel repulsion. But as the –ve induced charge is nearer, so the attractive force will be greater than the repulsive force. So the net force on the conductor due to positively charged body will be attractive. Similarly, we can prove for negatively charged body also.

From the above example we can conclude that. "A charged body can attract a neutral body."

If there is attraction between two bodies then one of them may be neutral. But if there is repulsion between two bodies, both must be charged (similarly charged). So **"repulsion is the sure test of electrification".**

**Example 2.** A positively charged body 'A' attracts a body 'B' then charge on body 'B' may be:

(A) positive (B) negative (C) zero (D) can't say

**Answer :** B, C

**Example 3.** Five styrofoam balls A, B, C, D and E are used in an experiment. Several experiments are performed on the balls and the following observations are made :

(i) Ball A repels C and attracts B.

(ii) Ball D attracts B and has no effect on E.

(iii) A negatively charged rod attracts both A and E.

For your information, an electrically neutral styrofoam ball is very sensitive to charge induction and gets attracted considerably, if placed nearby a charged body. What are the charges, if any, on each ball ?

| | A | B | C | D | E |
|---|---|---|---|---|---|
| (A) | + | – | + | 0 | + |
| (B) | + | – | + | + | 0 |
| (C) | + | – | + | 0 | 0 |
| (D) | – | + | – | – | 0 |

**Answer :** C

**Solution :** From (i), as A repels C, so both A and C must be charged similarly. Either both are +ve or both are –ve. As A also attract B, so charge on B should be opposite of A or B may be uncharged conductor.

From (ii) as D has no effect on E, so both D and E should be uncharged and as B attracts uncharged D, so B must be charged and D must be an uncharged conductor.

From (iii), a –vely charged rod attracts the charged ball A, so A must be +ve and from exp. (i) C must also be +ve and B must be –ve.

**Example 4.** Charge conservation is always valid. Is it also true for mass?

**Solution :** No, mass conservation is not always. In some nuclear reactions, some mass is lost and it is converted into energy.

**Example 5.** What are the differences between charging by induction and charging by conduction ?

**Solution :** Major differences between two methods of charging are as follows :

(i) In induction, two bodies are close to each other but do not touch each other while in conduction they touch each other. (Or they are connected by a metallic wire)

(ii) In induction, total charge of a body remains unchanged while in conduction it changes.

(iii) In induction, induced charge is always opposite in nature to that of source charge while in conduction charge on two bodies finally is of same nature.

**Example 6.** If a glass rod is rubbed with silk, it acquires a positive charge because :

(A) protons are added to it (B) protons are removed from it

(C) electrons are added to it (D) electrons are removed from it.

**Answer :** D

## 3. COULOMB'S LAW (INVERSE SQUARE LAW)

On the basis of experiments Coulomb established the following law known as Coulomb's law :
The magnitude of electrostatic force between two point charges is directly proportional to the product of charges and inversely proportional to the square of the distance between them.

i.e. $F \propto q_1q_2$ and $F \propto \frac{1}{r^2}$ $\Rightarrow F \propto \frac{q_1q_2}{r^2}$ $\Rightarrow$ $F = \frac{Kq_1q_2}{r^2}$

Important points regarding Coulomb's law :

**(i)** It is applicable only for point charges.

**(ii)** The constant of proportionality K in SI units in vacuum is expressed as $\frac{1}{4\pi\varepsilon_0}$ and in any other medium expressed as $\frac{1}{4\pi\varepsilon}$. If charges are dipped in a medium then electrostatic force on one charge is $\frac{1}{4\pi\varepsilon_0\varepsilon_r}\frac{q_1q_2}{r^2}$ where $\varepsilon_0$ and $\varepsilon$ are called permittivity of vacuum and absolute permittivity of the medium respectively. The ratio $\varepsilon/\varepsilon_0 = \varepsilon_r$ is called relative permittivity of the medium, which is a dimensionless quantity.

**(iii)** The value of relative permittivity $\varepsilon_r$ is constant for a medium and can have values between 1 to $\infty$. For vacuum, by definition it is equal to 1. For air it is nearly equal to 1 and may be taken to be equal to 1 for calculations. For metals, the value of $\varepsilon_r$ is $\infty$ and for water is 81. The material in which more charge can induce $\varepsilon_r$ will be higher.

**(iv)** The value of $\frac{1}{4\pi\varepsilon_0} = 9 \times 10^9$ Nm$^2$ C$^{-2}$ & $\varepsilon_0 = 8.855 \times 10^{-12}$ C$^2$/Nm$^2$.

Dimensional formula of $\varepsilon$ is $[M^{-1} L^{-3} T^4 A^2]$

**(v)** The force acting on one point charge due to the other point charge is always along the line joining these two charges. It is equal in magnitude and opposite in direction on two charges, irrespective of the medium in which they lie.

**(vi)** The force is conservative in nature i.e., work done by electrostatic force in moving a point charge along a closed loop of any shape is zero.

**(vii)** Since the force is a central force, in the absence of any other external force, angular momentum of one particle w.r.t. the other particle (in two particle system) is conserved.

**(viii)** In vector form formula can be given as below.

$\vec{F} = \frac{1}{4\pi\varepsilon_0\varepsilon_r}\frac{q_1q_2}{|\vec{r}|^3}\vec{r} = \frac{1}{4\pi\varepsilon_0\varepsilon_r}\frac{q_1q_2}{|\vec{r}|^2}\hat{r}$ ; ($q_1$ & $q_2$ are to be substituted with sign.)

Here, $\vec{r}$ is position vector of the test charge (on which force is to be calculated) with respect to the source charge (due to which force is to be calculated).

**Example 7.** Find out the electrostatic force between two point charges placed in air (each of +1 C) if they are separated by 1m.

**Solution :** $F_e = \frac{kq_1q_2}{r^2} = \frac{9\times10^9\times1\times1}{1^2} = 9\times10^9$ N

From the above result, we can say that 1 C charge is too large to realize. In nature, charge is usually of the order of $\mu$C

**Example 8.** A particle of mass m carrying charge $q_1$ is revolving around a fixed charge $-q_2$ in a circular path of radius r. Calculate the period of revolution and its speed also.

**Solution :** $\frac{1}{4\pi\varepsilon_0}\frac{q_1q_2}{r^2} = mr\omega^2 = \frac{4\pi^2 mr}{T^2}$,

$T^2 = \frac{(4\pi\varepsilon_0)r^2(4\pi^2 mr)}{q_1q_2}$ or $T = 4\pi r\sqrt{\frac{\pi\varepsilon_0 mr}{q_1q_2}}$

and also we can say that

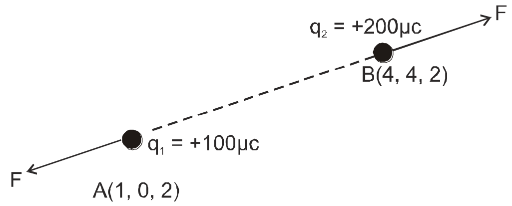
$\frac{q_1q_2}{4\pi\varepsilon_0 r^2} = \frac{mv^2}{r}$ $\Rightarrow$ $V = \sqrt{\frac{q_1q_2}{4\pi\varepsilon_0 mr}}$

**Example 9.** A point charge $q_A$ = + 100 μc is placed at point A (1, 0, 2) m and another point charge $q_B$= +200μc is placed at point B (4, 4, 2) m. Find :

(i) Magnitude of electrostatic interaction force acting between them

(ii) Find $\vec{F}_A$ (force on A due to B) and $\vec{F}_B$ (force on B due to A) in vector form

**Solution :** (i)




Value of F : $\left|\vec{F}\right| = \frac{kq_Aq_B}{r^2} = \frac{(9\times10^9)\ (100\times10^{-6})\ (200\ \times10^{-6})}{\left(\sqrt{(4-1)^2 + (4-0)^2 + (2-2)^2}\right)^2} = 7.2$ N

(ii) Force on B, $\vec{F}_B = \frac{kq_Aq_B}{|\vec{r}|^3}\vec{r} = \frac{(9\times10^9)(100\times10^{-6})(200\times10^{-6})}{\left(\sqrt{(4-1)^2+(4-0)^2+(2-2)^2}\right)^3}\left[(4-1)\hat{i}+(4-0)\hat{j}+(2-2)\hat{k}\right]$

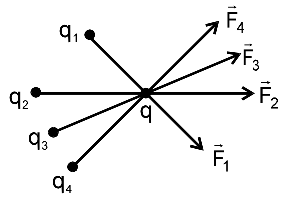
$= 7.2\left(\frac{3}{5}\hat{i}+\frac{4}{5}\hat{j}\right)$ N

Similarly $\vec{F}_A = 7.2\text{N}\left(-\frac{3}{5}\hat{i}-\frac{4}{5}\hat{j}\right)$ N

**Action($\vec{F}_A$)and Reaction ($\vec{F}_B$) are equal but in opposite direction.**

## 4. PRINCIPLE OF SUPERPOSITION

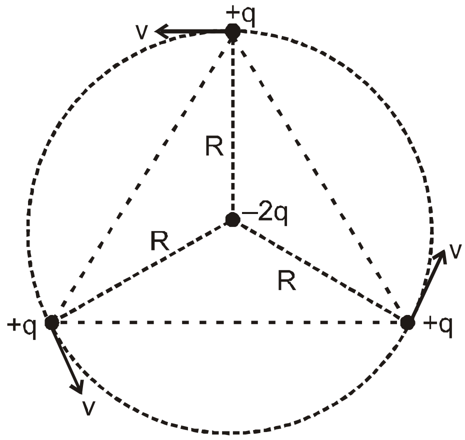
The electrostatic force is a two body interaction i.e. electrical force




The electrostatic force is a two body interaction i.e., electrical force between two point charges is independent of presence or absence of other charges and so the principle of superposition is valid i.e., force on charged particle due to number of point charges is the resultant of forces due to individual point charges. Therefore, force on a point test charge due to many charges is given by.

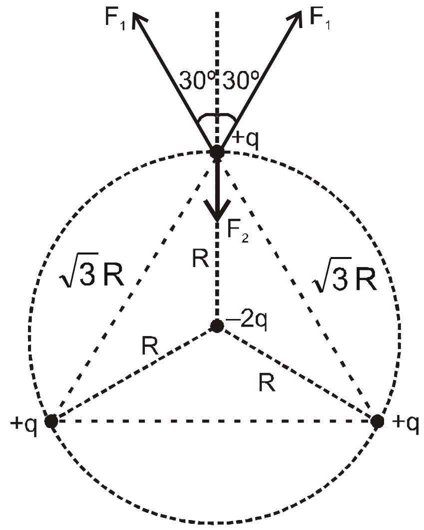
$\vec{F} = \vec{F}_1 + \vec{F}_2 + \vec{F}_3 + ...........$

**Example 10.** Three equal point charges of charge +q each are moving along a circle of radius R and a point charge –2q is also placed at the centre of circle (as shown in figure). If charges are revolving with constant and same speed in the circle then calculate speed of charges




**Solution :** $F_2 - 2F_1 \cos 30^o = \frac{mv^2}{R}$

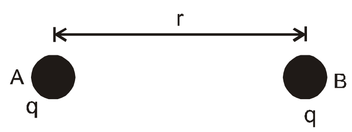
$\Rightarrow \frac{K(q)\ \ (2q)}{R^2} - \frac{2(Kq^2)}{(\sqrt{3}R)^2} \cos 30 = \frac{mv^2}{R}$




$\Rightarrow v = \sqrt{\frac{kq^2}{Rm}\left[2 - \frac{1}{\sqrt{3}}\right]}$

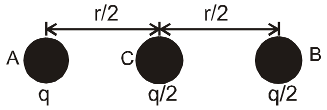
**Example 11.** Two equally charged identical small metallic spheres A and B repel each other with a force $2 \times 10^{-5}$ N when placed in air (neglect gravitational attraction). Another identical uncharged sphere C is touched to B and then placed at the mid point of line joining A and B. What is the net electrostatic force on C ?

**Solution :** Let, initially the charge on each sphere be q and separation between their centres be r. Then according to given problem : F




$= \frac{1}{4\pi\varepsilon_0} \frac{q \times q}{r^2} = 2 \times 10^{-5}$ N

When sphere C touches B, the charge of B i.e. q will distribute equally on B and C as sphere are identical how charges on spheres;




$q_B = q_C = (q/2)$

So sphere C will experience a force

$F_{CA} = \dfrac{1}{4\pi\varepsilon_0}\dfrac{q(q/2)}{(r/2)^2}$ = 2F along $\overrightarrow{AB}$ due to charge on A.

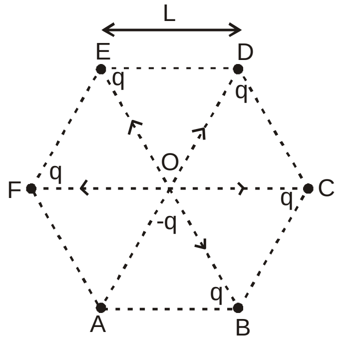
and, $F_{CB} = \dfrac{1}{4\pi\varepsilon_0}\dfrac{(q/2)(q/2)}{(r/2)^2}$ = F, along $\overrightarrow{BA}$ due to charge on B :

So the net force $F_C$ on C due to charges on A and B,

$F_C = F_{CA} - F_{CB} = 2F - F = 2 \times 10^{-5}$ N along $\overrightarrow{AB}$.

**Example 12.** Five point charges, each of value q are placed on five vertices of a regular hexagon of side L. What is the magnitude of the force on a point charge of value – q coulomb placed at the centre of the hexagon?

**Solution :** **Method-I :** If there had been a sixth charge +q at the remaining vertex of hexagon, force due to all the six charges on –q at O would have been zero (as the forces due to individual charges will balance each other), i.e., $\overrightarrow{F_R} = 0$



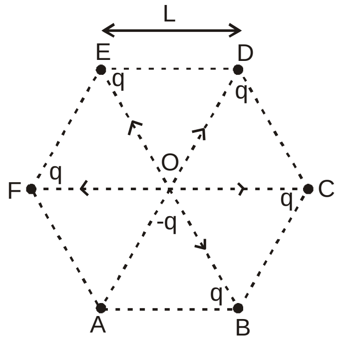

Now if $\vec{f}$ is the force due to sixth charge and $\vec{F}$ due to remaining five charges.

From $\vec{F} + \vec{f} = 0$ i.e. $\vec{F} = -\vec{f}$

or, $|F| = |f| = \dfrac{1}{4\pi\varepsilon_0}\dfrac{q \times q}{L^2} = \dfrac{1}{4\pi\varepsilon_0}\dfrac{q^2}{L^2}$

$\vec{F}_{Net} = \vec{F}_{OD} = \dfrac{1}{4\pi\varepsilon_0}\dfrac{q^2}{L^2}$ along OD

**Method-II :** In the diagram, we can see that force due to charge A and D are opposite to each other



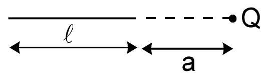

$\vec{F}_{oF} + \vec{F}_{OC} = \vec{0}$ ....(i)

Similarly $\vec{F}_{OB} + \vec{F}_{OE} = \vec{0}$ ....(ii)

So $\vec{F}_{oF} + \vec{F}_{OB} + \vec{F}_{OC} + \vec{F}_{OD} + \vec{F}_{OE} = \vec{F}_{Net}$

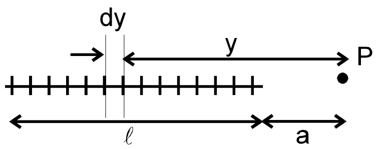
Using (i) and (ii) $\vec{F}_{Net} = \vec{F}_{OD} = \dfrac{1}{4\pi\varepsilon_0}\dfrac{q^2}{L^2}$ along OD.

**Example 13** A thin straight rod of length l carrying a uniformly distributed charge q is located in vacuum. Find the magnitude of the electric force on a point charge 'Q' kept as shown in the figure.




**Solution :** As the charge on the rod is not point charge, therefore, first we have to find force on charge Q due to charge over a very small part on the length of the rod. This part, called element of length dy can be considered as point charge.

Charge on element, $dq = \lambda dy = \dfrac{q}{\ell}dy$




Electric force on 'Q' due to element

$= \dfrac{K.dq.Q}{y^2} = \dfrac{K.Q.q.dy}{y^2.\ell}$

All forces are along the same direction,

$\therefore$ $F = \sum dF$. This sum can be calculated using integration,

therefore, $F = \int_{y=a}^{a+\ell} \frac{KQqdy}{y^2\ell} = \frac{KqQ}{\ell} = \left[-\frac{1}{y}\right]_a^{a+\ell} = \frac{KQ.q}{\ell}\left[\frac{1}{a} - \frac{1}{a+\ell}\right] = \frac{KQq}{a(a+\ell)}$

**Note : (1) The total charge of the rod cannot be considered to be placed at the centre of the rod as we do in mechanics for mass in many problems.**

**Note : (2)** If a >> l then, $F = \frac{KQq}{a^2}$

i.e. Behavior of the rod is just like a point charge.

# 5. ELECTROSTATIC EQUILIBRIUM

The point where the resultant force on a charged particle becomes zero is called equilibrium position.

**5.1 Stable Equilibrium:** A charge is initially in equilibrium position and is displaced by a small distance. If the charge tries to return back to the same equilibrium position then this equilibrium is called position of stable equilibrium.

**5.2 Unstable Equilibrium :** If charge is displaced by a small distance from its equilibrium position and the charge has no tendency to return to the same equilibrium position. Instead it goes away from the equilibrium position.

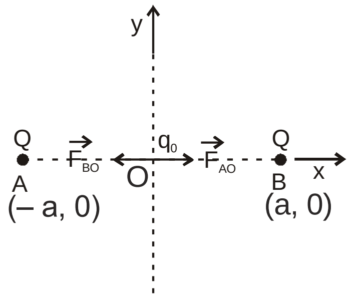
**5.3 Neutral Equilibrium :** If charge is displaced by a small distance and it is still in equilibrium condition then it is called neutral equilibrium.

**Example 14.** Two equal positive point charges 'Q' are fixed at points B(a, 0) and A(–a, 0). Another test charge $q_0$ is also placed at O(0, 0). Show that the equilibrium at 'O' is

(i) Stable for displacement along X-axis.

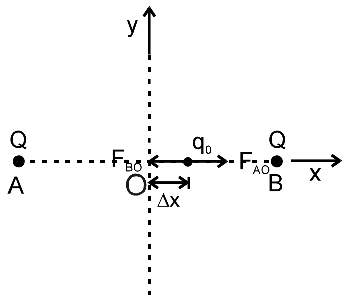
(ii) Unstable for displacement along Y-axis.

**Solution :** (i)




Initially $\vec{F}_{AO} + \vec{F}_{BO} = 0 \Rightarrow |\vec{F}_{AO}| = |\vec{F}_{BO}| = \frac{KQq_0}{a^2}$

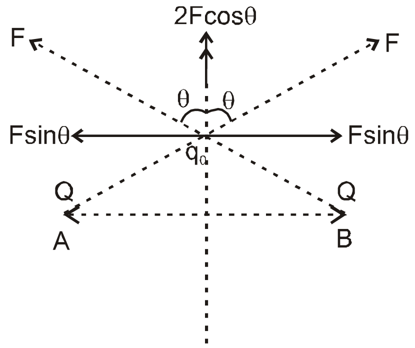
When charge is slightly shifted towards + x axis by a small distance $\Delta x$, then.




$|\vec{F}_{AO}| < |\vec{F}_{BO}|$

Therefore, the particle will move towards origin (its original position). Hence, the equilibrium is stable.

(ii) When charge is shifted along y axis:




After resolving components, net force will be along y axis So, the particle will not return to its original position & it is unstable equilibrium. Finally, the charge will move to infinity.

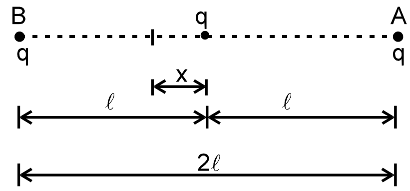
**Example 15.** Two point charges of charge $q_1$ and $q_2$ (both of same sign) and each of mass m are placed such that gravitational attraction between them balances the electrostatic repulsion. Are they in stable equilibrium? If not then what is the nature of equilibrium?

**Solution :** In given example : $\frac{Kq_1q_2}{r^2} = \frac{Gm^2}{r^2}$

We can see that irrespective of distance between them charges will main in equilibrium. If now distance is increased or decreased then there is no effect in their equilibrium. Therefore it is a neutral equilibrium.

**Example 16.** A particle of mass m and charge q is located midway between two fixed charged particles each having a charge q and a distance $2\ell$ apart. Prove that the motion of the particle will be SHM if it is displaced slightly along the line connecting them and released. Also find its time period.

**Solution :** Let the charge q at the mid-point is displaced slightly to the left. The force on the displaced charge q due to charge q at A,




$$F_1 = \frac{1}{4\pi\varepsilon_0}\frac{q^2}{(\ell+x)^2}$$

The force on the displaced charge q due to charge at B,

$$F_2 = \frac{1}{4\pi\varepsilon_0}\frac{q^2}{(\ell-x)^2}$$

Net restoring force on the displaced charge q.

$$F = F_2 - F_1 \text{ or } F = \frac{1}{4\pi\varepsilon_0}\frac{q^2}{(\ell-x)^2} - \frac{1}{4\pi\varepsilon_0}\frac{q^2}{(\ell+x)^2}$$

$$\text{or} \quad F = \frac{q^2}{4\pi\varepsilon_0}\left[\frac{1}{(\ell-x)^2} - \frac{1}{(\ell+x)^2}\right] = \frac{q^2}{4\pi\varepsilon_0}\frac{4\ell x}{(\ell^2-x^2)^2}$$

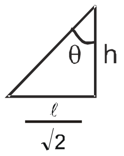
Since $\ell >> x$, $\therefore F = \frac{q^2\ell x}{\pi\varepsilon_0\ell^4}$ or $F = \frac{q^2x}{\pi\varepsilon_0\ell^3}$

Hence we see that $F \propto x$ and it is opposite to the direction of displacement. Therefore, the motion is SHM.

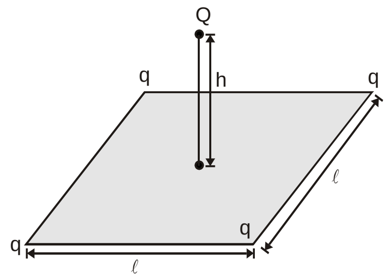
$$T = 2\pi\sqrt{\frac{m}{k}}, \text{ (here } k = \frac{q^2}{\pi \in_0 \ell^3}) \; T = 2\pi\sqrt{\frac{m\pi \in_0 \ell^3}{q^2}}$$

**Example 17.** Find out mass of the charge Q, so that it remains in equilibrium for the given configuration.

**Solution :**



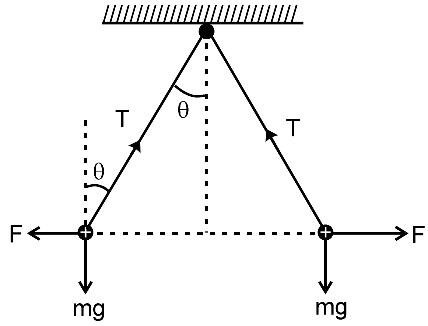

$\Rightarrow$ 4 F$\cos\theta$ = mg




$$\Rightarrow \ 4\times\frac{KQq}{\left(\frac{\ell^2}{2}+h^2\right)^{3/2}}\ h = mg \qquad \therefore\ m = \frac{4KQqh}{g\left(\frac{\ell^2}{2}+h^2\right)^{3/2}}$$

**Example 18.** Two identical charged spheres are suspended by strings of equal length. Each string makes an angle $\theta$ with the vertical. When suspended in a liquid of density $\sigma$ = 0.8 gm/cc, the angle remains the same. What is the dielectric constant of the liquid? (Density of the material of sphere is $\rho$ = 1.6 gm/cc.)

**Solution :** Initially as the forces acting on each ball are tension T, weight mg and electric force F, for its equilibrium along vertical




$T\cos\theta = mg$ .....(1)

and along horizontal

$T\sin\theta = F$ .....(2)

Dividing Eqn. (2) by (1), we have

$\tan\theta = \dfrac{F}{mg}$ .....(3)

When the balls are suspended in a liquid of density $\sigma$ and dielectric constant K, the electric force will become (1/K) times, i.e., F' = (F/K) while weight

$mg' = mg - F_B = mg - V\sigma g$ [as $F_B = V\sigma g$, where $\sigma$ is density of material of sphere]

i.e., $mg' = mg\left[1-\dfrac{\sigma}{\rho}\right]$ $\left[\text{as}\quad V=\dfrac{m}{\rho}\right]$

So, for equilibrium of ball,

$\tan\theta' = \dfrac{F'}{mg'} = \dfrac{F}{Kmg[1-(\sigma/\rho)]}$ ...(4)

According to given information $\theta' = \theta$; so from equations (4) and (3), we have :

$K = \dfrac{\rho}{(\rho-\sigma)} = \dfrac{(1.6)}{(1.6-0.8)} = 2$ **Ans.**

## 6. ELECTRIC FIELD

Electric field is the region around charged particle or charged body in which if another charge is placed, it experiences electrostatic force.

### 6.1 Electric field intensity $\vec{E}$ :

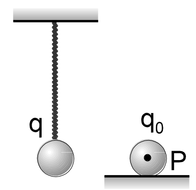
Electric field intensity at a point is equal to the electrostatic force experienced by a unit positive point charge both in magnitude and direction.

If a test charge $q_0$ is placed at a point in an electric field and experiences a force $\vec{F}$ due to some charges (called source charges), the electric field intensity at that point due to source charges is given by $\vec{E} = \frac{\vec{F}}{q_0}$

If the $\vec{E}$ is to be determined practically then the test charge $q_0$ should be small otherwise it will affect the charge distribution on the source which is producing the electric field and hence modify the quantity which is measured.

**Example 19.** A positively charged ball hangs from a long silk thread. We wish to measure E at a point P in the same horizontal plane as that of the hanging charge. To do so, we put a positive test charge $q_0$ at the point and measure $F/q_0$. Will $F/q_0$ be less than, equal to, or greater than E at the point in question?

**Solution :** When we try to measure the electric field at point P then after placing the test charge at P, it repels the source charge (suspended charge) and the measured value of electric field $E_{measured} = \frac{F}{q_0}$ will be less than the actual value $E_{act}$ ,that we wanted to measure.




## 6.2 Properties of electric field intensity $\vec{E}$ :

**(i)** It is a vector quantity. Its direction is the same as the force experienced by positive charge.

**(ii)** Direction of electric field due to positive charge is always away from it while due to negative charge, always towards it.

**(iii)** Its S.I. unit is Newton/Coulomb.

**(iv)** Its dimensional formula is $[MLT^{-3}A^{-1}]$

**(v)** Electric force on a charge q placed in a region of electric field at a point where the electric field intensity is $\vec{E}$ is given by $\vec{F} = q\vec{E}$.

Electric force on point charge is in the same direction of electric field on positive charge and in opposite direction on a negative charge.

**(vi)** It obeys the superposition principle, that is, the field intensity at a point due to a system of charges is vector sum of the field intensities due to individual point charges.

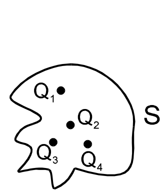
i.e. $\vec{E} = \vec{E}_1 + \vec{E}_2 + \vec{E}_3$ + .....

**(vii)** It is produced by source charges. The electric field will be a fixed value at a point unless we change the distribution of source charges.

**Example 20.** Electrostatic force experienced by –3μC charge placed at point 'P' due to a system 'S' of fixed point charges as shown in figure is $\vec{F} = (21\hat{i} + 9\hat{j})$ μN.

**(i)** Find out electric field intensity at point P due to S.

**(ii)** If now, 2μC charge is placed and –3 μC is removed at point P then force experienced by it will be.




**Solution :** **(i)** $\vec{F} = q\vec{E} \Rightarrow (21\hat{i} + 9\hat{j})\mu N = -3\mu C\,(\vec{E})$

$\Rightarrow \vec{E} = -7\hat{i} - 3\hat{j}\ \frac{N}{C}$

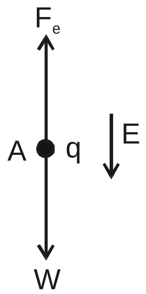
**(ii)** Since the source charges are not disturbed the electric field intensity at 'P' will remain same.

$\vec{F}_{2\mu C} = +2(\vec{E}) = 2(-7\hat{i} - 3\hat{j}) = (-14\hat{i} - 6\hat{j})\ \mu N$

**Example 21.** Calculate the electric field intensity which would be just sufficient to balance the weight of a particle of charge –10 $\mu$c and mass 10 mg. (Take g = 10 $ms^2$)

**Solution :** As force on a charge q in an electric field $\vec{E}$ is $\vec{F}_q = q\vec{E}$

So, according to given problem:

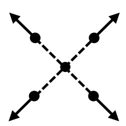
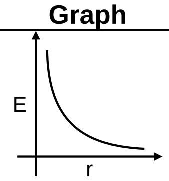
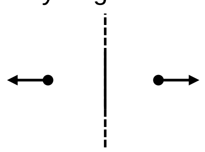
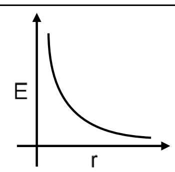
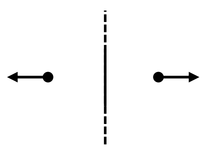
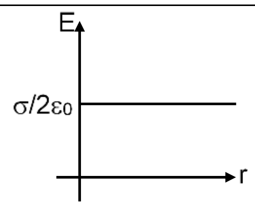


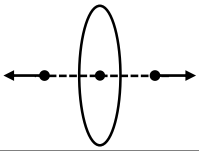
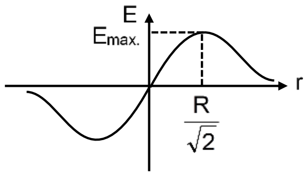
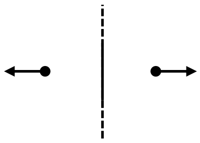
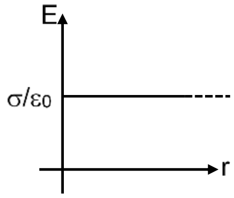
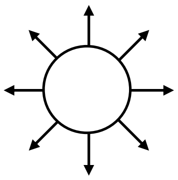
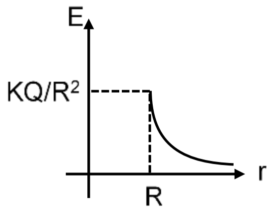
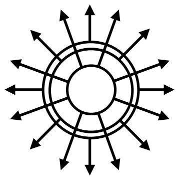
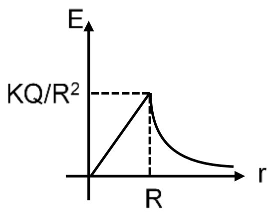

[W = weight of particle]

$|\vec{F}_q| = |\vec{W}|$ i.e., $|q|E = mg$

i.e., $E = \frac{mg}{|q|}$ = 10 N/C., in downward direction.

**List of formula for Electric Field Intensity due to various types of charge istribution :**

| Name/Type | Formula | Note | Graph |
|---|---|---|---|
| Point charge  | $\vec{E} = \frac{Kq}{\|\vec{r}\|^2} \cdot \hat{r}$ | * q is source charge. * $\vec{r}$ is vector drawn from source charge to the test point. Outwards due to +charges & inwards due to –charges. |   |
| Infinitely long line charge  | $\frac{\lambda}{2\pi\varepsilon_0 r}\hat{r} = \frac{2K\lambda\hat{r}}{r}$ | * q is linear charge density (assumed uniform) * r is perpendicular distance of point from line charge. * r is radial unit vector drawn from the charge to test point. |   |
| Infinite non-conducting thin sheet  | $\frac{\sigma}{2\varepsilon_0}\hat{n}$ | * $\sigma$ is surface charge density. (assumed uniform) * n is unit normal vector x = distance of point on the axis from centre of the ring. * electric field is always along the axis. |   |

| | | | |
|---|---|---|---|
| Uniformly charged ring  | $E = \dfrac{KQx}{(R^2+x^2)^{3/2}}$ <br> $E_{centre} = 0$ | * Q is total charge of the ring <br> * x = distance of point on the axis from centre of the ring. <br> * electric field is always along the axis. |   |
| Infinitely large charged conducting sheet  | $\dfrac{\sigma}{\varepsilon_0}\hat{n}$ | *$\sigma$ is the surface charge. density (assumed uniform) <br> * n is the unit vector perpendicular to the surface. |   |
| Uniformly charged hollow conducting / non-conducting / solid conducting sphere  | (i) for $r \geq R$ <br> $\vec{E} = \dfrac{kQ}{\lvert\vec{r}\rvert^2}\hat{r}$ <br> (ii) for $r < R$ <br> $E = 0$ | *R is radius of the sphere. <br> * $\vec{r}$ is vector drawn from centre of sphere to the point. <br> *Sphere acts like a point charge placed at centre for points outside the sphere. <br> * $\vec{E}$ is always along radial direction. <br> *Q is total charge (= $\sigma 4\pi R^2$.) <br> ($\sigma$ = surface charge density) |   |
| Uniformly charged solid non-conducting sphere (insulating material)  | (i) for $r \geq R$ <br> $\vec{E} = \dfrac{kQ}{\lvert\vec{r}\rvert^2}\hat{r}$ <br> (ii) for $r \leq R$ <br> $\vec{E} = \dfrac{kQ}{R^3}\vec{r}$ | * $\vec{r}$ is vector drawn from centre of sphere to the point <br> * Sphere acts like a point charge placed at the centre of points outside the sphere <br> * $\vec{E}$ is always along radial $dir^n$ <br> *Q is total charge $\left(\rho.\dfrac{4}{3}\pi R^3\right)$. <br> ($\rho$ = volume charge density) <br> * Inside the sphere $E \propto r$. <br> * Outside the sphere $E \propto 1/r^2$. |   |

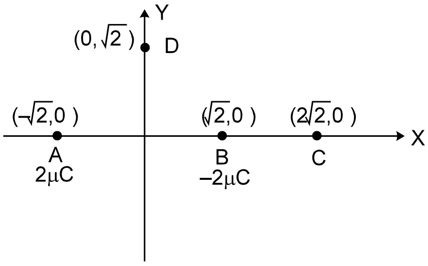
**Example 22.** Find out electric field intensity at point A (0, 1m, 2m) due to a point charge –20μC situated at point B( $\sqrt{2}$ m, 0, 1m).

**Solution :** $E = \frac{KQ}{|\vec{r}|^3}\vec{r} = \frac{KQ}{|\vec{r}|^2}\hat{r} \Rightarrow \vec{r}$ = P.V. of A – P.V. of B (P.V. = Position vector)

$= \left(-\sqrt{2}\hat{i} + \hat{j} + \hat{k}\right)$ $|\vec{r}| = \sqrt{(\sqrt{2})^2 + (1)^2 + (1)^2} = 2$

$E = \frac{9\times10^9 \times(-20\times10^{-6})}{8}\left(-\sqrt{2}\hat{i}+\hat{j}+\hat{k}\right) = -22.5\times10^3\left(-\sqrt{2}\hat{i}+\hat{j}+\hat{k}\right)$ N/C.

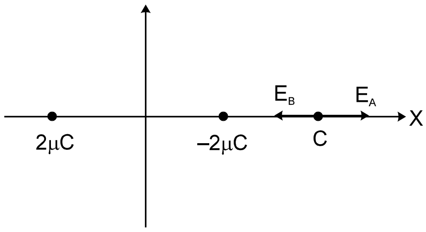
**Example 23.** Two point charges 2μc and – 2μc are placed at points A and B as shown in figure. Find out electric field intensity at points C and D. [All the distances are measured in meter].




**Solution :** Electric field at point C ($E_A$, $E_B$ are magnitudes only and arrows represent directions).

Electric field due to positive charge is away from it while due to negative charge, it is towards the charge.

It is clear that $E_B > E_A$.




$\therefore$ $E_{Net} = (E_B - E_A)$ towards negative X-axis

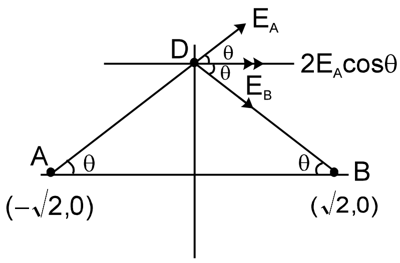
$= \frac{K(2\mu c)}{(\sqrt{2})^2} - \frac{K(2\mu c)}{(3\sqrt{2})^2}$ towards negative X-axis = 8000 $(-\hat{i})$ N/C

Electric field at point D :

Since magnitude of charges are same and also AD = BD

So, $E_A = E_B$

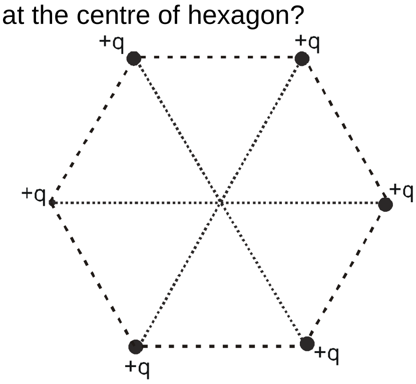
Vertical components of $\vec{E}_A$ and $\vec{E}_B$ cancel each other while horizontal components are in the same direction.




So, $E_{net} = 2E_A\cos\theta = \frac{2.K(2\mu c)}{2^2}\cos45°$

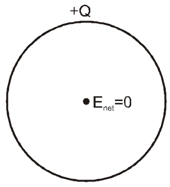
$= \frac{K\times10^{-6}}{\sqrt{2}} = \frac{9000}{\sqrt{2}}\hat{i}$ N/C.

**Example 24.** Six equal point charges are placed at the corners of a regular hexagon of side 'a'. Calculate electric field intensity at the centre of hexagon?




**Answer :** Zero (By symmetry)

**Similarly electric field due to a uniformly charged ring at the centre of ring :**




**Note : (i) Net charge on a conductor remains only on the outer surface of a conductor. This property will be discussed in the article of the conductor. (Article no.17)**

**(ii) On the surface of isolated spherical conductor charge is uniformly distributed.**

## 6.3 Electric field due to a uniformly charged ring and arc.

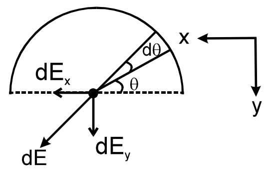
**Example 25.** Find out electric field intensity at the centre of a uniformly charged semicircular ring of radius R and linear charge density $\lambda$.

**Solution :** $\lambda$ = linear charge density.

The arc is the collection of large no. of point charges. Consider a part of ring as an element of length $Rd\theta$ which subtends an angle $d\theta$ at centre of ring and it lies between $\theta$ and $\theta + d\theta$

$\vec{dE} = dE_x\,\hat{i} + dE_y\,\hat{j}$ ; $E_x = \int dE_x = 0$ (due to symmetry)

& $E_y = \int dE_y = \int_0^{\pi} dE\sin\theta = \frac{K\lambda}{R}\int_0^{\pi}\sin\theta.d\theta = \frac{2K\lambda}{R}$

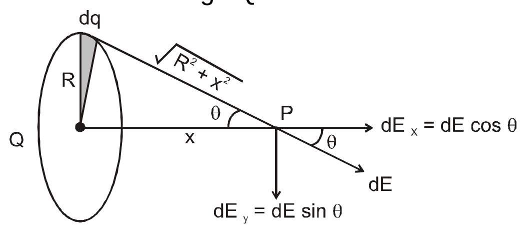



**Example 26.** Find out electric field intensity at the centre of uniformly charged quarter ring of radius R and linear charge density $\lambda$.

**Solution :** Refer to the previous question $\vec{dE} = dE_x\,\hat{i} + dE_y\,\hat{j}$ $\quad\therefore\quad$ on solving $E_{net} = \frac{K\lambda}{R}(\hat{i}+\hat{j})$

**By use of symmetry and from the formula of electric field due to half ring.**
**Above answer can be justified.**

**(ii)** Derivation of electric field intensity at a point on the axis at a distance x from centre of uniformly charged ring of radius R and total charge Q.



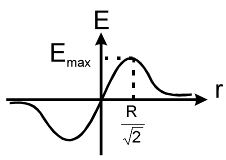

Consider an element of charge dq. Due to this element, the electric field at the point on axis, which is at a distance x from the centre of the ring is dE.

There are two components of this electric field

dE → $dE_X$, $dE_Y$

The y-component of electric field due to all the elements will be cancelled out to each other. So net electric field intensity at the point will be only due to X-component of each element.

$$E_{net} = \int dE_x = \int dE\cos\theta = \int_O^Q \frac{K(dq)}{R^2+x^2}\times\frac{x}{\sqrt{R^2+X^2}} = \frac{k\;\;x}{(R^2+x^2)^{3/2}}\int_O^Q dq$$




$$E_{net} = \frac{KQx}{[R^2+x^2]^{3/2}}$$

Graph for variation of E with r.

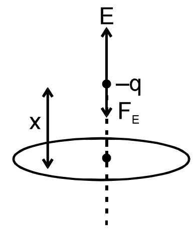
E will be max when $\frac{dE}{dx} = 0$, that is at $x = \frac{R}{\sqrt{2}}$ and $E_{max} = \frac{2KQ}{3\sqrt{3}\,R^2}$

**Case (i) :** if x>>R, $E = \frac{KQ}{x^2}$ Hence the ring will act like a point charge

**Case (ii) :** if x<<R, $E = \frac{KQ\,x}{R^3}$

**Example 27.** Positive charge Q is distributed uniformly over a circular ring of radius a. A point particle having a mass m and a negative charge –q, is placed on its axis at a distance y from the centre. Find the force on the particle. Assuming y << a, find the time period of oscillation of the particle if it is released from there. (Neglect gravity)

**Solution :** When the negative charge is shifted at a distance x from the centre of the ring along its axis then force acting on the point charge due to the ring :




$$F_E = qE \text{ (towards centre)} = q\left[\frac{KQy}{(a^2+y^2)^{3/2}}\right]$$

If a >> y then $a^2 + y^2 \simeq a^2$

$$\therefore \quad F_E = \frac{1}{4\pi\varepsilon_0}\frac{Qqy}{a^3} \text{ (Towards centre)}$$

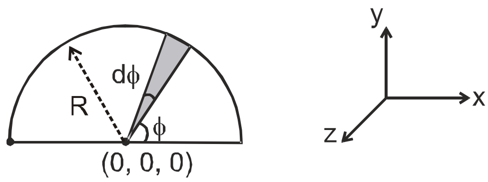
Since, restoring force $F_E \propto y$, therefore motion of charge the particle will be S.H.M.

Time period of SHM

$$T = 2\pi\sqrt{\frac{m}{k}} = 2\pi\sqrt{\frac{m}{\left(\frac{Qq}{4\pi\varepsilon_0 a^3}\right)}} = \left[\frac{16\pi^3\varepsilon_0 ma^3}{Qq}\right]^{1/2}$$

**Example 28.** Calculate electric field intensity at a point on the axis which is at distance x from the centre of half ring, having total charge Q distributed uniformly on it. The radius of half ring is R.

**Solution :** Consider an element of small angle $d\phi$ at an angle $\phi$ as shown.




Coordinates of element : $(R\cos\phi, R\sin\phi, 0)$

Coordinates of point : (0, 0, x)

Now electric field due to element :

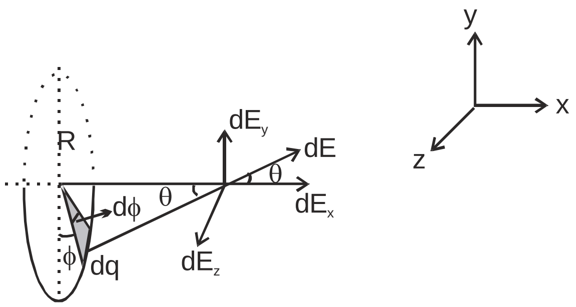
$$\overrightarrow{dE} = \frac{K(\lambda Rd\phi)\cdot[-R\cos\phi\,\hat{i} - R\sin\phi\,\hat{j} + x\hat{k}]}{(R^2\cos^2\phi + R^2\sin^2\phi + x^2)^{3/2}} \Rightarrow E_x = \Sigma dE_x = -\int_0^{\pi}\frac{K\lambda R^2\cos\phi\, d\phi}{(R^2+x^2)^{3/2}} = 0$$

$$E_y = \Sigma dE_y = -\int_0^{\pi}\frac{K\lambda R^2\sin\phi\, d\phi}{(R^2+x^2)^{3/2}} = \frac{2K\lambda R^2}{(R^2+x^2)^{3/2}} = \frac{2KQR}{\pi(R^2+x^2)^{3/2}}$$

$$E_z = \Sigma dE_z = \int_0^{\pi}\frac{K\lambda Rx\, d\phi}{(R^2+x^2)^{3/2}} = \frac{KQx}{(R^2+x^2)^{3/2}}$$

$$E_{net} = \sqrt{E_x^2 + E_y^2 + E_z^2} = \frac{KQ}{(R^2+x^2)^{3/2}}\sqrt{\frac{4R^2}{\pi^2} + x^2}$$

**Alternate solution :**




Consider an element of charge dq at an angle $\phi$ on circumference of half ring. Due to this element electric field at the point on axis, which is at a distance x from the centre of half ring is dE. This electric field can be resolved into three component.

dE

$dE_X$ $dE_Y$ $dE_Z$

$$E_Z = \int_{-\pi/2}^{\pi/2} dE\sin\theta\sin\phi = 0$$

$$E_X = \int_{-\pi/2}^{\pi/2} dE\cos\theta = \frac{KQx}{[R^2+x^2]^{3/2}} \qquad .....(1)$$

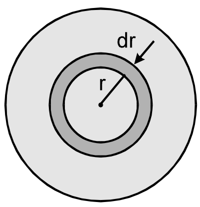
$$E_Y = \int_{-\pi/2}^{\pi/2} dE\sin\theta\cos\phi = \int \frac{Kdq}{R^2+x^2}\ \sin\theta.\cos\phi = \frac{2K\lambda R\sin\theta}{R^2+x^2} \qquad .....(2)$$

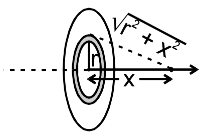
$\therefore$ $dq = \lambda R d\phi$, $\sin\theta = \dfrac{R}{\sqrt{R^2+x^2}}$ $\qquad \Rightarrow E_{net} = \sqrt{E_X^2+E_Y^2}$

**Example 29.** Derive the expression of electric field intensity at a point 'P' which is situated at a distance x on the axis of uniformly charged disc of radius R and surface charge density $\sigma$. Also, derive results for

(i) $x >> R$ (ii) $x << R$

**Solution :** The disc can be considered to be a collection of large number of concentric rings. Consider an element of the shape of rings of radius r and of width dr. Electric field due to this ring at P is




$$dE = \frac{K.\sigma 2\pi r.dr.x}{(r^2+x^2)^{3/2}}$$

Put, $r^2 + x^2 = y^2$

$2rdr = 2ydy$

$$\therefore\ dE = \frac{K.\sigma 2\pi y.dy.x}{y^3} = 2K\sigma\pi.x\ \frac{ydy}{y^3}$$

Electric field at P due to all rings is along the axis :

$$\therefore\ E = \int dE \Rightarrow E = 2K\sigma\pi x \int_x^{\sqrt{R^2+x^2}} \frac{1}{y^2}dy = 2K\rho\pi x.\left[-\frac{1}{y}\right]_x^{\sqrt{R^2+x^2}}$$

$$= 2K\sigma\pi x\left[+\frac{1}{x} - \frac{1}{\sqrt{R^2+x^2}}\right] = 2K\sigma\pi\left[1 - \frac{x}{\sqrt{R^2+x^2}}\right]$$

$$= \frac{\sigma}{2\varepsilon_0}\left[1 - \frac{x}{\sqrt{R^2+x^2}}\right] \text{ along the axis}$$

**Cases :(i)** If $x >> R$

$$E = \frac{\sigma}{2\varepsilon_0}\left[1 - \frac{x}{x\sqrt{\frac{R^2}{x^2}+1}}\right] = \frac{\sigma}{2\varepsilon_0}\left[1 - \left(1+\frac{R^2}{x^2}\right)^{-1/2}\right]$$

$$\frac{\sigma}{2\varepsilon_0} = \left[1 - 1 + \frac{1}{2}\frac{R^2}{x^2} + \text{higher order terms}\right] = \frac{\sigma}{4\varepsilon_0}\frac{R^2}{x^2} = \frac{\sigma\pi R^2}{4\pi\varepsilon_0 x^2} = \frac{Q}{4\pi\varepsilon_0 x^2}$$

i.e., behaviour of the disc is like a point charge.

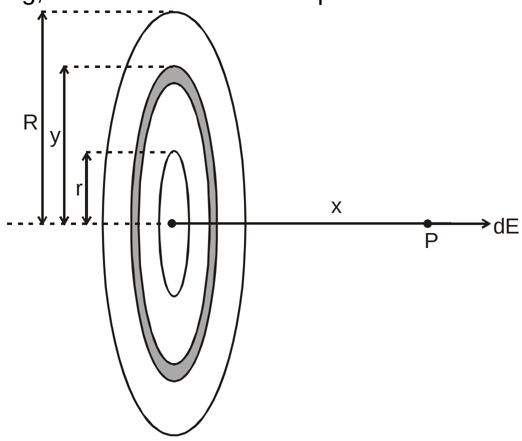
**(ii)** If $x << R$

$$E = \frac{\sigma}{2\varepsilon_0}[1-0] = \frac{\sigma}{2\varepsilon_0}$$

i.e., behaviour of the disc is like infinite sheet.

**Example 30.** Calculate electric field at a point on axis, which at a distance x from centre of uniformly charged disc having surface charge density $\sigma$ and R which also contains a concentric hole of radius r.

**Solution :** Consider a ring of radius y (r < y < R) and width dy concentric with disc and in the plane of the disc. Due to this ring, the electric field at the point P :




$$dE = \frac{K(dq)x}{[X^2 + Y^2]^{3/2}}$$

$$E_{net} = \int_r^R \frac{Kx.\sigma(2\pi y)dy}{[x^2+y^2]^{3/2}} \qquad [\because dq = \sigma 2\pi y dy]$$

$$E_{net} = \frac{2\pi\sigma kx}{2}\int_{x^2+r^2}^{x^2+R^2}\frac{dt}{t^{3/2}}, \text{put } x^2 + y^2 = t,\ 2y.\,dy = dt$$

$$= \frac{\sigma x}{2\varepsilon_0}\left[\frac{1}{\sqrt{x^2+r^2}} - \frac{1}{\sqrt{x^2+R^2}}\right] \quad \text{away from centre}$$

**Alternate method**

We can also use superposition principle to solve this problem.

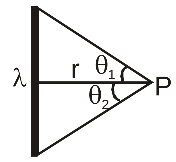
(i) Assume a disc without hole of radius R having surface charge density + $\sigma$.

(ii) Also assume a concentric disc of radius r in the same plane of first disc having charge density – $\sigma$. Now using derived formula in last example the net electric field at the centre is :

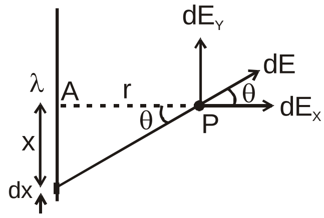
$$\vec{E}_{net} = \vec{E}_R + \vec{E}_r = \frac{\sigma x}{2\varepsilon_0}\left[\frac{1}{\sqrt{r^2+x^2}} - \frac{1}{\sqrt{R^2+x^2}}\right] \quad \text{away from centre.}$$

## 6.4 Electric field due to uniformly charged wire:

**(i) Line charge of finite length :** Derivation of expression for intensity of electric field at a point due to line charge of finite size of uniform linear charge density $\lambda$. The perpendicular distance of the point from the line charge is r and lines joining ends of line charge distribution make angle $\theta_1$ and $\theta_2$ with the perpendicular line.

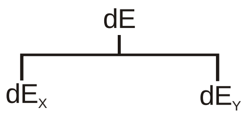



Consider a small element dx on line charge distribution at distance x from point A (see fig.). The charge of this element will be dq = $\lambda$dx. Due to this charge (dq), the intensity of electric field at the point P is dE.




Then $dE = \dfrac{K(dq)}{r^2+x^2} = \dfrac{K(\lambda dx)}{r^2+x^2}$

There will be two components of this field : $E_x = \int dE_x = \int dE\cos\theta = \int \dfrac{K\lambda dx}{r^2+x^2}.\cos\theta$




Assuming, $x = r\tan\theta \Rightarrow dx = r\sec^2\theta.\,d\theta$

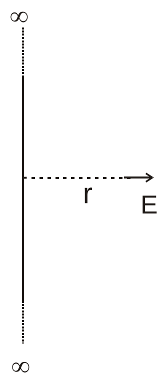
so $E_x = \displaystyle\int_{-\theta_2}^{+\theta_1} \frac{K\ \ \lambda r\sec^2\theta\cdot\cos\theta\cdot d\theta}{r^2+r^2\tan^2\theta} = \frac{K\lambda}{r}\int_{-\theta_2}^{+\theta_1}\cos\theta.d\theta = \frac{K\lambda}{r}\ [\sin\theta_1+\sin\theta_2]$ .....(1)

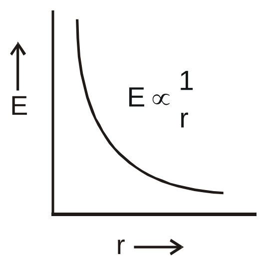
**Similarly y-component.** $E_y = \dfrac{K\lambda}{r}\displaystyle\int_{-\theta_2}^{+\theta_1}\sin\theta.d\theta = \frac{K\lambda}{r}\ [\cos\theta_2-\cos\theta_1]$

Net electric field at the point: $E_{net} = \sqrt{E_x^2+E_y^2}$

**(ii) We can derive a result for infinitely long line charge:** In above eq. (1) & (2), if we put $\theta_1 = \theta_2 = 90^o$, we can get required result.

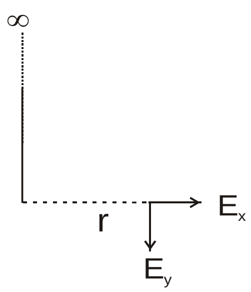
$E_{net} = E_x = \dfrac{2K\lambda}{r}$




**(iii) For Semi- infinite wire :** $\theta_1 = 90^o$ and $\theta_2 = 0^o$, so,

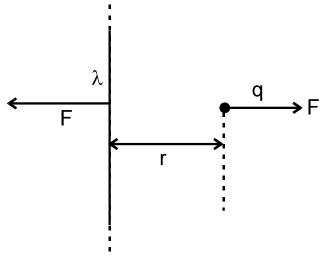
$E_x = \dfrac{K\lambda}{r}$ , $E_y = \dfrac{K\lambda}{r}$



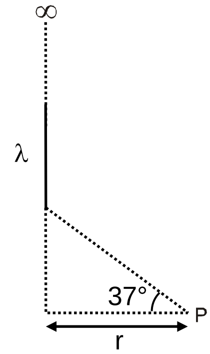
**Example 31.** A point charge q is placed at a distance r from a very long charged thread of uniform linear charge density $\lambda$. Find out total electric force experienced by the line charge due to the point charge. (Neglect gravity).

**Solution :** Force on charge q due to the thread, $F = \left(\frac{2K\lambda}{r}\right).q$

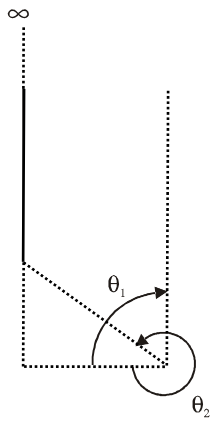
By Newton's III law, every action has equal and opposite reaction So, force on the thread $= \frac{2K\lambda}{r}.q$ (away from point charge)




**Example 32.** Figure shows a long wire having uniform charge density $\lambda$ as shown in figure. Calculate electric field intensity at point P.

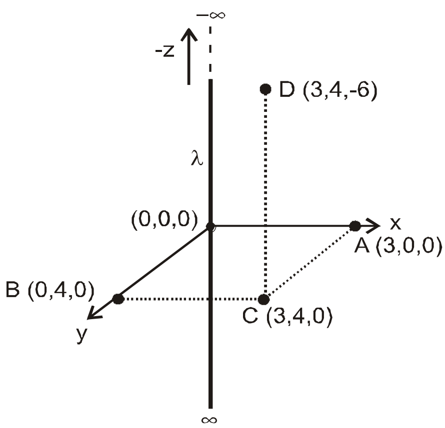



**Solution :** $\theta_1 = 90^o$ and $\theta_2 = 360^o - 37^o$ So




$E_x = \frac{K\lambda}{r}$ $[\sin\theta_1 + \sin\theta_2]$ ; $E_y = \frac{K\lambda}{r}$ $[\cos\theta_2 - \cos\theta_1]$

**Example 33.** Find electric field at point A, B, C, D due to infinitely long uniformly charged wire with linear charge density $\lambda$ and kept along z-axis (as shown in figure).Assume that all the parameters are in S.I. units.



**Solution :** $E_A = \frac{2K\lambda}{3}(\hat{i}) \Rightarrow E_B = \frac{2K\lambda}{4}(\hat{j})$

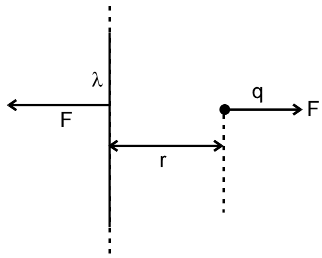
$E_C = \frac{2K\lambda}{5}\ \hat{OC} = \frac{2K\lambda}{5}\left(\frac{3\hat{i}+4\hat{j}}{5}\right) \qquad E_D = \frac{2K\lambda}{5}\left(\frac{3\hat{i}+4\hat{j}}{5}\right) \qquad \Rightarrow E_D = E_C$

**Example 34.** A point charge q is placed at a distance r from a very long charged thread of uniform linear charge density $\lambda$. Find out total electric force experienced by the line charge due to the point charge. (Neglect gravity).

**Solution :** Force on charge q due to the thread, $F = \left(\frac{2K\lambda}{r}\right).q$

By Newton's III law, every action has equal and opposite reaction, so force on the thread = $\frac{2K\lambda}{r}.q$

(away from point charge)

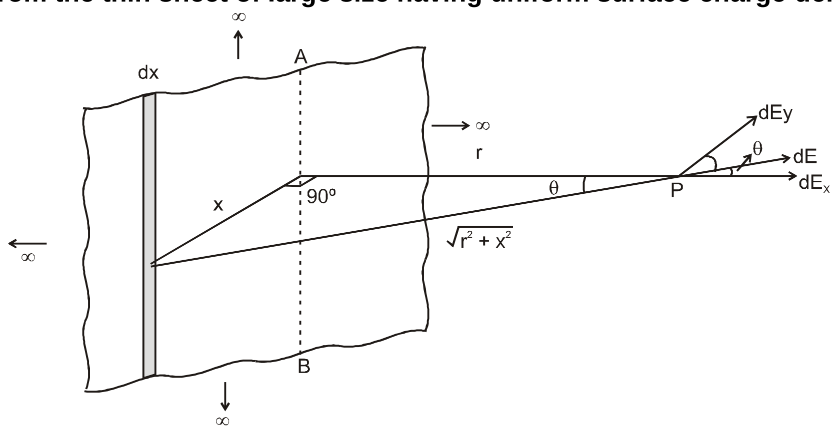



## 6.5 Electric field due to uniformly charged infinite sheet

$E_{net} = \frac{\sigma}{2\varepsilon_o}$ towards normal direction

### ELECTRIC FIELD DUE TO AN INFINITELY LARGE, UNIFORMLY CHARGED SHEET

**Derivation of expression for intensity of electric field at a point which is at a perpendicular distance r from the thin sheet of large size having uniform surface charge density $\sigma$.**




Assume a thin strip of width dx at distance x from line AB (see figure), which can be considered as a infinite line charge of charge density $\lambda = \sigma dx$

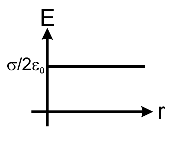
Due to this line charge the electric field intensity at point P will be $dE = \frac{\sigma K(dx)}{\sqrt{r^2+x^2}}$

Take another element similar to the first element on the other side of AB. Due to symmetry, Y-component of all such elements will be cancelled out.

So net electric field will be given by : $E_{net} = \int dE_x = \int dE\cos\theta = \int \frac{2K(\sigma dx)}{\sqrt{r^2+x^2}} \times \cos\theta$

Assume, $x = r\tan\theta \Rightarrow dx = r\sec^2\theta.\, d\theta$

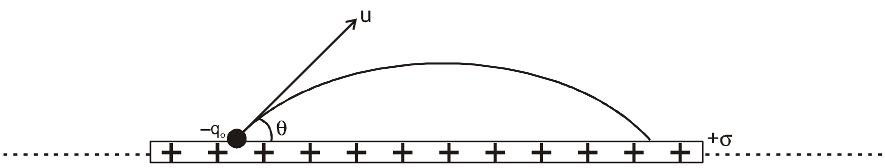
$\therefore \quad E_{net} = 2K\sigma \int_{-\pi/2}^{+\pi/2} \frac{r\sec^2\theta.d\theta.\cos\theta}{\sqrt{r^2+r^2\tan^2\theta}} = \frac{\sigma}{2\varepsilon_o}$ away from sheet




**Note : (1) The direction of electric field is always perpendicular to the sheet.**

**(2) The magnitude of electric field is independent of distance from sheet.**

**Example 35.** An infinitely large plate of surface charge density $+\sigma$ is lying in horizontal xy-plane. A particle having charge $-q_o$ and mass m is projected from the plate with velocity u making an angle $\theta$ with sheet. Find :

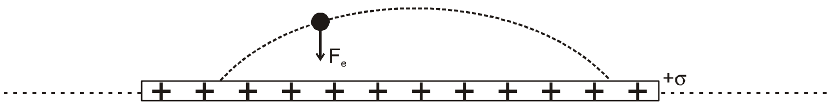



(i) The time taken by the particle to return on the plate..

(ii) Maximum height achieved by the particle.

(iii) At what distance will it strike the plate (Neglect gravitational force on the particle)

**Solution :**




Electric force acting on the particle $F_e = q_oE$ : $F_e = (q_o)\left(\dfrac{\sigma}{2\varepsilon_o}\right)$ downward

So, acceleration of the particle : $a = \dfrac{F_e}{m} = \dfrac{q_o\sigma}{2\varepsilon_o m}$ = uniform

This acceleration will act like 'g' (acceleration due to gravity)

So, the particle will perform projectile motion.

(i) $T = \dfrac{2u\sin\theta}{g} = \dfrac{2u\sin\theta}{\left(\dfrac{q_o\sigma}{2\varepsilon_o m}\right)}$ (ii) $H = \dfrac{u^2\sin^2\theta}{2g} = \dfrac{u^2\sin^2\theta}{2\left(\dfrac{q_o\sigma}{2\varepsilon_o m}\right)}$

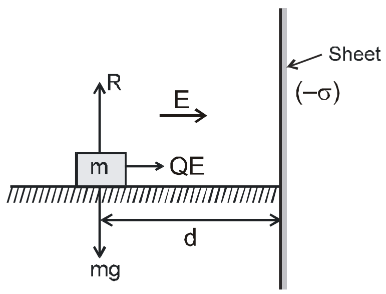
(iii) $R = \dfrac{u^2\sin 2\theta}{g} = \dfrac{u^2\sin 2\theta}{\left(\dfrac{q_o\sigma}{2\varepsilon_o m}\right)}$

**Example 36.** A block having mass m and charge Q is resting on a frictionless plane at a distance d from fixed large non-conducting infinite sheet of uniform charge density $-\sigma$ as shown in Figure. Assuming that collision of the block with the sheet is perfectly elastic, find the time period of oscillatory motion of the block. Is it SHM?

**Solution :** The situation is shown in Figure. Electric force produced by sheet will accelerate the block towards the sheet producing an acceleration. Acceleration will be uniform because electric field E due to the sheet is uniform.

$a = \dfrac{F}{m} = \dfrac{QE}{m}$, where $E = \sigma/2\varepsilon_0$

As initially the block is at rest and acceleration is constant, from second equation of motion, time taken by the block to reach the wall




$d = \dfrac{1}{2}at^2$ i.e., $t = \sqrt{\dfrac{2L}{a}} = \sqrt{\dfrac{2md}{QE}} = \sqrt{\dfrac{4md\varepsilon_0}{Q\sigma}}$

As collision with the wall is perfectly elastic, the block will rebound with same speed and as now its motion is opposite to the acceleration, it will come to rest after traveling same distance d in same time t.

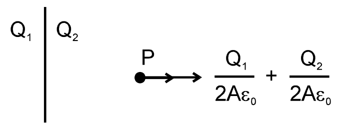
After stopping, it will again be accelerated towards the wall and so the block will execute oscillatory motion with 'span' d and time period.

$$T = 2t = 2\sqrt{\frac{2md}{QE}} = 2\sqrt{\frac{4md\varepsilon_0}{Q\sigma}}$$

However, as the restoring force F = QE is constant and not proportional to displacement x, the motion is not simple harmonic.

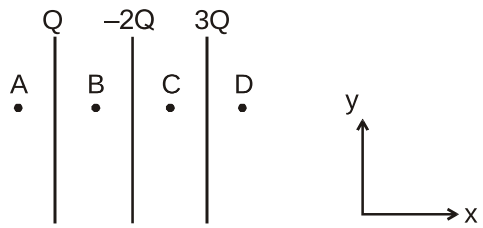
**Example 37.** If an isolated infinite sheet contains charge $Q_1$ on its one surface and charge $Q_2$ on its other surface then prove that electric field intensity at a point in front of sheet will be $\frac{Q}{2A\varepsilon_O}$, where $Q = Q_1 + Q_2$

**Solution :** Electric field at point P : $\vec{E} = \vec{E}_{Q_1} + \vec{E}_{Q_2} = \frac{Q_1}{2A\varepsilon_0}\hat{n} + \frac{Q_2}{2A\varepsilon_0}\hat{n} = \frac{Q_1 + Q_2}{2A\varepsilon_0}\hat{n} = \frac{Q}{2A\varepsilon_0}\hat{n}$

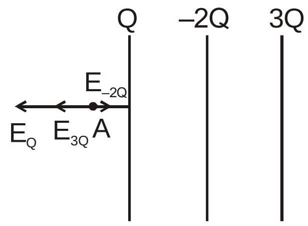



[This shows that the resultant field due to a sheet depends only on the total charge of the sheet and not on the distribution of charge on individual surfaces].

**Example 38.** Three large conducting parallel sheets are placed at a finite distance from each other as shown in figure. Find out electric field intensity at points A, B, C & D.

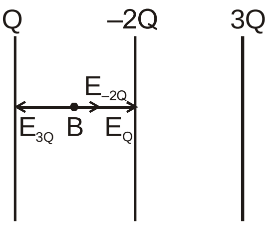



**Solution :** **For point A :**



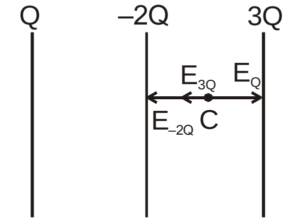

$$\vec{E}_{net} = \vec{E}_Q + \vec{E}_{3Q} + \vec{E}_{-2Q} = -\frac{Q}{2A\varepsilon_0}\hat{i} - \frac{3Q}{2A\varepsilon_0}\hat{i} + \frac{2Q}{2A\varepsilon_0}\hat{i} = -\frac{Q}{A\varepsilon_0}\hat{i}$$

**For point B:**



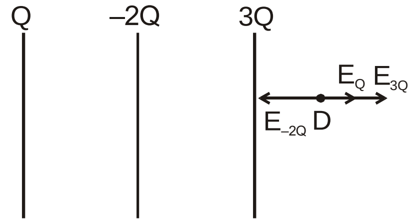

$$\vec{E}_{net} = \vec{E}_{3Q} + \vec{E}_{-2Q} + \vec{E}_Q = -\frac{3Q}{2A\varepsilon_0}\hat{i} + \frac{2Q}{2A\varepsilon_0}\hat{i} + \frac{Q}{2A\varepsilon_0}\hat{i} = \vec{0}$$

**For point C :**




$$\vec{E}_{net} = \vec{E}_Q + \vec{E}_{3Q} + \vec{E}_{-2Q} = +\frac{Q}{2A\varepsilon_0}\hat{i} - \frac{3Q}{2A\varepsilon_0}\hat{i} - \frac{Q}{2A\varepsilon_0}\hat{i} = -\frac{2Q}{A\varepsilon_0}\hat{i}$$

**For point D :**



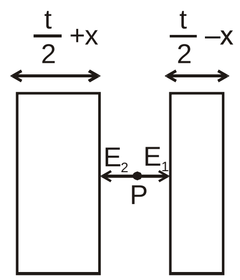

$$\vec{E}_{net} = \vec{E}_Q + \vec{E}_{3Q} + \vec{E}_{-2Q} = +\frac{Q}{2A\varepsilon_0}\hat{i} + \frac{3Q}{2A\varepsilon_0}\hat{i} - \frac{2Q}{2A\varepsilon_0}\hat{i} = \frac{2Q}{A\varepsilon_0}\hat{i}$$

**Example 39.** Determine and draw the graph of electric field due to infinitely large non-conducting sheet of thickness t and uniform volume charge density $\rho$ as a function of distance x from its symmetry plane.

(a) $x \le \frac{t}{2}$ (b) $x \ge \frac{t}{2}$

**Solution :** We can consider two sheets of thickness $\left(\frac{t}{2} - x\right)$ and $\left(\frac{t}{2} + x\right)$

Where the point P lies inside the sheet.

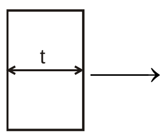



Now, net electric field at point P :

$E = E_1 - E_2 = \frac{Q_1}{2A\varepsilon_0} - \frac{Q_2}{2A\varepsilon_0}$ [$Q_1$ : charge of left sheet; $Q_2$ : charge of right sheet.]

$$= \frac{A\rho\left(\frac{t}{2} + x\right) - \rho A\left(\frac{t}{2} - x\right)}{2A\varepsilon_0} = \frac{\rho x}{\varepsilon_0}$$

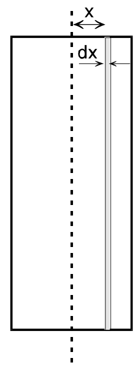
For point which lies outside the sheet we can consider a complete sheet of thickness t




$$\frac{Q}{2A\varepsilon_0}$$

$$E = \frac{\sigma t A}{2A\varepsilon_0} = \frac{\sigma t}{2\varepsilon_0}$$

**Alternate :** We can assume thick sheet to be made of large number of uniformly charged thin sheets. Consider an elementary thin sheet of width dx at a distance x from symmetry plane.



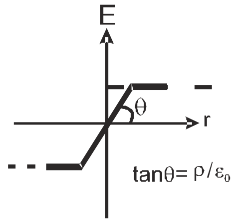

Charge in sheet = $\rho A dx$ (A : assumed area of sheet)

Surface charge density, $\sigma = \frac{\rho A dx}{A}$

so, electric field intensity due to elementary sheet : $dE = \frac{\rho dx}{2\varepsilon_0}$

(a) When $x < \frac{t}{2} \Rightarrow$ $E_{Net} = \int_{-t/2}^{x} \frac{\rho dx}{2\varepsilon_0} - \int_{x}^{t/2} \frac{\rho dx}{2\varepsilon_0} = \frac{\rho x}{\varepsilon_0}$

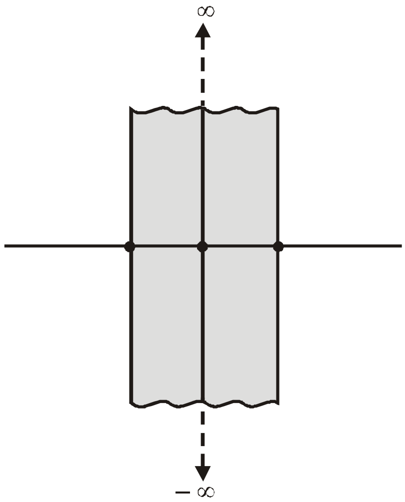
(b) When $x > \frac{t}{2} \Rightarrow$ $E_{Net} = \int_{-t/2}^{t/2} \frac{\rho dx}{2\varepsilon_0} = \frac{\rho t}{2\varepsilon_0}$




**Example 40.** Thin infinite sheet of width w contains uniform charge distribution σ. Find out electric field intensity at following points :

(a) A point which lies in the same plane at a distance d from one of its edge.

(b) A point which is on the symmetry plane of sheet at a perpendicular distance d from it.



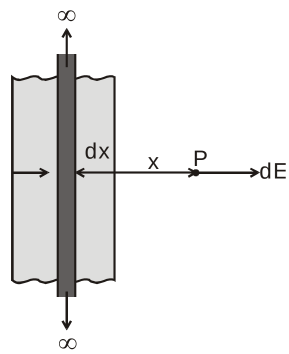

**Solution :** (a) Consider a thin strip of width dx. Linear charge density of strip : $\lambda = \sigma dx$

So, electric field due to this strip at point P

$$dE = \frac{2k\sigma dx}{x}$$

$$E_{net} = \int_{d}^{d+w} \frac{\sigma}{2\pi\varepsilon_0}\frac{dx}{x}$$

$$= \frac{\sigma}{2\pi\varepsilon_0}\ell n\left(\frac{d+w}{d}\right)$$



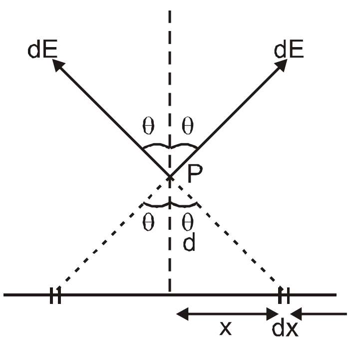

(b) Consider a thin strip of width dx. Linear charge density of strip :

$\lambda = \sigma dx$

$\therefore$ $E_p = \int 2dE\cos\theta$

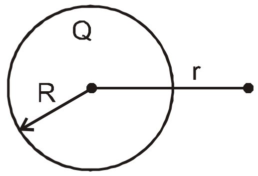
or $E_p = 2.\int_0^{w/2} \frac{\sigma dx}{2\pi\varepsilon_0\sqrt{d^2+x^2}}.\frac{d}{\sqrt{d^2+x^2}}$

$$= \frac{\sigma d}{\pi\varepsilon_0}\int_0^{w/2}\frac{dx}{d^2+x^2} = \frac{\sigma}{\pi\varepsilon_0}\tan^{-1}\frac{w}{2d}$$



## 6.6 Electric field due to uniformly charged spherical shell

$E = \frac{KQ}{r^2}$ $r \geq R$ $\Rightarrow$ For the outside points & point on the surface the uniformly charged spherical shell behaves as a point charge placed at the centre

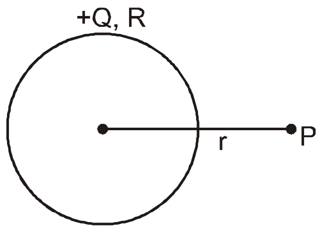



$E = 0$ $r < R$

Electric field due to spherical shell outside it is always along the radial direction.

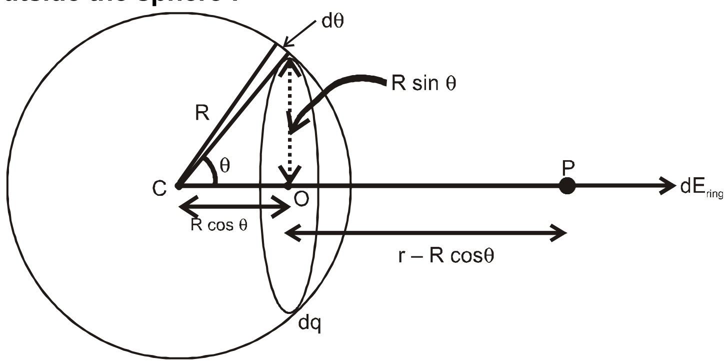
### Finding electric field due to a uniformly charged spherical shell :

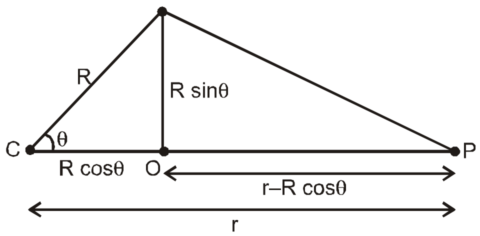
Suppose we have a spherical shell of radius R and charge +Q uniformly distributed on its surface. We have to find electric field at a point P, which is at a distance 'r' from the centre of the sphere.




For this, we can divide the shell into thin rings. Let's consider a ring making an angle $\theta$ with the axis and subtending a small angle $d\theta$. Its width will be $Rd\theta$. (arc = radius × angle = $Rd\theta$).

**For the points outside the sphere :**




Electric field due to this small ring element :

$$dE = \frac{Kdqx}{[(\text{ring radius})^2 + x^2]^{3/2}} \quad .....(1)$$

So, total electric field $E_{net} = \int \frac{Kdq\,x}{[(\text{ring radius})^2 + x^2]^{3/2}}$

Here, radius of the ring element = R $\sin\theta$ & x = axial distance of point P from the ring = r – R cos $\theta$

Area of the ring element = (length) (width) = ($2\pi$ (radius of the ring) ) $Rd\theta$ = ($2\pi R \sin\theta$) $Rd\theta$

dq = charge of the small ring element. We can find dq by unitary method.

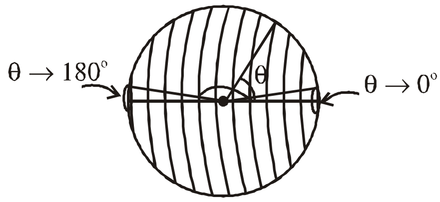
In $4\pi R^2$ Area, charge is Q

In unit Area, charge is $\frac{Q}{4\pi R^2}$ .

In ($2\pi R \sin\theta$) $Rd\theta$ Area, charge = $\frac{Q}{4\pi R^2} \times (2\pi R \sin\theta)\, Rd\theta = dq$

Putting values of r and dq in equation ..(1)

We get $E_{out} = \displaystyle\int_{\theta=0}^{\theta=\pi} \frac{K\left(\frac{Q}{4\pi R^2}\times 2\pi(R\sin\theta)Rd\theta\right)(r-R\cos\theta)}{\left[(R\sin\theta)^2+(r-R\cos\theta)^2\right]^{3/2}}$




(The first ring will make angle $\theta$ = 0 and the last ring will make $\theta$ = 180º. So, limit will be from $\theta$ = 0 to $\theta$ = 180º)
Steps of integration : From above integral :

$$E_{out} = \frac{KQ}{2}\int_0^{\pi}\frac{(r-R\cos\theta)\sin\theta d\theta}{(R^2+r^2-2Rr\cos\theta)^{3/2}}$$

Now, let $z^2 = R^2 + r^2 - 2Rr\cos\theta \Rightarrow \quad 2zdz = 0 + 0 - 2Rr(-\sin\theta)\,d\theta$

$\therefore \quad zdz = Rr\sin\theta\, d\theta \quad \& \quad \cos\theta = \dfrac{R^2+r^2-z^2}{2Rr}$

Now, when $\theta = 0 \quad \rightarrow \quad z = (r - R)$

$\theta = \pi \quad \rightarrow \quad z = r + R$

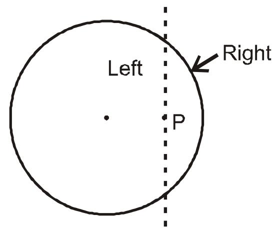
$$\therefore \quad E_{out} = \frac{KQ}{2}\int_{r-R}^{r+R}\frac{\left[r-R\frac{(R^2+r^2-z^2)}{2Rr}\right]\frac{zdz}{Rr}}{z^3} = \frac{KQ}{2}\int_{r-R}^{r+R}\frac{(2Rr^2-R^3-Rr^2+Rz^2)dz}{2R^2r^2z^2}$$

$$= \frac{KQ}{4R^2r^2}\left[\int_{r-R}^{r+R}\frac{Rr^2dz}{z^2}-\int_{r-R}^{r+R}\frac{R^3dz}{z^2}+\int_{r-R}^{r+R}Rdz\right] = \frac{KQ}{4R^2r^2}\left[-\frac{Rr^2}{z}+\frac{R^3}{z}+Rz\right]_{r-R}^{r+R}$$

On solving above, we will get : $E_{out} = \dfrac{KQ}{r^2}$ if r > R

**For the points inside the sphere:**

Now lets derive the electric field due to uniformly charged solid sphere at a point 'P' inside it. The sphere is divided into two parts, the rings on the left part of point 'P' will produce electric field towards right and the rings on right part will produce electric field towards left and $E_{net} = E_{right} - E_{left}$. For this, limit of integration is divided into two parts.




$$E_{net} = \int_{\theta=0}^{\theta=\cos^{-1}\left(\frac{r}{R}\right)}\begin{pmatrix}\text{Electric}\\ \text{field due}\\ \text{to rings}\\ \text{of right}\\ \text{part}\end{pmatrix} - \int_{\theta=\cos^{-1}\left(\frac{r}{R}\right)}^{\theta=\pi}\begin{pmatrix}\text{Electric}\\ \text{field due}\\ \text{to rings}\\ \text{of left}\\ \text{part}\end{pmatrix}$$

As $\quad z^2 = R^2 + r^2 - 2rR\cos\theta$

When $\quad \theta = \cos^{-1}\left(\dfrac{r}{R}\right) \quad \Rightarrow \quad z = \sqrt{R^2-r^2}$

When $\quad \theta = 0 \quad \Rightarrow \quad z = R - r$

When $\quad \theta = \pi \quad \Rightarrow \quad z = R + r$

From the result of previous case and just by changing limits we can write

$$E_{in} = \left[-\frac{Rr^2}{z} + \frac{R^3}{z} + Rz\right]_{R-r}^{\sqrt{R^2-r^2}} - \left[-\frac{Rr^2}{z} + \frac{R^3}{z} + Rz\right]_{\sqrt{R^2-r^2}}^{R+r}$$

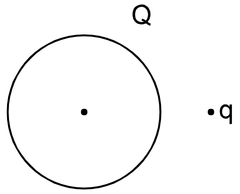
On solving this expression, we will get E and E = 0 if r < R.

Finding electric field due to shell by integration is very lengthy, so we will not use this method. The given hand-out was just for knowledge. The best method to find E due to shell is by Gauss theorem which we will study later.

**Example 41.** Figure shows a uniformly charged sphere of radius R and total charge Q. A point charge q is situated outside the sphere at a distance r from centre of sphere. Find out the following :

(i) Force acting on the point charge q due to the sphere.

(ii) Force acting on the sphere due to the point charge.




**Solution :** (i) Electric field at the position of point charge $\vec{E} = \frac{KQ}{r^2}\hat{r}$

So, $\vec{F} = \frac{KqQ}{r^2}\hat{r}$ $\quad |\vec{F}| = \frac{KqQ}{r^2}$

(ii) Since we know that every action has equal and opposite reaction so $\vec{F}_{sphere} = \frac{KqQ}{r^2}\hat{r}$

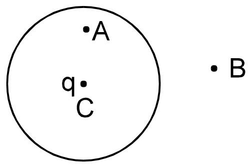
$\vec{F}_{sphere} = \frac{KqQ}{r^2}$

**Example 42.** Figure shows a uniformly charged thin sphere of total charge Q and radius R. A point charge q is also situated at the centre of the sphere. Find out the following :

(i) Force on charge q

(ii) Electric field intensity at A.

(iii) Electric field intensity at B.




**Solution :** (i) Electric field at the centre of the uniformly charged hollow sphere = 0

So force on charge q = 0

(ii) Electric field at A

$\vec{E}_A = \vec{E}_{Sphere} + \vec{E}_q = 0 + \frac{Kq}{r^2}$ ; r = CA

E due to sphere = 0 , because point lies inside the charged hollow sphere.

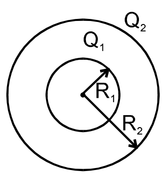
(iii) Electric field $\vec{E}_B$ at point B $= \vec{E}_{Sphere} + \vec{E}_q = \frac{KQ}{r^2}.\hat{r} + \frac{Kq}{r^2}.\hat{r} = \frac{K(Q+q)}{r^2}.\hat{r}$ ; r = CB

**Note :** Here, we can also assume that the total charge of sphere is concentrated at the centre, for calculation of electric field at B.

**Example 43.** Two concentric uniformly charged spherical shells of radius $R_1$ and $R_2$ ($R_2 > R_1$) have total charges $Q_1$ and $Q_2$ respectively. Derive an expression of electric field as a function of r for following positions.

(i) $r < R_1$ (ii) $R_1 \le r < R_2$

(iii) $r \ge R_2$




**Solution :** (i) For $r < R_1$, therefore, point lies inside both the spheres

$E_{net} = E_{Inner} + E_{outer} \quad = 0 + 0$

(ii) For $R_1 \le r < R_2$, point lies outside inner sphere but inside outer sphere :

$\therefore \quad E_{net} = E_{inner} + E_{outer}$

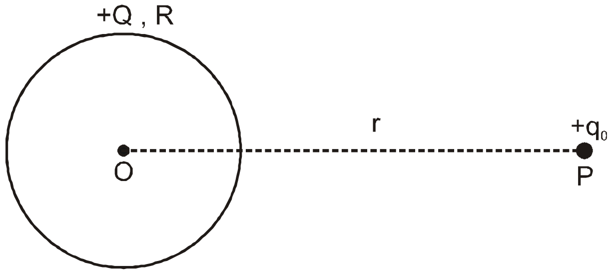
$$\frac{KQ_1}{r^2}\hat{r} + 0 = \frac{KQ_1}{r^2}\hat{r}$$

(iii) For $r \ge R_2$

point lies outside inner as well as outer sphere.

Therefore, $E_{Net} = E_{inner} + E_{outer} = \frac{KQ_1}{r^2}\hat{r} + \frac{KQ_2}{r^2}\hat{r} = \frac{K(Q_1+Q_2)}{r^2}\hat{r}$

**Example 44.** A spherical shell having charge +Q (uniformly distributed) and a point charge + $q_0$ are placed as shown. Find the force between shell and the point charge($r >> R$).

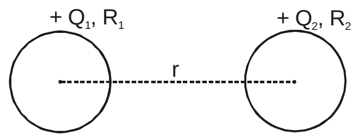



(i) Force on the point charge + $q_0$ due to the shell = $q\vec{E}_{0shell} = (q_0)\left(\frac{KQ}{r^2}\right)\hat{r} = \frac{KQq_0}{r^2}\hat{r}$ where $\hat{r}$, is unit vector along OP.

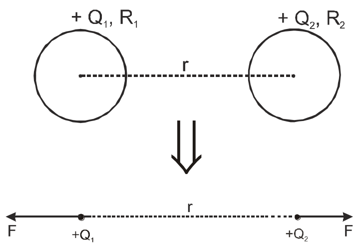
From action - reaction principle, force on the shell due to the point charge will also be : $F_{shell} = \frac{KQq_0}{r^2}(-\hat{r})$

**Conclusion :** To find the force on a hollow sphere due to outside charges, we can replace the sphere by a point charge kept at centre.

**Example 45.** Find force acting between two shells of radius $R_1$ and $R_2$ which have uniformly distributed charges $Q_1$ and $Q_2$ respectively and distance between their centers is r.




**Solution :** The shells can be replaced by point charges kept at centre, so force between them

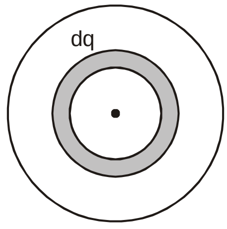



$$F = \frac{KQ_1Q_2}{r^2}$$

## 6.7 Electric field due to uniformly charged solid sphere:

**Derive an expression for electric field due to solid sphere of radius R and total charge Q which is uniformly distributed in the volume, at a point which is at a distance r from centre for given two cases.**

**(i) $r \ge R$ (ii) $r \le R$**

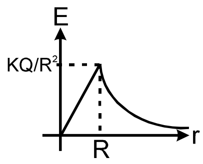



Assume an elementary concentric shell of charge dq. Due to this shell, the electric field at the point (r > R) will be:

$dE = \dfrac{Kdq}{r^2}$ [from above result of hollow sphere]

$E_{net} = \int dE = \dfrac{KQ}{r^2}$

For r < R, there will be no electric field due to shell of radius greater than r, so electric field at the point will be present only due to shells having radius less than r.




$E'_{net} = \dfrac{KQ'}{r^2}$

Here, $Q' = \dfrac{Q}{\frac{4}{3}\pi R^3} \times \dfrac{4}{3}\pi r^3 = \dfrac{Qr^3}{R^3}$

$\therefore$ $E'_{net} = \dfrac{KQ'}{r^2} = \dfrac{KQr}{R^3}$ ; away from the centre.

**Note : The electric field inside and outside the sphere is always in radial direction.**

**Example 46.** A solid non conducting sphere of radius R and uniform volume charge density $\rho$ has its centre at origin. Find out electric field intensity in vector form at following positions :

(i) (R/2, 0, 0) (ii) $\left(\dfrac{R}{\sqrt{2}}, \dfrac{R}{\sqrt{2}}, 0\right)$ (iii) (R, R, 0)

**Solution :** (i) At (R/2, 0, 0) : Distance of point from centre = $\sqrt{(R/2)^2 + 0^2 + 0^2}$ = R/2 < R, so point lies inside the sphere, so $\vec{E} = \dfrac{p\vec{r}}{3\varepsilon_0} = \dfrac{\rho}{3\varepsilon_0}\left[\dfrac{R}{2}\hat{i}\right]$

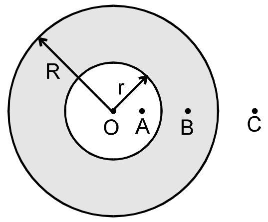
(ii) At $\left(\dfrac{R}{\sqrt{2}}, \dfrac{R}{\sqrt{2}}, 0\right)$ distance of point from centre $= \sqrt{(R/\sqrt{2})^2 + (R/\sqrt{2})^2 + 0^2}$ = R = R, so point lies at the surface of sphere, therefore

$$\vec{E} = \frac{KQ}{R^3}\vec{r} = \frac{K\frac{4}{3}\pi R^3 \rho}{R^3} = \left[\frac{R}{\sqrt{2}}\hat{i} + \frac{R}{\sqrt{2}}\hat{j}\right] = \frac{\rho}{3\varepsilon_0}\left[\frac{R}{\sqrt{2}}\hat{i} + \frac{R}{\sqrt{2}}\hat{j}\right]$$

(iii) The point is outside the sphere

So, $\vec{E} = \dfrac{KQ}{r^3}\vec{r} = \dfrac{K\frac{4}{3}\pi R^3 \rho}{(\sqrt{2}R)^3}\left[R\hat{i} + R\hat{j}\right] = \dfrac{\rho}{6\sqrt{2}\varepsilon_0}\left[R\hat{i} + R\hat{j}\right]$

**Example 47.** A Uniformly charged solid non-conducting sphere of uniform volume charge density $\rho$ and radius R is having a concentric spherical cavity of radius r. Find out electric field intensity at following points, as shown in the figure :




(i) Point A (ii) Point B

(iii) Point C (iv) Centre of the sphere

**Solution :** **Method-I :**

(i) For point A : We can consider the solid part of sphere to be made of large number of spherical shells which have uniformly distributed charge on its surface. Now, since point A lies inside all spherical shells so electric field intensity due to all shells will be zero.

$\vec{E}_A = 0$

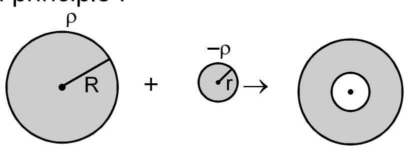
(ii) For point B : All the spherical shells for which point B lies inside will make electric field zero at point B. So electric field will be due to charge present from radius r to OB.

$$\text{So, } \vec{E}_B = \frac{K\frac{4}{3}\pi(OB^3 - r^3)\rho}{OB^3}\,\overrightarrow{OB} = \frac{\rho}{3\varepsilon_0}\frac{[OB^3 - r^3]}{OB^3}\,\overrightarrow{OB}$$

(iii) For point C, similarly we can say that for all the shell points C lies outside the shell

$$\text{So, } \vec{E}_C = \frac{K[\frac{4}{3}\pi(R^3 - r^3)]}{[OC]^3}\,\overrightarrow{OC} = \frac{\rho}{3\varepsilon_0}\frac{R^3 - r^3}{[OC]^3}\,\overrightarrow{OC}$$

**Method-II :** We can consider that the spherical cavity is filled with charge density $\rho$ and also $-\rho$, thereby making net charge density zero after combining. We can consider two concentric solid spheres: One of radius R and charge density $\rho$ and other of radius r and charge density $-\rho$. Applying superposition principle :




(i) $$\vec{E}_A = \vec{E}_\rho + \vec{E}_{-\rho} = \frac{\rho(\overrightarrow{OA})}{3\varepsilon_0} + \frac{[-\rho(\overrightarrow{OA})]}{3\varepsilon_0} = 0$$

(ii) $$\vec{E}_B = \vec{E}_\rho + \vec{E}_{-\rho} = \frac{\rho(\overrightarrow{OB})}{3\varepsilon_0} + \frac{K\left[\frac{4}{3}\pi r^3(-\rho)\right]}{(OB)^3}\overrightarrow{OB}$$

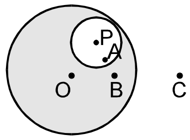
$$= \left[\frac{\rho}{3\varepsilon_0} - \frac{r^3\rho}{3\varepsilon_0(OB)^3}\right]\overrightarrow{OB} = \frac{\rho}{3\varepsilon_0}\left[1 - \frac{r^3}{OB^3}\right]\overrightarrow{OB}$$

(iii) $$\vec{E}_C = \vec{E}_\rho + \vec{E}_{-\rho} = \frac{K\left(\frac{4}{3}\pi R^3\rho\right)}{OC^3}\overrightarrow{OC} + \frac{K\left(\frac{4}{3}\pi r^3(-\rho)\right)}{OC^3}\overrightarrow{OC} = \frac{\rho}{3\varepsilon_0(OC)^3}[R^3 - r^3]\overrightarrow{OC}$$

(iv) $$\vec{E}_O = \vec{E}_\rho + \vec{E}_{-\rho} = 0 + 0 = 0$$

**Example 48.** In above question, if cavity is not concentric and centered at point P then repeat all the steps.

**Solution :** Again assume $\rho$ and $-\rho$ in the cavity, (similar to the previous example) :




(i) $\vec{E_A} = \vec{E_\rho} + \vec{E_{-\rho}} = \frac{\rho[\vec{OA}]}{3\varepsilon_0} + \frac{(-\rho)\vec{PA}}{3\varepsilon_0}$

$\frac{\rho}{3\varepsilon_0}[\vec{OA} - \vec{PA}] = \frac{\rho}{3\varepsilon_0} \quad [\vec{OP}]$

**Note :** Here, we can see that the electric field intensity at point A is independent of position of point A inside the cavity. Also the electric field is along the line joining the centres of the sphere and the spherical cavity.

(ii) $\vec{E_B} = \vec{E_\rho} + \vec{E_{-\rho}} = \frac{\rho(\vec{OB})}{3\varepsilon_0} + \frac{K[\frac{4}{3}\pi r^3(-\rho)]}{[PB]^3}\ \vec{PB}$

(iii) $\vec{E_C} = \vec{E_\rho} + \vec{E_{-\rho}} = \frac{K[\frac{4}{3}\pi R^3\rho]}{[OC]^3}\ \vec{OC} + \frac{K[\frac{4}{3}\pi r^3(-\rho)]}{[PC]^3}\ \vec{PC}$

(iv) $\vec{E_O} = \vec{E_\rho} + \vec{E_{-\rho}} = 0 + \frac{K[\frac{4}{3}\pi r^3(-\rho)]}{[PO]^3}\ \vec{PO}$

**Example 49.** A non-conducting solid sphere has volume charge density that varies as $\rho = \rho_0\, r$, where $\rho_0$ is a constant and r is distance from centre. Find out electric field intensities at following positions.

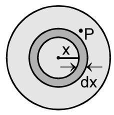
(i) $r < R$ (ii) $r \geq R$

**Solution :** **Method I :**

(i) For $r < R$ :

The sphere can be considered to be made of large number of spherical shells. Each shell has uniform charge density on its surface. So the previous results of the spherical shell can be used. Consider a shell of radius x and thickness dx as an element. Charge on shell $dq = (4\pi x^2 dx)\rho_0 x$.

$\therefore$ Electric field intensity at point P due to shell, $dE = \frac{Kdq}{x^2}$




Since all the shell will have electric field in same direction

$$\therefore \quad E = \int_0^R dE = \int_0^r dE + \int_r^R dE$$

Due to shells which lie between region $r < x \leq R$, electric field at point P will be zero.

$$\therefore \quad |\vec{E}| = \int_0^r \frac{Kdq}{r^2} + 0 = \int_0^r \frac{K.4\pi x^2 dx \rho_0 x}{r^2} = \frac{4\pi K\rho_0}{r^2}\left[\frac{x^4}{4}\right]_0^r = \frac{\rho_0 r^2}{4\varepsilon_0}$$

(ii) For $r \geq R$, $\vec{E} = \int_0^R dE = \int_0^R \frac{K.4\pi x^2 dx \rho_0 x}{r^2} = \frac{\rho_0 R^4}{4\varepsilon_0 r^2}\hat{r}$

**Method II :**

(i) The sphere can be considered to be made of large number of spherical shells. Each shell has uniform charge density on its surface. So the previous results of the spherical shell can be used. we can say that all the shells for which point lies inside will make electric field zero at that point,

$$\text{So} \quad \vec{E}_{(r<R)} = \frac{K\int_0^r (4\pi x^2 dx)\rho_0 x}{r^2} = \frac{\rho_0 r^2}{4\varepsilon_0}\hat{r}$$

(ii) Similarly, for $r \geq R$, all the shells will contribute in electric field. Therefore :

$$\vec{E}_{(r<R)} = \frac{K\int_0^R (4\pi x^2 dx)\rho_0 x}{r^2} = \frac{\rho_0 R^4}{4\varepsilon_0 r^2}\hat{r}$$

## 7. ELECTRIC POTENTIAL :

In electrostatic field, the electric potential (due to some source charges) at a point P is defined as the work done by external agent in taking a unit point positive charge from a reference point (generally taken at infinity) to that point P without changing its kinetic energy.

### 7.1 Mathematical representation :

If $(W_{\infty \to P})_{ext}$ is the work required in moving a point charge q from infinity to a point P, the electric potential of the point P is

$$V_p = \left.\frac{W_{\infty\to p})_{ext}}{q}\right]_{\Delta K=0} = \frac{(-W_{elc})_{\infty\to p}}{q}$$

**Note :** **(i)** $(W_{\infty \to P})_{ext}$ can also be called as the work done by external agent against the electric force on a unit positive charge due to the source charge.

**(ii)** Write both W and q with proper sign.

### 7.2 Properties :

**(i)** Potential is a scalar quantity, its value may be positive, negative or zero.

**(ii)** S.I. Unit of potential is volt = $\frac{\text{joule}}{\text{coulmb}}$ and its dimensional formula is $[M^1L^2T^{-3}I^{-1}]$.

**(iii)** Electric potential at a point is also equal to the negative of the work done by the electric field in taking the point charge from reference point (i.e. infinity) to that point.

**(iv)** Electric potential due to a positive charge is always positive and due to negative charge it is always negative except at infinity. (Taking $V_\infty = 0$).

**(v)** Potential decreases in the direction of electric field.

**(vi)** $V = V_1 + V_2 + V_3 + .......$

### 7.3 Use of potential :

If we know the potential at some point (in terms of numerical value or in terms of formula) then we can find out the work done by electric force when charge moves from point 'P' to $\infty$ by the formula

$W_{ep})_{p\to\infty} = qV_p$

**Example 50** A charge 2μC is taken from infinity to a point in an electric field, without changing its velocity. If work done against electrostatic forces is – 40μJ, then find the potential at that point.

**Solution :** $V = \frac{W_{ext}}{q} = \frac{-40\mu J}{2\mu C} = -20\text{ V}$

**Example 51** When charge 10 $\mu$C is shifted from infinity to a point in an electric field, it is found that work done by electrostatic forces is –10 $\mu$J. If the charge is doubled and taken again from infinity to the same point without accelerating it, then find the amount of work done by electric field and against electric field.

**Solution :** $W_{ext})_{\infty \to p} = -W_{el})_{\infty \to p} = W_{el})_{p \to \infty} = 10\ \mu J$

because $\Delta KE = 0$

$\therefore \quad V_p = \dfrac{(W_{ext})_{\infty \to p}}{20\mu C} = \dfrac{10\mu J}{10\mu C} = 1V$

So, if now the charge is doubled and taken from infinity then

$1 = \dfrac{w_{ext})_{\infty \to p}}{20\mu C} \quad \Rightarrow \quad$ or $\quad W_{ext})_{\infty \to P} = 20\ \mu J$

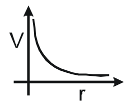
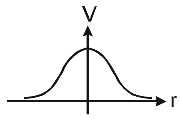
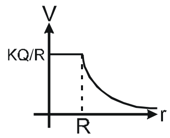
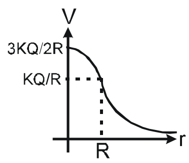
$\Rightarrow \quad W_{el})_{\infty \to P} = -20\ \mu J$

**Example 52** A charge 3$\mu$C is released from rest from a point P where electric potential is 20 V then its kinetic energy when it reaches infinity is :

**Solution :** $W_{el} = \Delta K = K_f - 0$

$\therefore \quad W_{el})_{P \to \infty} = qV_P = 60\ \mu J \qquad$ So, $K_f = 60\ \mu J$

## Electric Potential due to various charge distributions are given in table.

| Name / Type | Formula | Note | Graph |
|---|---|---|---|
| Point charge | $\frac{Kq}{r}$ | * q is source charge.<br>* r is the distance of the point from the point charge. |   |
| Ring (uniform/nonuniform charge distribution) | at centre: $\frac{KQ}{R}$<br>at the axis: $\frac{KQ}{\sqrt{R^2+x^2}}$ | * Q is source chage.<br>* x is the distance of the point on the axis from centre f ring |   |
| Uniformly charged hollow conducting/nonconducting /solid conducting sphere | for $r \geq R$, $V = \frac{kQ}{r}$<br>for $r \leq R$, $V = \frac{kQ}{R}$ | * R is radius of sphere<br>* r is the distance from centre of sphere to the point<br>* Q is total charge = $\sigma 4\pi R^2$. |   |
| Uniformly charged solid nonconducting sphere . | For $r \geq R$, $V = \frac{kQ}{r}$<br>For $r \leq R$<br>$V = \frac{KQ(3R^2 - r^2)}{2R^3}$<br>$= \frac{\rho}{6\varepsilon_0}(3R^2 - r^2)$ | * R is radius of sphere<br>* r is distance from centre to the point<br>* $V_{centre} = \frac{3}{2} V_{surface}$.<br>* Q is total charge = $\rho \frac{4}{3}\pi R^3$.<br>* Inside the sphere potential varies parabolically<br>* Outside the sphere potential varies hyperbolically. |   |
| Infinite line charge | Not defined | * Absolute potential is not defined.<br>* Potential difference between two points is given by formula : $v_B - v_A = -2K\lambda \ \ln(r_B / r_A)$ | |
| Infinite nonconducting thin sheet | Not defined | * Absolute potential is not defined.<br>* Potential difference between two points is given by formula<br>$V_B - V_A = -\frac{\sigma}{2\varepsilon_0}(r_B - r_A)$ | |
| Infinite charged conducting thin sheet | Not defined | * Absolute potential is not defined.<br>* Potential difference between two points is given by formula<br>$V_B - V_A = -\frac{\sigma}{\varepsilon_0}(r_B - r_A)$ | |

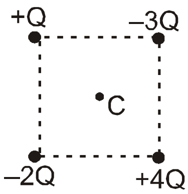
## 7.4 Potential due to a point charge :

Derivation of expression for potential due to point charge Q, at a point which is at a distance r from the point charge.


From definition of potential

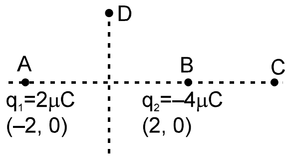
$$V = \frac{W_{ext(\infty \to p)}}{q_o} = \frac{-\int_{\infty}^{r} (q_0\vec{E}) \cdot d\vec{r}}{q_o} = -\int_{\infty}^{r} \vec{E} \cdot \vec{dr} \qquad \Rightarrow \qquad V = -\int_{\infty}^{r} \frac{KQ}{r^2}(-dr)\cos 180^o = \frac{KQ}{r}$$

**Example 53.** Four point charges are placed at the corners of a square of side $\ell$. Calculate potential at the centre of square.




**Solution :** V = 0 at 'C'. [Use $V = \frac{Kq}{r}$]

**Example 54.** Two point charges 2μC and – 4μC are situated at points (–2m, 0m) and (2 m, 0 m) respectively. Find out potential at point C(4 m, 0 m) and D(0 m, $\sqrt{5}$ m).

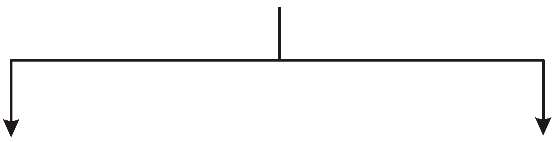



**Solution :** Potential at point C

$$V_C = V_{q_1} + V_{q_2} = \frac{K(2\mu C)}{6} + \frac{K(-4\mu C)}{2} = \frac{9 \times 10^9 \times 2 \times 10^{-6}}{6} - \frac{9 \times 10^9 \times 4 \times 10^{-6}}{2} = -15000 \text{ V.}$$

$$\text{Similarly, } V_D = V_{q_1} + V_{q_2} = \frac{K(2\mu C)}{\sqrt{(\sqrt{5})^2 + 2^2}} + \frac{K(-4\mu C)}{\sqrt{(\sqrt{5})^2 + 2^2}} = \frac{K(2\mu C)}{3} + \frac{K(-4\mu C)}{3} = -6000 \text{ V.}$$

**Finding potential due to continuous charges**



**If formula of E is tough, then we take a small element and integrate**

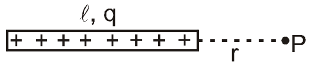
**V =** $\int dv$

**If formula of E is easy then, we use**

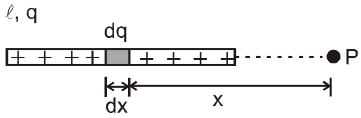
**V =** $-\int_{r\to\infty}^{r=r} \vec{E}.d\vec{r}$

**(i.e. for sphere, plate, infinite wire etc.)**

**Example 55.** A rod of length $\ell$ is uniformly charged with charge q. Calculate potential at point P.




**Solution :** Take a small element of length dx, at a distance x from left end. Potential due to this small element

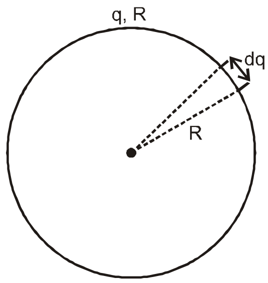



$dV = \dfrac{K(dq)}{x}$ $\therefore$ Total potential $\Rightarrow$ $V = \displaystyle\int_{x=0}^{x=\ell} \frac{kdq}{x}$

Where $dq = \dfrac{q}{\ell}dx$ $\Rightarrow$ $V = \displaystyle\int_{x=r}^{x=r+\ell} \frac{K\left(\frac{q}{\ell}dx\right)}{x} = \frac{Kq}{\ell}\log_e\left(\frac{\ell+r}{r}\right)$

## 7.5 Potential due to a ring :

**(i) Potential at the centre of uniformly charged ring :**



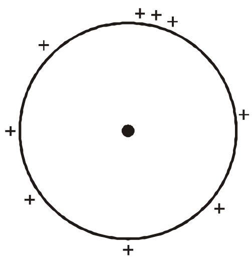

Potential due to the small element dq

$dV = \dfrac{Kdq}{R}$

$\therefore$ Net potential : $V = \displaystyle\int \frac{Kdq}{R}$

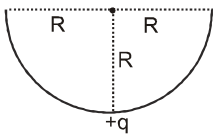
$\therefore$ $V = \dfrac{K}{R}\displaystyle\int dq = \frac{Kq}{R}$

**(ii) For non-uniformly charged ring potential at the center is**



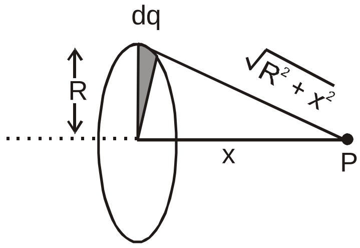
$V = \dfrac{Kq_{total}}{R}$

**(iii) Potential due to half ring at center is :**




$V = \dfrac{Kq}{R}$

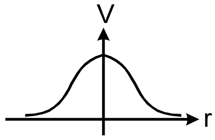
**(iv) Potential at the axis of a ring :** Calculation of potential at a point on the axis which is a distance x from centre of uniformly charged (total charge Q) ring of radius R.



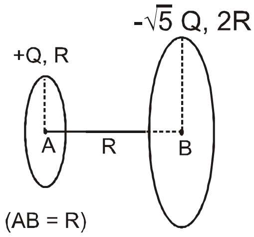

Consider an element of charge dq on the ring. Potential at point P due to charge dq will be

$$dV = \frac{K(dq)}{\sqrt{R^2 + x^2}}$$

$\therefore$ Net potential at point P due to all such element will be : $V = \int dv = \frac{KQ}{\sqrt{R^2 + x^2}}$




**Example 56.** Figure shows two rings having charges Q and $-\sqrt{5}$ Q. Find Potential difference between A and B i.e. $(V_A - V_B)$.




**Solution :** $V_A = \frac{KQ}{R} + \frac{K\left(-\sqrt{5}\ Q\right)}{\sqrt{(2R)^2 + (R)^2}}$ ; $V_B = \frac{K\left(-\sqrt{5}\ Q\right)}{2\ R} + \frac{K(\ Q)}{\sqrt{(R)^2 + (R)^2}}$

From above, we can easily find $V_A - V_B$.

**Example 57.** A point charge $q_0$ is placed at the centre of uniformly charged ring of total charge Q and radius R. If the point charge is slightly displaced with negligible force along axis of the ring then find out its speed when it reaches a large distance.

**Solution :** Only electric force is acting on $q_0$

$\therefore$ $W_{el} = \Delta K = \frac{1}{2} mv^2 - 0$ $\Rightarrow$ Now $W_{el})_{c\to\infty} = q_0 V_c = q_0 \frac{KQ}{R}$.

$\therefore$ $\frac{Kq_0Q}{R} = \frac{1}{2} mv^2$ $\Rightarrow$ $v = \sqrt{\frac{2Kq_0Q}{mR}}$

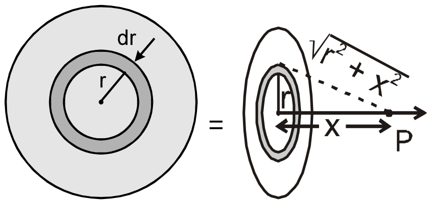
## 7.6 Potential due to uniformly charged disc :

$V = \frac{\sigma}{2\varepsilon_0}\left(\sqrt{R^2 + x^2} - x\right)$, where $\sigma$ is the charge density and x is the distance of the point on the axis from the center of the disc & R is the radius of disc.

### Finding potential due to a uniformly charged disc:

A disc of radius 'R' has surface charge density (charge/area) = $\sigma$. We have to find potential at its axis, at point 'P' which is at a distance x from the centre.

For this, we can divide the disc into thin rings and let's consider a thin ring of radius r and thickness dr. Suppose charge on the small ring element = dq. Potential due to this ring at point 'P' is:




$$dV = \frac{Kdq}{\sqrt{r^2 + x^2}}$$

So, net potential : $V_{net} = \int \frac{Kdq}{\sqrt{r^2+x^2}}$

Here, $\sigma$ = charge/area = $\frac{dq}{d(area)}$

So, dq = $\sigma$ × d (area) = $\sigma\,(2\pi r\,dr)$

(Here, d (area) = area of the small ring element = (length of ring) × (width of the ring) = ($2\pi$r) . (dr)

So, $V_{net} = \int_{r=0}^{r=R} \frac{K\sigma(2\pi r\,dr)}{\sqrt{r^2 + x^2}}$

To integrate it, let $r^2 + x^2 = y^2$

$2r\,dr = 2y\,dy$. Substituting we will get :

$$V_{net} = \int_{r=0}^{r=R} \frac{1}{4\pi\varepsilon_0}\frac{\sigma(2\pi)y\,dy}{y} \qquad \Rightarrow \qquad V_{net} = \frac{\sigma}{2\varepsilon_0}[y]_{r=0}^{r=R}$$

$$V_{net} = \frac{\sigma}{2\varepsilon_0}\left(\sqrt{r^2+x^2}\right)_{r=0}^{r=R} \qquad \Rightarrow \qquad V_{net} = \frac{\sigma}{2\varepsilon_0}\left(\sqrt{R^2+x^2} - x\right)$$

If a test charge $q_0$ is placed at point P, then potential energy of this charge $q_0$ due to the disc = U = $q_0V$

$$\Rightarrow \quad U = q_0\left[\frac{\sigma}{2\varepsilon_0}\left(\sqrt{R^2+x^2} - x\right)\right]$$

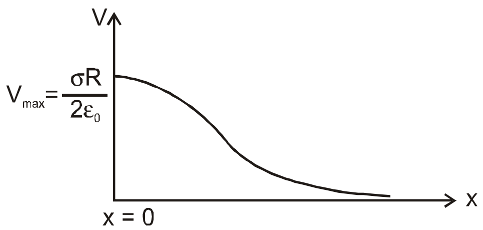
**Graph of V v/s x :** $V = \frac{\sigma}{2\varepsilon_0}\left(\sqrt{R^2+x^2} - x\right)$

at x = 0, $V = \frac{\sigma R}{2\varepsilon_0}$ to check whether V will increase with x or decrease, lets multiply and divide by conjugate.

$$V = \frac{\sigma}{2\varepsilon_0}\left(\sqrt{R^2+x^2} - x\right) \times \frac{\left(\sqrt{R^2+x^2}+x\right)}{\left(\sqrt{R^2+x^2}+x\right)}$$

$$\Rightarrow \quad V = \frac{\sigma R^2}{2\varepsilon_0}\left(\frac{1}{\left(\sqrt{R^2+x^2}+x\right)}\right)$$

Now, we can say that as x $\uparrow \Rightarrow$ V$\downarrow$ so curve will be



## 7.7 Potential Due To Uniformly Charged Spherical shell :

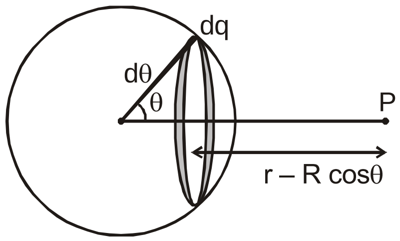
**Derivation of expression for potential due to uniformly charged hollow sphere of radius R and total charge Q, at a point which is at a distance r from centre for the following situation**

(i) $r > R$ (ii) $r < R$

Assume a ring of width $Rd\theta$ at angle $\theta$ from X axis (as shown in figure). Potential due to the ring at the point P will be

$$dV = \frac{K(dq)}{\sqrt{(r - R\cos\theta)^2 + (R\sin\theta)^2}}$$

Where $dq = 2\pi R \sin\theta\,(Rd\theta)\sigma$ where $Q = 4\pi R^2\sigma$




then, net potential $V = \int dV = \frac{KQ}{2}\int_0^{\pi} \frac{\sin\theta . d\theta}{\sqrt{(r - R\cos\theta)^2 + (R\sin\theta)^2}}$

Solving this eq. we find $V = \frac{KQ}{r}$ (for $r > R$)

& $V = \frac{KQ}{R}$ for ($r < R$)

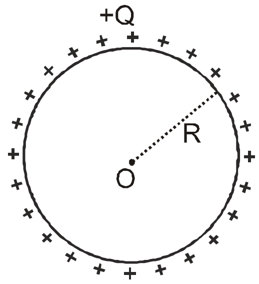
**Alternate Method :** As the formula of E is easy, we use $V = -\int_{r \to \infty}^{r = r} \vec{E}\cdot d\vec{r}$

**(i) At outside point ($r \geq R$) :** $V_{out} = -\int_{r \to \infty}^{r = r} \left(\frac{KQ}{r^2}\right) dr \Rightarrow V_{out} = \frac{KQ}{r} = \frac{KQ}{(\text{Distance from centre})}$

**For outside point, the hollow sphere acts like a point charge.**

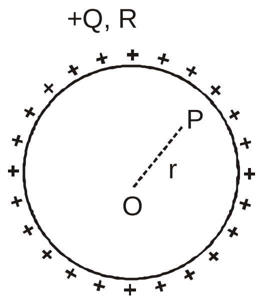
**(ii) Potential at the centre of the sphere (r = 0) :** As all the charges are at a distance R from the centre,

So, $V_{centre} = \frac{KQ}{R} = \frac{KQ}{(\text{Radius of the sphere})}$




**(iii) Potential at inside point (r < R) :** Suppose we want to find potential at point P, inside the sphere.

$\therefore$ Potential difference between Point P and O :




$$V_P - V_O = -\int_O^P \vec{E}_{in}\cdot d\vec{r} \quad \text{Where, } E_{in} = 0$$

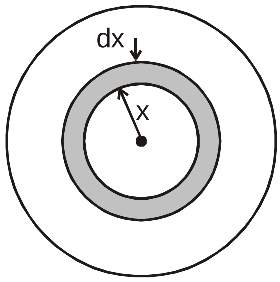
So $V_P - V_O = 0 \Rightarrow V_P = V_O = \frac{KQ}{R} \Rightarrow V_{in} = \frac{KQ}{R} = \frac{KQ}{(\text{Radius of the sphere})}$

## 7.8 Potential Due To Uniformly Charged Solid Sphere :

Derivation of expression for potential due to uniformly charged solid sphere of radius R and total charge Q (distributed in volume), at a point which is at a distance r from centre for the following situations.

(i) $r \geq R$ (ii) $r \leq R$

Consider an elementary shell of radius x and width dx




(i) For $r \geq R$ : $V = \int_0^R \frac{K \cdot 4\pi x^2 dx\rho}{r} = \frac{KQ}{r}$

(ii) For $r \leq R$ : $V = \int_0^r \frac{K \cdot 4\pi x^2 dx\rho}{r} + \int_r^R \frac{K4\pi x^2 dx\rho}{x}$

$= \frac{KQ}{2R^3}(3R^2 - r^2) \Rightarrow \left(\rho = \frac{Q}{\frac{4}{3}\pi R^3}\right)$

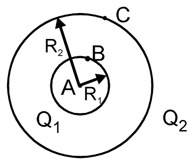
**From definition of potential**

(i) For $r \geq R$ : $V = -\int_\infty^r \frac{KQ}{r^2}\hat{r} \cdot dr = \frac{KQ}{r}$

(ii) For $r \leq R$ : $V = \int_\infty^R \frac{KQ}{r^2}.dr - \int_R^r \frac{KQr}{R^3} dr$

$V = \frac{KQ}{R} - \frac{KQ}{2R^3}[r^2 - R^2] = \frac{KQ}{2R^3}[2R^2 - r^2 + R^2] = \frac{KQ}{2R^3}(3R^2 - r^2)$

**Example 58.** Two concentric spherical shells of radius $R_1$ and $R_2$ ($R_2 > R_1$) are having uniformly distributed charges $Q_1$ and $Q_2$ respectively. Find out potential




(i) at point A
(ii) at surface of smaller shell (i.e. at point B)
(iii) at surface of larger shell (i.e. at point C)
(iv) at $r \leq R_1$
(v) at $R_1 \leq r \leq R_2$
(vi) at $r \geq R_2$

**Solution :** Using the results of hollow sphere as given in the table 7.4.

(i) $V_A = \frac{KQ_1}{R_1} + \frac{KQ_2}{R_2}$

(ii) $V_B = \frac{KQ_1}{R_1} + \frac{KQ_2}{R_2}$

(iii) $V_C = \frac{KQ_1}{R_2} + \frac{KQ_2}{R_2}$

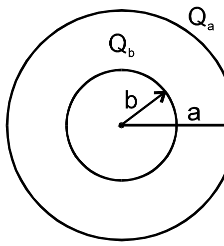
(iv) for $r \leq R_1$, $V = \frac{KQ_1}{R_1} + \frac{KQ_2}{R_2}$

(v) for $R_1 \leq r \leq R_2$, $V = \frac{KQ_1}{r} + \frac{KQ_2}{R_2}$

(vi) for $r \geq R_2$ , $V = \frac{KQ_1}{r} + \frac{KQ_2}{r}$

**Example 59.** Two hollow concentric non-conducting spheres of radius the a and b (a > b) contain charges $Q_a$ and $Q_b$ respectively. Prove that potential difference between the two spheres is independent of charge on outer sphere. If outer sphere is given an extra charge, is there any change in potential difference?

**Solution :** $V_{\text{inner sphere}} = \dfrac{KQ_b}{b} + \dfrac{KQ_a}{a}$




$V_{\text{outer sphere}} = \dfrac{KQ_b}{a} + \dfrac{KQ_a}{a}$

$V_{\text{inner sphere}} - V_{\text{outer sphere}} = \dfrac{KQ_b}{b} - \dfrac{KQ_b}{a}$

$\therefore \quad \Delta V = KQ_b \left[\dfrac{1}{b} - \dfrac{1}{a}\right]$

Which is independent of charge on outer sphere. If outer sphere in given any extra charge, then there will be no change in potential difference.

## 8. POTENTIAL DIFFERENCE

The potential difference between two points A and B is work done by external agent against electric field in taking a unit positive charge from A to B with no change in kinetic energy between initial and final points ie. $\Delta K = 0$ or $K_i = K_f$

### (a) Mathematical representation :

If $(W_{A\to B})_{ext}$ = Work done by external agent against electric field in taking the unit charge from A to B

Then, $V_B - V_A = \left.\dfrac{(W_{A\to B})_{ext}}{q}\right)_{\Delta K=0} = \dfrac{-(W_{A\to B})_{electric}}{q} = \dfrac{U_B - U_A}{q} = \dfrac{-\int_A^B \vec{F}_e \cdot d\vec{r}}{q} = -\int_A^B \vec{E} \cdot d\vec{r}$

**Note : Take W and q both with sign**

### (b) Properties :

**(i)** The difference of potential between two points is called potential difference. It is also called voltage.

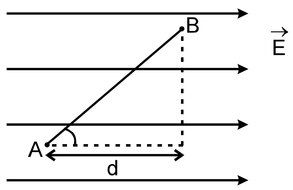
**(ii)** Potential difference is a scalar quantity. Its S.I. unit is also volt.

**(iii)** If $V_A$ and $V_B$ be the potential of two points A and B, then work done by an external agent in taking the charge q from A to B is $(W_{ext})_{AB} = q\,(V_B - V_A)$ or $(W_{el})_{AB} = q\,(V_A - V_B)$.

**(iv)** Potential difference between two points is independent of reference point.

## 8.1 Potential difference in a uniform electric field :

$\mathbf{V_B - V_A = -}\ \vec{E} \cdot \overrightarrow{AB}$




$\Rightarrow \quad V_B - V_A = -|E|\,|AB| \cos\theta$

$= -|E|\,d$

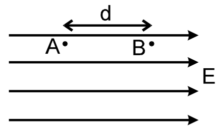
$= -Ed$

where d = effective distance between A and B along electric field.

or we can also say that $E = \dfrac{\Delta V}{\Delta d}$

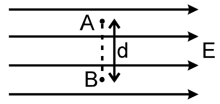
**Special Cases :**

**Case-1.** Line AB is parallel to electric field.




$\therefore \quad V_A - V_B = Ed$

**Case-2.** Line AB is perpendicular to electric field.




$\therefore \quad V_A - V_B = 0 \quad \Rightarrow \quad V_A = V_B$

**Note : In the direction of electric field potential always decreases.**

**Example 60.** 1μC charge is shifted from A to B and it is found that work done by an external force is 40μJ in doing so against electrostatic forces, then find potential difference $V_A - V_B$

**Solution :** $(W_{AB})_{ext} = q(V_B - V_A) \quad \Rightarrow \quad 40\ \mu J = 1\mu C\ (V_B - V_A) \quad \Rightarrow \quad V_A - V_B = -40\ V$

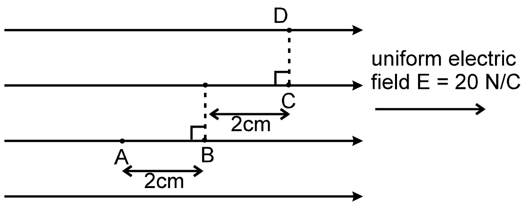
**Example 61.** A uniform electric field is present in the positive x-direction. If the intensity of the field is 5N/C then find the potential difference ($V_B - V_A$) between two points A (0m, 2 m) and B (5 m ,3 m)

**Solution :** $V_B - V_A = -\vec{E}\cdot\overrightarrow{AB} = -(5\hat{i}).(5\hat{i}+\hat{j}) = -25V.$

The electric field intensity in uniform electric field, $E = \dfrac{\Delta V}{\Delta d}$

Where $\Delta V$ = potential difference between two points.
$\Delta d$ = effective distance between the two points.
(projection of the displacement along the direction of electric field.)

**Example 62.** Find out following




(i) $V_A - V_B$ (ii) $V_B - V_C$ (iii) $V_C - V_A$ (iv) $V_D - V_C$
(v) $V_A - V_D$ (vi) Arrange the order of potential for points A, B, C and D.

**Solution :** (i) $|\Delta V_{AB}| = Ed = 20 \times 2 \times 10^{-2} = 0.4$ so, $V_A - V_B = 0.4\ V$

because **In the direction of electric field potential always decreases.**

(ii) $|\Delta V_{BC}| = Ed = 20 \times 2 \times 10^{-2} = 0.4$ so, $V_B - V_C = 0.4\ V$

(iii) $|\Delta V_{CA}| = Ed = 20 \times 4 \times 10^{-2} = 0.8$ so, $V_C - V_A = -0.8\ V$

because **In the direction of electric field potential always decreases.**

(iv) $|\Delta V_{DC}| = Ed = 20 \times 0 = 0$ so, $V_D - V_C = 0$

because the effective distance between D and C is zero.

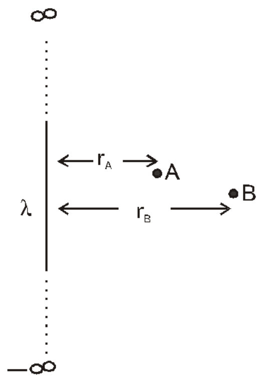
(v) $|\Delta V_{AD}| = Ed = 20 \times 4 \times 10^{-2} = 0.8$ so, $V_A - V_D = 0.8\ V$

because **In the direction of electric field potential always decreases.**

(vi) The order of potential is : $V_A > V_B > V_C = V_D$.

## 8.2 Potential difference due to infinitely long wire :

Derivation of expression for potential difference between two points, which have perpendicular distance $r_A$ and $r_B$ from infinitely long line charge of uniform linear charge density $\lambda$. :

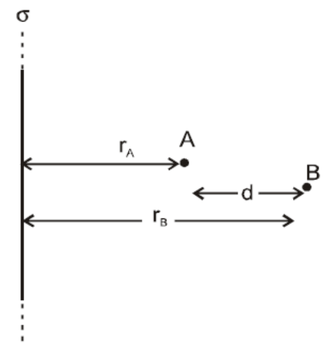



From definition of potential difference : $V_{AB} = V_B - V_A = -\int_{r_A}^{r_B} \vec{E} \cdot \overrightarrow{dr} = -\int_{r_A}^{r_B} \frac{2K\lambda}{r} \hat{r} \cdot \overrightarrow{dr}$

$\therefore \quad V_{AB} = -2K\lambda \, \ell n\left(\frac{r_B}{r_A}\right)$

## 8.3 Potential difference due to infinitely long thin sheet:

Derivation of expression for potential difference between two points, having separation d in the direction perpendicularly to a very large uniformly charged thin sheet of uniform surface charge density $\sigma$.




Let the points A and B have perpendicular distance $r_A$ and $r_B$ respectively then from definition of potential difference.

$V_{AB} = V_B - V_A = -\int_{r_A}^{r_B} \vec{E} \cdot \overrightarrow{dr} = -\int_{r_A}^{r_B} \frac{\sigma}{2\varepsilon_0} \hat{r} \cdot \overrightarrow{dr} \Rightarrow V_{AB} = -\frac{\sigma}{2\varepsilon_o}(r_B - r_A) = -\frac{\sigma d}{2\varepsilon_o}$

# 9. EQUIPOTENTIAL SURFACE :

**9.1 Definition :** If potential of a surface (imaginary or physically existing) is same throughout, then such surface is known as an equipotential surface.

## 9.2 Properties of equipotential surfaces :

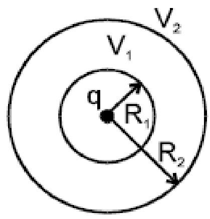
**(i)** When a charge is shifted from one point to another point on an equipotential surface, then work done against electrostatic forces is zero.

**(ii)** Electric field is always perpendicular to equipotential surfaces.

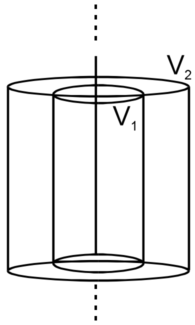
**(iii)** Two equipotential surfaces do not cross each other.

## 9.3 Examples of equipotential surfaces :

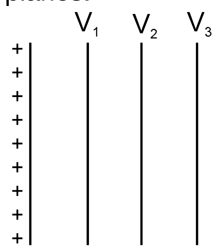
**(i)** **Point charge :** Equipotential surfaces are concentric and spherical as shown in figure. In figure, we can see that sphere of radius $R_1$ has potential $V_1$ throughout its surface and similarly for other concentric sphere potential is same.




**(ii) Line charge :** Equipotential surfaces have curved surfaces as that of coaxial cylinders of different radii.

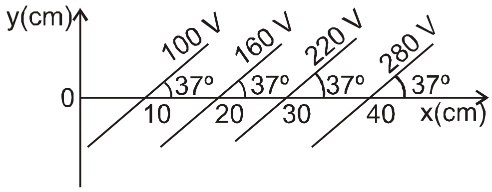



**(iii) Uniformly charged large conducting / non conducting sheets :**
Equipotential surfaces are parallel planes.




**Note :** In uniform electric field equi-potential surfaces are always parallel planes.

**Example 63.** Some equipotential surfaces are shown in figure. What can you say about the magnitude and the direction of the electric field ?

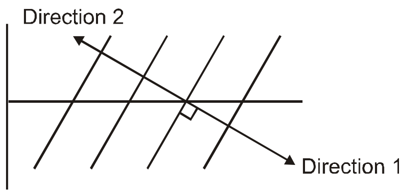



**Solution :** Here, we can say that the electric will be perpendicular to equipotential surfaces.

Also, $|\vec{E}| = \dfrac{\Delta V}{\Delta d}$

Where, $\Delta V$ = potential difference between two equipotential surfaces.

$\Delta d$ = perpendicular distance between two equipotential surfaces.

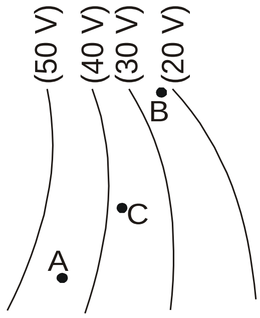



So $|\vec{E}| = \dfrac{60}{(10\sin 37^\circ)\times 10^{-2}}$ = 1000 V/m

Now there are two perpendicular directions: either direction 1 or direction 2 as shown in figure, but since we know that in the direction of electric field, electric potential decreases, so the correct direction is direction 2.

Hence E = 1000 V/m, making an angle 127° with the x-axis

**Example 64.** Figure shows some equipotential surfaces produce by some charges. At which point, the value of electric field is greatest?

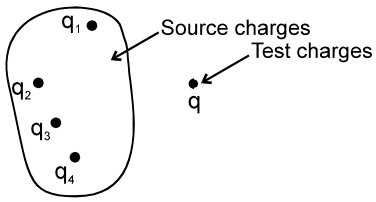



**Solution :** E is larger where equipotential surfaces are closer. ELOF are $\perp$ to equipotential surfaces. In the figure, we can see that for point B, they are closer so E at point B is maximum

# 10. ELECTROSTATIC POTENTIAL ENERGY

## 10.1 Electrostatic potential energy of a point charge due to many charges :

The electrostatic potential energy of a point charge at a point in electric field is the work done in taking the charge from reference point (generally at infinity) to that point without change in kinetic energy.




Its Mathematical formula is $U = W_{\infty \to P)ext}\big|_{\Delta K = 0} = qV = -W_{P \to \infty)ele}$

Here, q is the charge whose potential energy is being calculated and V is the potential at its position due to the source charges.

**Note : Always put q and V with sign.**

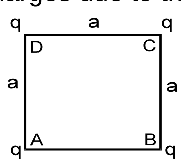
## 10.2 Properties :

**(i)** Electric potential energy is a scalar quantity but may be positive, negative or zero.

**(ii)** Its unit is same as unit of work or energy i.e., joule (in S.I. system). Some times energy is also given in electron-volts. Where, $1eV = 1.6 \times 10^{-19}$ J

**(iii)** Electric potential energy depends on reference point. (Generally Potential Energy at $r = \infty$ is taken zero)

**Example 65.** The four identical charges (q each) are placed at the corners of a square of side a. Find the potential energy of one of the charges due to the remaining charges.




**Solution :** The electric potential of point A due to the charges placed at B, C and D is

$$V = \frac{1}{4\pi\varepsilon_0}\frac{q}{a} + \frac{1}{4\pi\varepsilon_0}\frac{q}{\sqrt{2}a} + \frac{1}{4\pi\varepsilon_0}\frac{q}{a} = \frac{1}{4\pi\varepsilon_0}\left(2 + \frac{1}{\sqrt{2}}\right)\frac{q}{a}$$

$\therefore$ Potential energy of the charge at A is = qV = $\frac{1}{4\pi\varepsilon_0}\left(2+\frac{1}{\sqrt{2}}\right)\frac{q^2}{a}$.

**Example 66.** A particle of mass 40 mg and carrying a charge 5 × $10^{-9}$ C is moving directly towards a fixed positive point charge of magnitude $10^{-8}$ C. When it is at a distance of 10 cm from the fixed point charge it has speed of 50 cm/s. At what distance from the fixed point charge will the particle come momentarily to rest? Is the acceleration constant during the motion?

**Solution :** If the particle comes to rest momentarily at a distance r from the fixed charge, then from conservation of energy, we have ?

$$\frac{1}{2}mu^2 + \frac{1}{4\pi\varepsilon_0}\frac{Qq}{a} = \frac{1}{4\pi\varepsilon_0}\frac{Qq}{r}$$

Substituting the given data, we get: $\frac{1}{2}$ × 40 × $10^{-6}$ × $\frac{1}{2}$ × $\frac{1}{2}$ = 9 × $10^9$ × 5 × $10^{-8}$ ×$10^{-9}$ $\left[\frac{1}{r}-10\right]$

or, $\frac{1}{r}$–10 = $\frac{5\times10^{-6}}{9\times5\times10^{-8}} = \frac{100}{9}$ $\Rightarrow$ $\frac{1}{r}$ = $\frac{190}{9}$ $\Rightarrow$ r = $\frac{9}{190}$ m

or, i.e., r = 4.7 × $10^{-2}$ m. As here, F = $\frac{1}{4\pi\varepsilon_0}\frac{Qq}{r^2}$ So, acc. = $\frac{F}{m} \propto \frac{1}{r^2}$

i.e., Acceleration is not constant during the motion.

**Example 67.** A proton moves from a large distance with a speed u m/s directly towards a free proton originally at rest. Find the distance of closest approach for the two protons in terms of mass of proton m and its charge e.

**Solution :** As here the particle at rest is free to move, when one particle approaches the other, due to electrostatic repulsion other will also start moving and so the velocity of first particle will decrease while of other will increase and at closest approach, both will move with same velocity. So, if v is the common velocity of each particle at closest approach, then by 'conservation of momentum' of the two proton system.

mu = mv + mv i.e., v = $\frac{1}{2}$u

And by conservation of energy, $\frac{1}{2}$ mu$^2$ = $\frac{1}{2}$ mv$^2$ + $\frac{1}{2}$mv$^2$ + $\frac{1}{4\pi\varepsilon_0}\frac{e^2}{r}$

$\Rightarrow$ $\frac{1}{2}$mu$^2$ – m$\left(\frac{u}{2}\right)^2$ = $\frac{1}{4\pi\varepsilon_0}\frac{e^2}{r}$ [as v = $\frac{u}{2}$]

$\Rightarrow$ $\frac{1}{4}$mu$^2$ = $\frac{e^2}{4\pi\varepsilon_0 r}$ $\Rightarrow$ r = $\frac{e^2}{\pi m\varepsilon_0 u^2}$

## 11. ELECTROSTATIC POTENTIAL ENERGY OF A SYSTEM OF CHARGES

**(This concept is useful when more than one charges move.)**

It is the work done by an external agent against the internal electric field required to make a system of charges in a particular configuration from infinite separation without change in their kinetic energies.

### 11.1 Types of system of charges :

(i) Point charge system (ii) Continuous charge system.

### 11.2 Derivation for a system of point charges:

(i) Keep all the charges at infinity. Now bring the charges one by one to its corresponding position and find work required. PE of the system is algebraic sum of all the works.

Let $W_1$ = Work done in bringing first charge.

$W_2$ = Work done in bringing second charge against force due to 1$^{st}$ charge.

$W_3$ = Work done in bringing third charge against force due to 1st and 2nd charge.

PE = $W_1 + W_2 + W_3$ + ...... . (This will contain $\frac{n(n-1)}{2}$ = ${}^nC_2$ terms)

(ii) Method of calculation (to be used in problems) :

U = sum of the interaction energies of the charges.

= $(U_{12} + U_{13} + ........ + U_{1n}) + (U_{23} + U_{24} + ........ + U_{2n}) + (U_{34} + U_{35} + ........ + U_{3n})$ .....

(iii) Method of calculation useful for symmetrical point charge systems.

Find PE of each charge due to rest of the charges.

If $U_1$ = PE of first charge due to all other charges.

= $(U_{12} + U_{13} + ........ + U_{1n})$

$U_2$ = PE of second charge due to all other charges.

= $(U_{21} + U_{23} + ........ + U_{2n})$ then U = PE of the system = $\frac{U_1 + U_2 + ....U_n}{2}$

**Example 68.** Find out potential energy of the two point charge system having charges $q_1$ and $q_2$ separated by distance r.

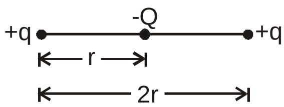
**Solution :** Let both the charges be placed at a very large separation initially.

Let, $W_1$ = work done in bringing charge $q_1$ in absence of $q_2 = q_1(V_f - V_i) = 0$

$W_2$ = work done in bringing charge $q_2$ in presence of $q_1 = q_2(V_f - V_i) = q_2(Kq_1/r - 0)$

$\therefore$ PE = $W_1 + W_2 = 0 + Kq_1q_2/r = Kq_1q_2/r$

**Example 69.** Figure shows an arrangement of three point charges. The total potential energy of this arrangement is zero. Calculate the ratio $\frac{q}{Q}$.

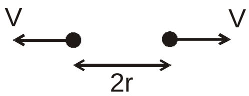



**Solution :** $U_{sys} = \frac{1}{4\pi\varepsilon_0}\left[\frac{-qQ}{r} + \frac{(+q)(+q)}{2r} + \frac{Q(-q)}{r}\right] = 0$

$-Q + \frac{q}{2} - Q = 0$ or $2Q = \frac{q}{2}$ or $\frac{q}{Q} = \frac{4}{1}$.

**Example 70.** Two point charges, each of mass m and charge q are released when they are at a distance r from each other. What is the speed of each charged particle when they are at a distance 2r ?

**Solution :** According to momentum conservation, both the charge particles will move with same speed. Now applying energy conservation:




$0 + 0 + \frac{Kq^2}{r} = 2\frac{1}{2}mv^2 + \frac{Kq^2}{2r} \Rightarrow v = \sqrt{\frac{Kq^2}{2rm}}$

**Example 71.** Two charged particles each having equal charges $2 \times 10^{-5}$ C are brought from infinity to within a separation of 10 cm. Calculate the increase in potential energy during the process and the work required for this purpose.

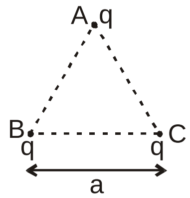
**Solution :** $\Delta U = U_f - U_i = U_f - 0 = U_f$

We have to simply calculate the electrostatic potential energy of the given system of charges

$\therefore \Delta U = U_f = \frac{1}{4\pi\varepsilon_0}\frac{q_1q_2}{r} = \frac{9\times10^9\times2\times10^{-5}\times2\times10^{-5}\times100}{10}$ J = 36 J

$\therefore$ Work required = 36 J = equal to change in potential energy of system

**Example 72.** Three equal charges q each are placed at the corners of an equilateral triangle of side a.

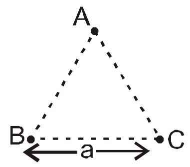



(i) Find out potential energy of charge system.

(ii) Calculate work required to decrease the side of triangle to a/2.

(iii) If the charges are released from the shown position and each of them has same mass m then find the speed of each particle when they lie on triangle of side 2a.

**Solution :** **(i) Method I** (Derivation)




Assume all the charges are at infinity initially. Work done in putting charge q at corner A

$\Rightarrow$ $W_1 = q(v_f - v_i) = q(0 - 0)$

Since potential at A is zero in absence of charges, work done in putting q at corner B in presence of charge at A :

$$\Rightarrow W_2 = \left(\frac{Kq}{a} - 0\right) q = \frac{Kq^2}{a}$$

Similarly work done in putting charge q at corner C in presence of charge at A and B.

$$\Rightarrow W_3 = q(v_f - v_i) = q\left[\left(\frac{Kq}{a} + \frac{Kq}{a}\right) - 0\right] = \frac{2Kq^2}{a}$$

So, net potential energy PE = $W_1 + W_2 + W_3 = 0 + \frac{Kq^2}{a} + \frac{2Kq^2}{a} = \frac{3Kq^2}{a}$

**Method II** (using direct formula) :

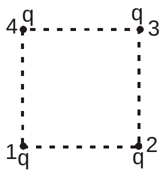
$$U = U_{12} + U_{13} + U_{23} = \frac{Kq^2}{a} + \frac{Kq^2}{a} + \frac{Kq^2}{a} = \frac{3Kq^2}{a}$$

**(ii)** Work required to decrease the sides $W = U_f - U_i = \frac{3Kq^2}{a/2} - \frac{3Kq^2}{a} = \frac{3Kq^2}{a}$ Joules

**(iii)** Work done by electrostatic forces = Change is kinetic energy of particles.

$$U_i - U_f = K_f - K_i \qquad \Rightarrow \frac{3Kq^2}{a} - \frac{3Kq^2}{2a} = 3\left(\frac{1}{2}mv^2\right) - 0 \Rightarrow v = \sqrt{\frac{Kq^2}{am}}$$

**Example 73.** Four identical point charges q each are placed at four corners of a square of side a. Find out potential energy of the charge system




**Solution :** **Method 1 (using direct formula) :** $U = U_{12} + U_{13} + U_{14} + U_{23} + U_{24} + U_{34}$

$$= \frac{Kq^2}{a} + \frac{Kq^2}{a\sqrt{2}} + \frac{Kq^2}{a} + \frac{Kq^2}{a} + \frac{Kq^2}{a\sqrt{2}} + \frac{Kq^2}{a} = \left[\frac{4Kq^2}{a} + \frac{2Kq^2}{a\sqrt{2}}\right] = \frac{2Kq^2}{a}\left[2 + \frac{1}{\sqrt{2}}\right]$$

**Method 2** [Using, $U = \frac{1}{2}(U_1 + U_2 + ......)$] **:**

$U_1$ = total P.E. of charge at corner 1 due to all other charges.

$U_2$ = total P.E. of charge at corner 2 due to all other charges.

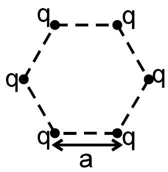
$U_3$ = total P.E. of charge at corner 3 due to all other charges.

$U_4$ = total P.E. of charge at corner 4 due to all other charges.

Since, due to symmetry, $U_1 = U_2 = U_3 = U_4$

$$U_{Net} = \frac{U_1 + U_2 + U_3 + U_4}{2} = 2U_1 = 2\left[\frac{Kq^2}{a} + \frac{Kq^2}{a} + \frac{Kq^2}{\sqrt{2}a}\right] = \frac{2Kq^2}{a}\left[2 + \frac{1}{\sqrt{2}}\right]$$

**Example 74.** Six equal point charges q each are placed at six corners of a hexagon of side a. Find out potential energy of charge system




**Solution :** $U_{Net} = \frac{U_1 + U_2 + U_3 + U_4 + U_5 + U_6}{2}$

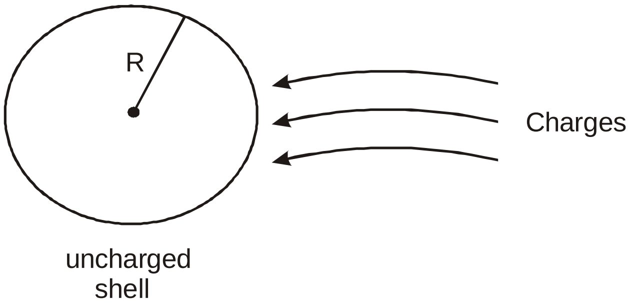
Due to symmetry $U_1 = U_2 = U_3 = U_4 = U_5 = U_6$ So $U_{net} = 3U_1 = \frac{3Kq^2}{a}\left[2 + \frac{2}{\sqrt{3}} + \frac{1}{2}\right]$

## 11.3 Derivation of electric potential energy for continuous charge system :

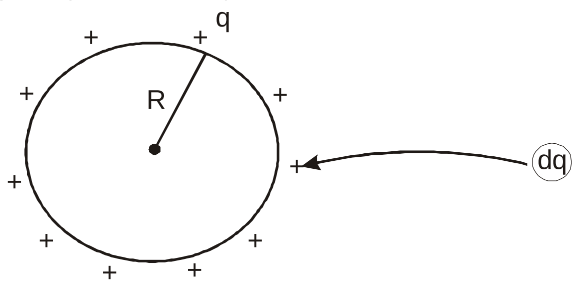
This energy is also known as self energy.

**(i) Finding P.E., (Self Energy) of a uniformly Charged spherical shell :**

For this, lets use method 1 : Take an uncharged shell. Now bring charges one by one from infinity to the surface fo the shell. The work required in this process will be stored as potential Energy.




Suppose, we have given charge q to the sphere and now we are giving extra charge dq to it. Work required to bring dq charge from infinity to the shell is




$dW = (dq)(V_f - V_i)$ $\Rightarrow$ $dW = (dq)\left(\frac{Kq}{R} - 0\right) = \frac{Kq}{R}dq$

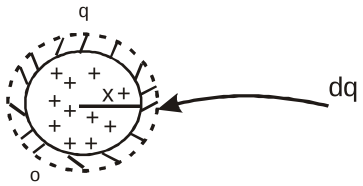
$\Rightarrow$ Total work required to give charge Q is $W = \int_{q=o}^{q=Q} \frac{Kq}{R}dq = \frac{KQ^2}{2R}$

This work will be stored in the form of P.E. (self energy)

So, P.E. of a charged spherical shell : $U = \frac{KQ^2}{2R}$

**(ii) Self energy of uniformly charged solid sphere :**

In this case we have to assemble a solid charged sphere. So we bring the charges one-by-one from infinity to the sphere so that the size of the sphere increases.




Suppose we have given charge q to the sphere, and its radius becomes 'x' .Now we are giving extra charge dq to it, which will increase its radius by 'dx'

$\therefore$ Work required to bring charge dq from infinity to the sphere

$$= dq\,(V_f - V_i) = (dq)\left(\frac{Kq}{x} - o\right) = \frac{Kqdq}{x}$$

$\therefore$ Total work required to give charge Q to the sphere

$W = \int \frac{Kqdq}{x}$, where $q = \rho\left(\frac{4}{3}\pi x^3\right)$

& $dq = \rho\,(4\,\pi x^2\,dx)$ $\Rightarrow$ $W = \int_{x=o}^{x=R} K\frac{\rho\left(\frac{4}{3}\pi x^3\right)\rho(4\pi x^2 dx)}{x}$

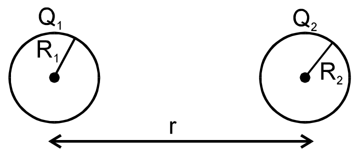
Solving, well get : $W = \frac{3}{5}\frac{KQ^2}{R} = U_{self}$ (for a solid sphere)

**Example 75.** A spherical shell of radius R with uniform charge q is expanded to a radius 2R. Find the work performed by the electric forces and external agent against electric forces in this process.

**Solution :** $W_{ext} = U_f - U_i = \frac{q^2}{16\pi\varepsilon_0 R} - \frac{q^2}{8\pi\varepsilon_0 R} = -\frac{q^2}{16\pi\varepsilon_0 R}$

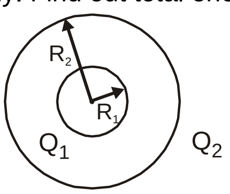
$W_{elec} = U_i - U_f = \frac{q^2}{8\pi\varepsilon_0 R} - \frac{q^2}{16\pi\varepsilon_0 R} = \frac{q^2}{16\pi\varepsilon_0 R}$

**Example 76.** Two non-conducting hollow uniformly charged spheres of radii $R_1$ and $R_2$ with charge $Q_1$ and $Q_2$ respectively are placed at a distance r. Find out total energy of the system.




**Solution :** $U_{total} = U_{self} + U_{Interaction} = \frac{Q_1^2}{8\pi\varepsilon_0 R_1} + \frac{Q_2^2}{8\pi\varepsilon_0 R_2} + \frac{Q_1 Q_2}{4\pi\varepsilon_0 r}$

**Example 77.** Two concentric spherical shells of radius $R_1$ and $R_2$ ($R_2 > R_1$) are having uniformly distributed charges $Q_1$ and $Q_2$ respectively. Find out total energy of the system.




**Solution :** $U_{total} = U_{self\ 1} + U_{self\ 2} + U_{Interaction} = \frac{Q_1^2}{8\pi\varepsilon_0 R_1} + \frac{Q_2^2}{8\pi\varepsilon_0 R_2} + \frac{Q_1 Q_2}{4\pi\varepsilon_0 R_2}$

## 11.4 Energy density :

**Def :** Energy density is defined as energy stored in unit volume in any electric field. Its mathematical formula is given as following :

Energy density = $\frac{1}{2}\varepsilon E^2$

where E = electric field intensity at that point

. $\varepsilon = \varepsilon_0\varepsilon_r$ electric permittivity of medium

**Example 78.** Find out energy stored in an imaginary cubical volume of side a in front of a infinitely large non-conducting sheet of uniform charge density $\sigma$.

**Solution :** Energy stored : $U = \int \frac{1}{2}\varepsilon_0 E^2 dV$ ; where dV is small volume

$\therefore \quad U = \frac{1}{2}\varepsilon_0 E^2 \int dV \qquad \because \quad$ E is constant .

$\therefore \quad U = \frac{1}{2}\varepsilon_0 \frac{\sigma^2}{4\varepsilon_0^2} . a^3 . = \frac{\sigma^2 a^3}{8\varepsilon_0}$

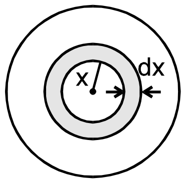
**Example 79.** Find out energy stored in the electric field of uniformly charged thin spherical shell of total charge Q and radius R.

**Solution :** We know that electric field inside the shell is zero, so the energy is stored only outside the shell, which can be calculated by using energy density formula.

$U_{self} = \int_{x=R}^{x\to\infty} \frac{1}{2}\varepsilon_0 E^2 dV$

Consider an elementary shell of thickness dx and radius x (x > R).

Volume of the shell = $(4\pi x^2 dx) = dV$




$$U = \int_R^\infty \frac{1}{2}\varepsilon_0 \left[\frac{KQ}{x^2}\right]^2 . 4\pi x^2\, dx = \frac{1}{2}\varepsilon_0\, K^2Q^2 4\pi \int_R^\infty \frac{1}{x^2} dx$$

$$= \frac{4\pi\varepsilon_0}{2} \frac{Q^2}{(4\pi\varepsilon_0)^2} . \left(\frac{1}{R}\right) = \frac{Q^2}{8\pi\varepsilon_0 R} = \frac{KQ^2}{2R} .$$

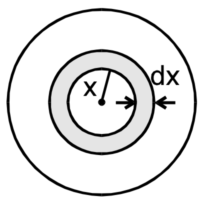
**Example 80.** Find out energy stored inside a solid non-conducting sphere of total charge Q and radius R. [Assume charge is uniformly distributed in its volume.]

**Solution :** We can consider solid sphere to be made of large number of concentric spherical shells. Also electric field intensity at the location of any particular shell is constant.

$\therefore \quad U_{inside} = \int_0^R \frac{1}{2}\varepsilon_0 E^2 dV$

Consider an elementary shell of thickness dx and radius x.

Volume of the shell = $(4\pi x^2 dx)$

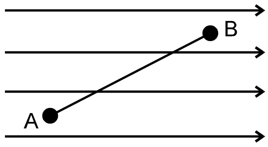



$$U_{inside} = \int_0^R \frac{1}{2}\varepsilon_0 \left[\frac{KQx}{R^3}\right]^2 .4\pi x^2\, dx = \frac{1}{2}\varepsilon_0 \frac{K^2Q^2 4\pi}{R^6} \int_0^R x^4 dx$$

$$= \frac{4\pi\varepsilon_0}{2R^6} \frac{Q^2}{(4\pi\varepsilon_0)^2} . \frac{R^5}{5} = \frac{Q^2}{40\pi\varepsilon_0 R} = \frac{KQ^2}{10R} .$$

# 12. RELATION BETWEEN ELECTRIC FIELD INTENSITY AND ELECTRIC POTENTIAL :

## 12.1 For uniform electric field :




(i) Potential difference between two points A and B

$V_B - V_A = -\vec{E}\cdot\overrightarrow{AB}$

## 12.2 Non uniform electric field

(i) $E_x = -\dfrac{\partial V}{\partial x}$ , $E_y = -\dfrac{\partial V}{\partial y}$ , $E_z = -\dfrac{\partial V}{\partial z}$

$$\Rightarrow \quad \vec{E} = E_x\hat{i} + E_y\hat{j} + E_z\hat{k} = -\left[\hat{i}\frac{\partial}{\partial x}V + \hat{j}\frac{\partial}{\partial y}V + \hat{k}\frac{\partial}{\partial z}V\right] = -\left[\hat{i}\frac{\partial}{\partial x} + \hat{j}\frac{\partial}{\partial y} + \hat{k}\frac{\partial}{\partial z}\right]V = -\nabla V = -\text{grad } V$$

Where, $\dfrac{\partial V}{\partial x}$ = derivative of V with respect to x (keeping y and z constant)

$\dfrac{\partial V}{\partial y}$ = derivative of V with respect to y (keeping z and x constant)

$\dfrac{\partial V}{\partial z}$ = derivative of V with respect to z (keeping x and y constant)

## 12.3 If electric potential and electric field depends only on one coordinate, say r :

(i) $\vec{E} = -\dfrac{\partial V}{\partial r}\hat{r}$ where, $\hat{r}$ is a unit vector along increasing r.

(ii) $\int dV = -\int \vec{E}\cdot\vec{dr} \quad \Rightarrow \quad V_B - V_A = -\int_{r_A}^{r_B} \vec{E}.\vec{dr}$

$\vec{dr}$ is along the increasing direction of r.

(iii) The potential of a point $V = -\int_{\infty}^{r} \vec{E}.\vec{dr}$

**Example 81.** A uniform electric field is along x-axis. The potential difference $V_A - V_B = 10$ V is between two points A (2m, 3m) and B (4m, 8m). Find the electric field intensity.

**Solution :** $E = \dfrac{\Delta V}{\Delta d} = \dfrac{10}{2} = 5$ V / m. (It is along + ve x-axis)

**Example 82.** $V = x^2 + y$. Find $\vec{E}$.

**Solution :** $\dfrac{\partial V}{\partial x}$ c = 2x, $\dfrac{\partial V}{\partial y} = 1$ and $\dfrac{\partial V}{\partial z} = 0$

$$\vec{E} = -\left(\hat{i}\frac{\partial V}{\partial x} + \hat{j}\frac{\partial V}{\partial y} + \hat{k}\frac{\partial V}{\partial z}\right) = -(2x\hat{i} + \hat{j})$$

$\therefore$ Electric field is non-uniform.

**Example 83.** For given $\vec{E} = 2x\hat{i} + 3y\hat{j}$, find the potential at (x, y) if V at origin is 5 volts.

**Solution :** $\int_5^v dV = -\int \vec{E}\cdot\overline{dr} = -\int_0^x E_x dx - \int_0^y E_y dy \Rightarrow V - 5 = -\dfrac{2x^2}{2} - \dfrac{3y^2}{2} \Rightarrow V = -\dfrac{2x^2}{2} - \dfrac{3y^2}{2} + 5.$

# 13. ELECTRIC DIPOLE

## 13.1 Electric Dipole

If two point charges, equal in magnitude 'q' and opposite in sign separated by a distance 'a' such that the distance of field point r>>a, the system is called a dipole. The electric dipole moment is defined as a vector quantity having magnitude p = (q × a) and direction from negative charge to positive charge.

**Note:** [In chemistry, the direction of dipole moment is assumed to be from positive to negative charge.] The C.G.S unit of electric dipole moment is **debye** which is defined as the dipole moment of two equal and opposite point charges each having charge $10^{-10}$ Franklin and separation of 1 Å, i.e.,

$$1 \text{ debye (D)} = 10^{-10} \times 10^{-8} = 10^{-18} \text{ Fr} \times \text{cm}$$

or $$1 \text{ D} = 10^{-18} \times \frac{C}{3\times10^{9}} \times 10^{-2} \text{ m} = 3.3 \times 10^{-30} \text{ C} \times \text{m}.$$

S.I. Unit is coulomb × metre = C . m

**Example 84.** A system has two charges $q_A = 2.5 \times 10^{-7}$ C and $q_B = -2.5 \times 10^{-7}$ C located at points A : (0, 0, – 0.15 m) and B ; (0, 0, + 0.15 m) respectively. What is the net charge and electric dipole moment of the system ?

**Solution :** Net charge = $2.5 \times 10^{-7} - 2.5 \times 10^{-7} = 0$

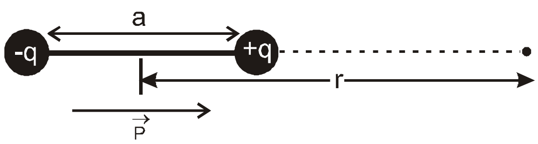
Electric dipole moment, P = (Magnitude of charge) × (Separation between charges)

$= 2.5 \times 10^{-7}$ [0.15 + 0.15] C m = $7.5 \times 10^{-8}$ C m

The direction of dipole moment is from B to A.

## 13.2 Electric Field Intensity Due to Dipole :

**(i) At the axial point :**



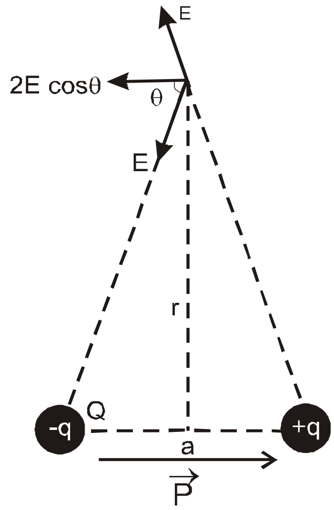

$$\vec{E} = \frac{Kq}{\left(r-\frac{a}{2}\right)^2} - \frac{Kq}{\left(r+\frac{a}{2}\right)^2} \text{ (along the } \vec{P}) = \frac{Kq(2ra)}{\left(r^2-\frac{a^2}{4}\right)^2}\hat{P}$$

If r >> a then, $\vec{E} = \frac{Kq2ra}{r^4}\hat{P} = \frac{2K\vec{P}}{r^3}$,

As the direction of electric field at axial position is along the dipole moment ( $\vec{P}$ )

So, $\vec{E}_{axial} = \frac{2K\vec{P}}{r^3}$

**(ii) Electric field at perpendicular Bisector (Equatorial Position)**




$E_{net} = 2\,E\cos\theta$ (along $-\hat{P}$)

$$\vec{E}_{net} = 2\left(\frac{Kq}{\left(\sqrt{r^2+\left(\frac{a}{2}\right)^2}\right)^2}\right)\frac{\frac{a}{2}}{\sqrt{r^2+\left(\frac{a}{2}\right)^2}}(-\hat{P})$$

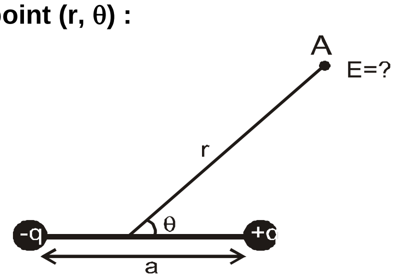
$$= \frac{Kqa}{\left(r^2+\left(\frac{a}{2}\right)^2\right)^{3/2}}(-\hat{P})$$

If $r >> a$ then $\vec{E}_{net} = \frac{KP}{r^3}(-\hat{P})$

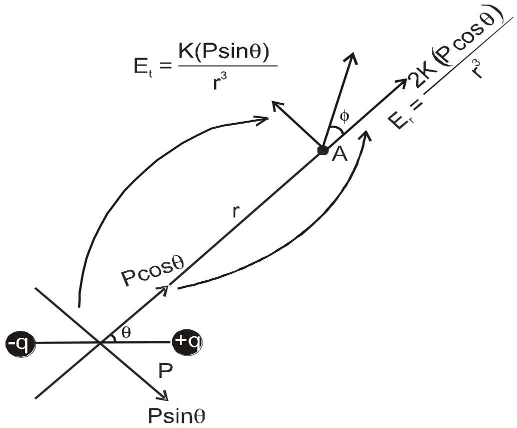
As the direction of $\vec{E}$ at equatorial position is opposite of $\vec{P}$ so we can write in vector form:

$$\vec{E}_{eqt} = -\frac{K\vec{P}}{r^3}$$

**(iii) Electric field at general point (r, θ) :**




For this, let's resolve the dipole moment $\vec{P}$ into components.




One component is along radial line (=P cosθ) and other component is $\perp_r$ to the radial line (= Psinθ)

From the given figure $E_{net} = \sqrt{E_r^2+E_t^2} = \sqrt{\left(\frac{2KP\cos\theta}{r^3}\right)^2+\left(\frac{KP\sin\theta}{r^3}\right)^2} = \frac{KP}{r^3}\sqrt{1+3\cos^2\theta}$

$$\tan\phi = \frac{E_t}{E_r} = \frac{\frac{KP\sin\theta}{r^3}}{\frac{2KP\cos\theta}{r^3}} = \frac{\tan\theta}{2}$$

$$E_{net} = \frac{KP}{r^3}\sqrt{1+3\cos^2\theta} \quad ; \quad \tan\phi = \frac{\tan\theta}{2}$$

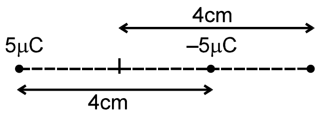
**Example 85.** The electric field due to a short dipole at a distance r, on the axial line, from its mid point is the same as that of electric field at a distance r', on the equatorial line, from its mid-point. Determine the ratio $\frac{r}{r'}$.

**Solution :** $\frac{1}{4\pi\varepsilon_0}\frac{2p}{r^3} = \frac{1}{4\pi\varepsilon_0}\frac{p}{r'^3}$ or $\frac{2}{r^3} = \frac{1}{r'^3}$ or $\frac{r^3}{r'^3} = 2$ or, $\frac{r}{r'} = 2^{1/3}$

**Example 86.** Two charges, each of 5 μC but opposite in sign, are placed 4 cm apart. Calculate the electric field intensity of a point that is at a distance 4 cm from the mid point on the axial line of the dipole.

**Solution :** We cannot use formula of short dipole here because distance of the point is comparable to the distance between the two point charges.

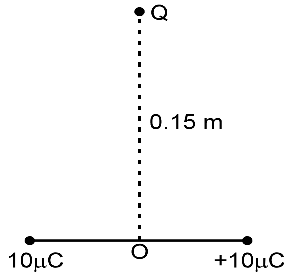
$q = 5 \times 10^{-6}$ C, $a = 4 \times 10^{-2}$ m, $r = 4 \times 10^{-2}$ m




$E_{res} = E_+ + E_- = \frac{K(5\mu C)}{(2cm)^2} - \frac{K(5\mu C)}{(6\ \ cm)^2} = \frac{144}{144\times10^{-8}}$ NC$^{-1}$ = $10^8$ NC$^{-1}$

**Example 87.** Two charges ±10μC are placed $5 \times 10^{-3}$ m apart as shown in figure. Determine the electric field at a point Q which is 0.15 m away from O, on the equatorial line.

**Solution :** In the given problem, r >> a

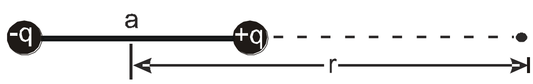



$\therefore \quad E = \frac{1}{4\pi\varepsilon_0}\frac{p}{r^3} = \frac{1}{4\pi\varepsilon_0}$ or $E = 9 \times 10^9 \times \frac{10\times10^{-6}\times5\times10^{-3}}{0.15\times0.15\times0.15}$ NC$^{-1}$ = $1.33 \times 10^5$ NC$^{-1}$

## 13.3 Electric Potential due to a small dipole :

**(i) Potential at axial position :**

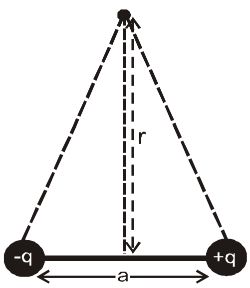
$$V = \frac{Kq}{\left(r-\frac{a}{2}\right)} + \frac{K(-q)}{\left(r+\frac{a}{2}\right)} \Rightarrow V = \frac{Kqa}{\left(r^2-\left(\frac{a}{2}\right)^2\right)}$$




If r >> a then $V = \frac{Kqa}{r^2}$ ; where, qa = p

$$\boxed{\therefore\ V_{axial} = \frac{KP}{r^2}}$$

**(ii) Potential at equatorial position :**




$$V = \frac{Kq}{\sqrt{r^2+\left(\frac{a}{2}\right)^2}} + \frac{K(-q)}{\sqrt{r^2+\left(\frac{a}{2}\right)^2}} = 0 \qquad \boxed{V_{eqt} = 0}$$

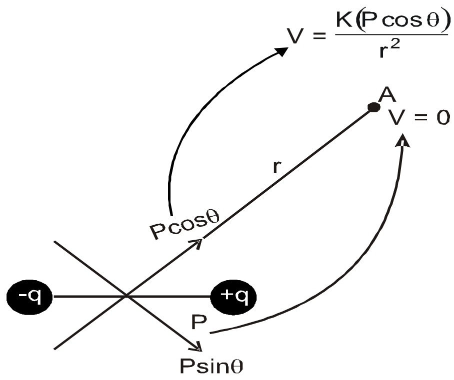
**(iii) Potential at general point (r, θ) :** Lets resolve the dipole moment $\vec{P}$ into components : Pcosθ along radial line and Psinθ $\perp_r$ to the radial line

For the Pcosθ component, the point A is an axial point,

So, potential at A due to Pcosθ = $\frac{K(P\cos\theta)}{r^2}$

And for Psinθ component, the point A is an equatorial point,

So potential at A due to Psinθ = 0

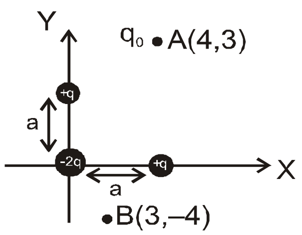



$V_{net} = \frac{K(P\cos\theta)}{r^2}$

$\therefore\ V = \frac{K\left(\vec{P}.\vec{r}\right)}{r^3}$

**Example 88.** (i) Find potential at point A and B due to the small charge - system fixed near origin. (Distance between the charges is negligible).

(ii) Find work done to bring a test charge $q_0$ from point A to point B, slowly. All parameters are in S.I. units.




**Solution :** (i) Dipole moment of the system is $\vec{P} = (qa)\,\hat{i} + (qa)\,\hat{j}$

Potential at point A due to the dipole

$$V_A = K\,\frac{(\vec{P}\cdot\vec{r})}{r^3} = \frac{K[(qa)\hat{i}+(qa)\hat{j}]\cdot(4\hat{i}+3\hat{j})}{5^3} = \frac{k(qa)}{125}\quad(7)$$

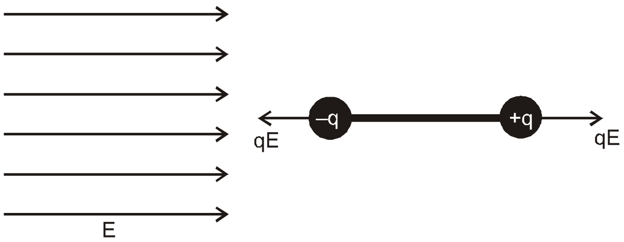
$$\Rightarrow\ V_B = \frac{K\left[(qa)\hat{i}+(qa)\hat{j}\right]\cdot\left(3\hat{i}-4\hat{j}\right)}{(5)^3} = \frac{-K(qa)}{125}$$

(ii) $$W_{A\to B} = U_B - U_A = q_0\,(V_B - V_A) = q_0\left[-\frac{K(qa)}{125} - \left(\frac{K(qa)(7)}{125}\right)\right]$$

$$\Rightarrow\ W_{A\to B} = \frac{-Kq q_0 a}{125}(8)$$

## 13.4 Dipole in uniform electric field

**(i) Dipole is placed along electric field :**




In this case, $F_{net} = 0$, $\tau_{net} = 0$, so it is an equilibrium state. And it is a stable equilibrium position.

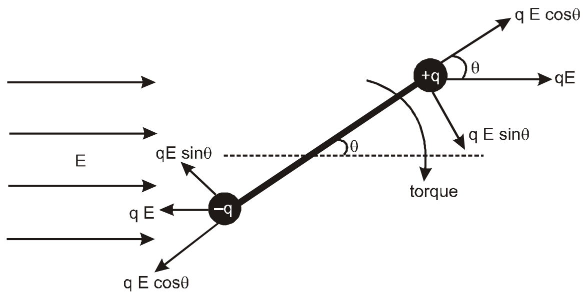
**(ii) If the dipole is placed at angle $\theta$ from $\vec{E}$ :**

In this case $F_{net} = 0$ but

Net torque $\tau = (qE\sin\theta)\ (a)$

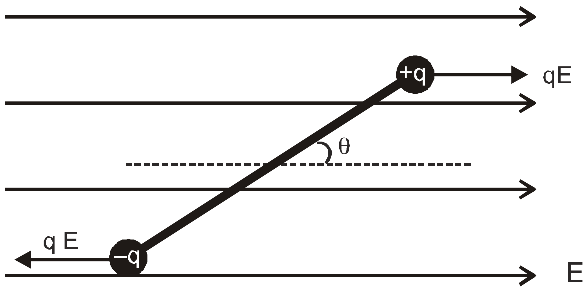
Here $qa = P \Rightarrow \tau = PE\sin\theta$

in vector form : $\vec{\tau} = \vec{P} \times \vec{E}$




**Example 89.** A dipole is formed by two point charge –q and +q, each of mass m, and both the point charges are connected by a rod of length $\ell$ and mass m. This dipole is placed in uniform electric field $\vec{E}$. If the dipole is disturbed by a small angle $\theta$ from stable equilibrium position, prove that its motion will be almost SHM. Also find its time period.

**Solution :** If the dipole is disturbed by $\theta$ angle, $\tau_{net} = -PE\sin\theta$ (Here -ve sign indicates that direction of torque is opposite to $\theta$)




If $\theta$ is very small, $\sin\theta \approx \theta$

$\therefore \quad \tau_{net} = -(PE)\theta$

$\tau_{net} \propto (-\theta)$ so motion will be almost SHM & C = PE (where, $P = q\ell$)

$$\because \quad T = 2\pi\sqrt{\frac{I}{C}}$$

$$\therefore \quad T = 2m^{-}\sqrt{\frac{\frac{m\ell}{12} + 2m\left(\frac{\ell}{2}\right)^2}{P.E}} = 2\pi\sqrt{\frac{\frac{m\ell}{12} + \frac{m\ell^2}{2}}{q\ell E}} = 2\pi\sqrt{\frac{7m\ell^2}{12q\ell E}} = 2\pi\sqrt{\frac{7m\ell}{12qE}}$$

$$T = \pi\sqrt{\frac{7m\ell}{3qE}}$$

**(iii) Potential energy of a dipole placed in uniform electric field :**

$$U_B - U_A = -\int_A^B \vec{F}\cdot\vec{dr} \qquad \text{(for translational motion)}$$

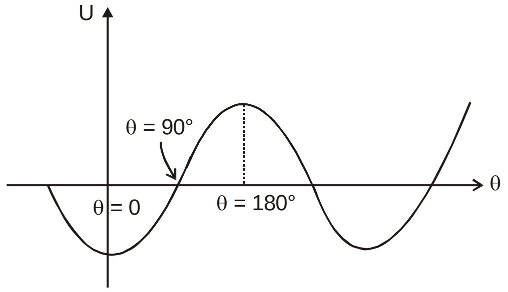
Here, $$U_B - U_A = -\int_A^B \vec{\tau}.\vec{d\theta} \qquad \text{(for rotational motion)}$$

In the case of dipole, at $\theta = 90°$, P.E. is assumed to be zero.

$$U_\theta - U_{90°} = -\int_{\theta=90°}^{\theta=\theta} (-PE\sin\theta)\ (d\theta) \qquad \text{(As the direction of torque is opposite of } \theta)$$

$U_\theta - 0 = -PE\cos\theta$

$\theta = 90°$ is chosen as reference, so that the lower limit comes out to be zero.



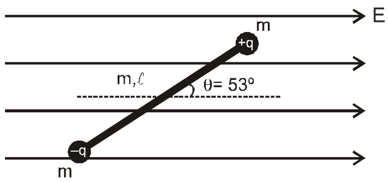

$$\boxed{U_\theta = -\vec{P}\cdot\vec{E}}$$

From the potential energy curve, we can conclude :

(i) At $\theta = 0$, there is minimum of P.E. so it is a stable equilibrium position.

(ii) At $\theta = 180°$, there is maxima of P.E. so it is a position of unstable equilibrium.

**Example 90.** Two point masses of mass m and equal and opposite charge of magnitude q are attached on the corners of a non-conducting uniform rod of mass m and the system is released from rest in uniform electric field E as shown in figure from $\theta = 53°$




(i) Find angular acceleration of the rod just after releasing

(ii) What will be angular velocity of the rod when it passes through stable equilibrium.

(iii) Find work required to rotate the system by 180°.

**Solution :** (i) $\tau_{net} = PE\sin 53° = I\alpha$

$$\therefore \quad \alpha = \frac{(q\ell)\ E\ \left(\frac{4}{5}\right)}{\frac{m\ell^2}{12} + m\left(\frac{\ell}{2}\right)^2 + m\left(\frac{\ell}{2}\right)^2} = \frac{48qE}{35\ m\ell}$$

(ii) From energy conservation : $K_i + U_i = K_f + U_f$

$$\therefore \quad 0 + (-PE\cos 53°) = \frac{1}{2} I\omega^2 + (-PE\cos 0°)$$

where $I = \frac{m\ell^2}{12} + m\left(\frac{\ell}{2}\right)^2 + m\left(\frac{\ell}{2}\right)^2 = \frac{7m\ell^2}{12}$ $\quad\therefore\quad$ $\frac{1}{2}I\omega^2 = PE\,(1-3/5) = \frac{2}{5}PE$

$\therefore\ \frac{1}{2} \times \frac{7m\ell^2}{12} \times \omega^2 = \frac{2}{5}\, q\ell E$ $\quad$ or $\quad \omega = \sqrt{\frac{48qE}{35\,m\ell}}$

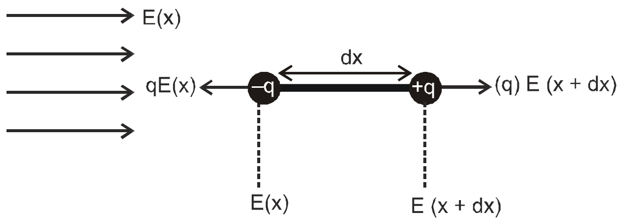
(iii) $\because\ W_{ext} = U_f - U_i$

$\therefore\ W_{ext} = (-PE\cos(180° + 53°)) - (-PE\cos 53°)$

or $W_{ext} = (q\ell)E\left(\frac{3}{5}\right) + (q\ell)E\left(\frac{3}{5}\right)$ $\quad\Rightarrow\quad W_{ext} = \left(\frac{6}{5}\right)q\ell E$

## 13.5 Dipole in non-uniform electric field :

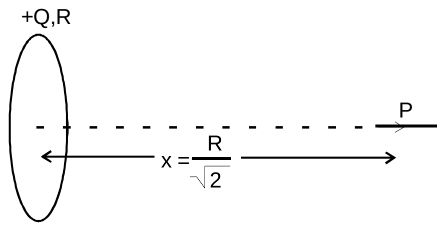
If the dipole is placed along $\vec{E}$, (shown in figure)




Then, Net force on the dipole : $F_{net} = qE(x + dx) - qE(x)$

$F_{net} = q\,\frac{E(x+dx)-E(x)}{dx}\,(dx)$ ; here $(q(dx) = P)$ $\quad \boxed{\therefore\ F_{net} = P\left(\frac{dE}{dx}\right)}$

**Example 91.** A short dipole is placed on the axis of a uniformly charged ring (total charge –Q, radius R) at a distance $\frac{R}{\sqrt{2}}$ from centre of ring as shown in figure. Find the Force on the dipole due to the ring



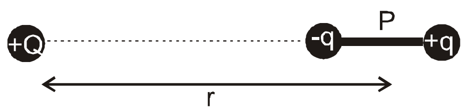

**Solution :** $\because\ F = P\left(\frac{dE}{dx}\right)$

$\therefore\ F = P\frac{d}{dx}\left(\frac{KQx}{(R^2+x^2)^{3/2}}\right)$ ; (at $x = \frac{R}{\sqrt{2}}$)

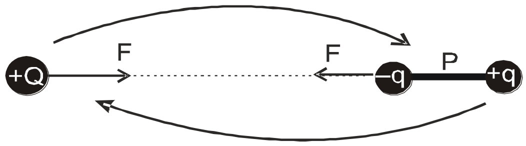
Solving we get, $F = 0$

## 13.6 Force between a dipole and a point charge :

**Example 92.** A short dipole of dipole moment P is placed near a point charge Q as shown in figure. Find force on the dipole due to the point charge



**Solution :**




Force on the point charge due to the dipole F = (Q) $E_{dipole}$

$$F = (Q)\left(\frac{2KP}{r^3}\right) \text{ (towards right)}$$

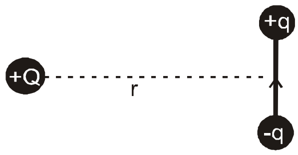
From action reaction concept, force on the dipole due to point charge will be equal to the force on charge due to dipole

$$F = \frac{2KPQ}{r^3} \text{ (towards left)}$$

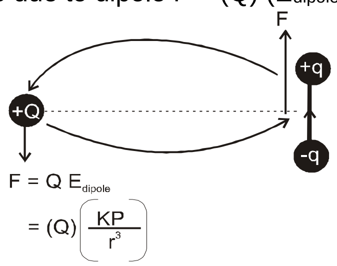
is force on dipole due to point Charge

## 13.7 Force between two dipoles :

**Example 93.** A short dipole of dipole moment P is placed near a point charge Q as shown in figure. Find force on the dipole due to the point charge.



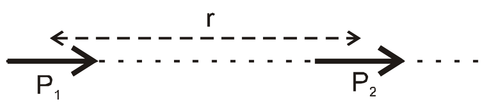

**Solution :** Force on the point charge due to dipole F = (Q) ($E_{dipole}$)



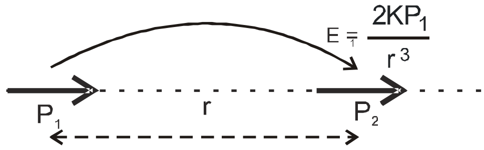

$$F = (Q)\left(\frac{KP}{r^3}\right)(\downarrow)$$

So force on the dipole due to the point charge will also be $F = \left(\frac{KPQ}{r^3}\right)(\uparrow)$ (but in opposite direction) as shown

**Example 94.** Find force on short dipole $P_2$ due to short dipole $P_1$ if they are placed at a distance r a part as shown in figure.




**Solution :** Force on $P_2$ due to $P_1$

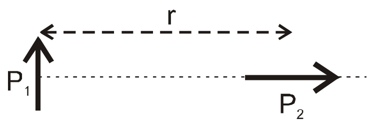



$$F_2 = (P_2)\left(\frac{dE_1}{dr}\right)$$

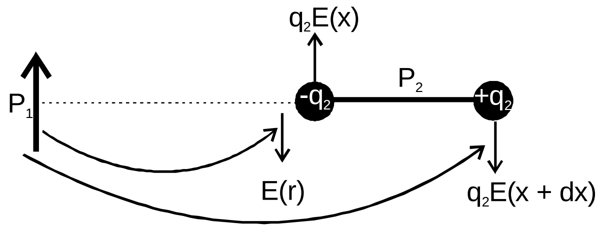
$$\therefore \quad F_2 = (P_2)\left(\frac{d}{dr}\left(\frac{2\ \ kP_1}{r^3}\right)\right) \quad \text{or} \quad F_2 = -\frac{6KP_1P_2}{r^4}$$

Here – sign indicates that this force will be attractive (opposite to r)

**Example 95.** Find force on short dipole $P_2$ due to short dipole $P_1$ if they are placed a distance r apart as shown in figure.




**Solution :**




$F_{net} = q_2 E(x + dx) - q_2 E(x)$

$F_{net} = q_2 \left(\frac{E(x+dx)-E(x)}{dx}\right)dx$

or $F_{net} = (P_2)\left(\frac{dE}{dx}\right)$

(Usually this formula is valid when the dipole is placed along $\vec{E}$. However, in this case also, we are getting the same formula)

$\therefore$ $F_{net} = (P_2)\left(\frac{d}{dr}\left(\frac{KP_1}{r^3}\right)\right)$

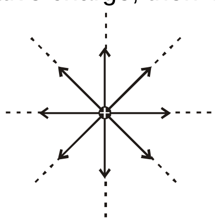
$\Rightarrow$ $F_{net} = \frac{3KP_1P_2}{r^4}$ (in magnitude) & (direction upwards)

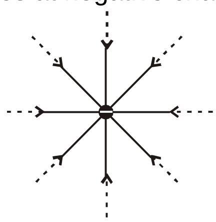
# 14. ELECTRIC LINES OF FORCE (ELOF)

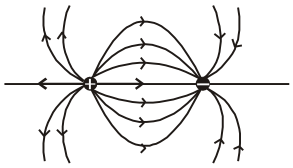
The line of force in an electric field is an imaginary line, the tangent to which at any point on it represents the direction of electric field at the given point.

## 14.1 Properties :

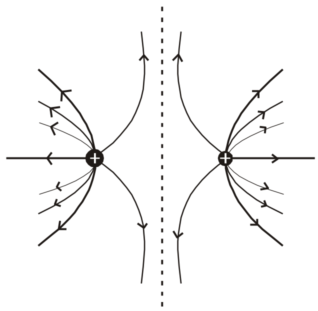
**(i)** Line of force originates out from a positive charge and terminates on a negative charge. If there is only one positive charge, then lines starts from positive charge and terminates at $\infty$. If there is only one negative charge, then lines starts from $\infty$ and terminates at negative charge.


ELOF of Isolated positive charge


ELOF of Isolated negative charge

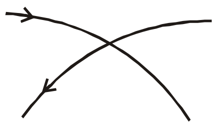


ELOF due to positive and negative charges



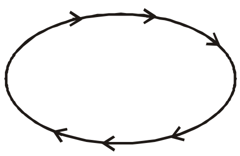
ELOF due to two positive charges

**(ii)** Two lines of force never intersect each other because there cannot be two directions of $\vec{E}$ at a single Point

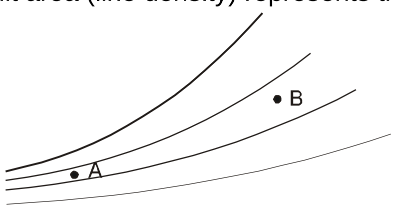


impossilble

**(iii)** Electric lines of force produced by static charges do not form closed loop.
If lines of force make a closed loop, then work done to move a +q charge along the loop will be non-zero. So it will not be conservative field. So these type of lines of force are not possible in electrostatics.



**(iv)** The Number of lines per unit area (line density) represents the magnitude of electric field.




$E_A > E_B$

If lines are dense $\Rightarrow$ E will be more
If Lines are rare $\Rightarrow$ E will be less and if E = O, no line of force will be found there

**(v)** Number of lines originating (terminating) at a charge is proportional to the magnitude of charge

**Example 96.** If number of electric lines of force from charge q are 10, then find out number of electric lines of force from 2q charge.

**Solution :** No. of ELOF $\propto$ charge

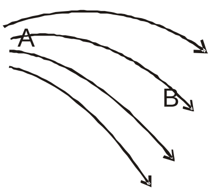
$10 \propto q \quad \Rightarrow \quad 20 \propto 2q$

So, number of ELOF will be 20.

**(vi)** Electric lines of force end or start perpendicularly on the surface of a conductor.

**(vii)** Electric lines of force never enter into conductors.

**Example 97.** Some electric lines of force are shown in figure. For points A and B

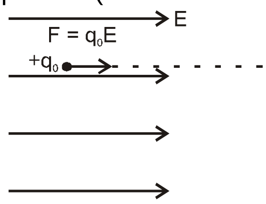



(A) $E_A > E_B$ (B) $E_B > E_A$ (C) $V_A > V_B$ (D) $V_B > V_A$

**Solution :** Lines are more dense at A, so $E_A > E_B$ In the direction of Electric field, potential decreases so $V_A > V_B$

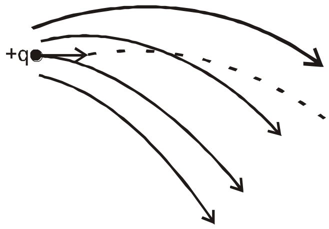
**Example 98.** If a charge is released in electric field, will it follow lines of force?

**Solution :** **Case I :** If lines of force are parallel (in uniform electric field) :




In this type of field, if a charge is released, force on it will be $q_0E$ and its direction will be along. So the charge will move in a straight line, along the lines of force.

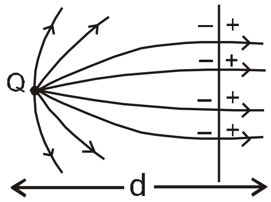
**Case II :** If lines of force are curved (in non-uniform electric field) :




The charge will not follow lines of force

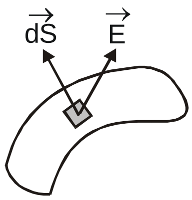
**Example 99.** A charge + Q is fixed at a distance d in front of an infinite metal plate. Draw the lines of force indicating the directions clearly.

**Solution :** There will be induced charge on two surfaces of conducting plate, so ELOF will start from +Q charge and terminate at conductor and then will again start from other surface of conductor.




## 15. ELECTRIC FLUX

Consider some surface in an electric field $\vec{E}$. Let us select a small area element $\vec{dS}$ on this surface. The electric flux of the field over the area element is given by $d\phi_E = \vec{E}.\vec{dS}$




**Direction of $\vec{dS}$ is normal to the surface. It is along $\hat{n}$**

or $d\phi_E = EdS\cos\theta$

or $d\phi_E = (E\cos\theta)\, dS$

or $d\phi_E = E_n\, dS$

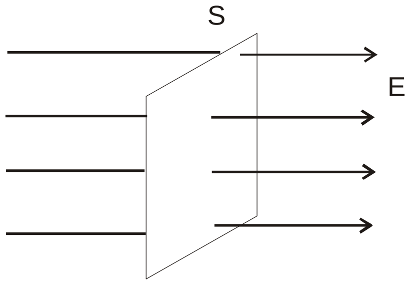
where $E_n$ is the component of electric field in the direction of $\vec{dS}$.

The electric flux over the whole area is given by $\phi_E = \int_S \vec{E}.\vec{dS} = \int_S E_n dS$

If the electric field is uniform over that area then $\phi_E = \vec{E}\cdot\vec{S}$

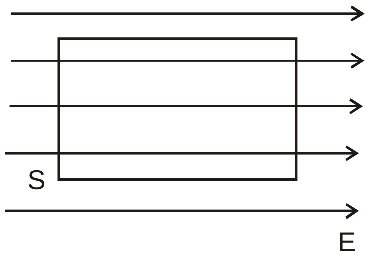
**Special Cases :**

**Case I :** If the electric field is normal to the surface, then angle of electric field $\vec{E}$ with normal will be zero




So $\phi$ = ES cos 0 or $\phi$ = ES

**Case II :** If electric field is parallel of the surface (grazing), then angle made by $\vec{E}$ with normal = 90º




So $\phi$ = ES cos 90º = 0

## 15.1 Physical Meaning :

The electric flux through a surface inside an electric field represents the total number of electric lines of force crossing the surface. It is a property of electric field

## 15.2 Unit

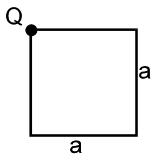
**(i)** The SI unit of electric flux is $Nm^2\,C^{-1}$ (Gauss) or $J\,m\,C^{-1}$.

**(ii)** Electric flux is a scalar quantity. (It can be positive, negative or zero)

**Example 100.** The electric field in a region is given by $\vec{E} = \frac{3}{5}E_0\hat{i} + \frac{4}{5}E_0\hat{j}$ with $E_0 = 2.0 \times 10^3$ N/C. Find the flux of this field through a rectangular surface of area $0.2m^2$ parallel to the Y–Z plane.

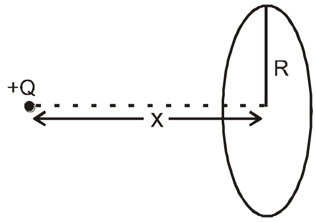
**Solution :** $\phi_E = \vec{E}\cdot\vec{S} = \left(\frac{3}{5}E_0\hat{i} + \frac{4}{5}E_0\hat{j}\right).\left(0.2\hat{i}\right) = 240\frac{N-m^2}{C}$

**Example 101.** A point charge Q is placed at the corner of a square of side a, then find the flux through the square.

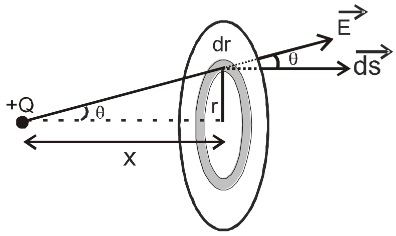



**Solution :** The electric field due to Q at any point of the square will be along the plane of square and the electric field lines are perpendicular to square ; so $\phi$ = 0. In other words we can say that no line is crossing the square so, flux = 0.

**Example 102.** Find the electric flux due to a point charge 'Q' through the circular region of radius R if the charge is placed on the axis of ring at a distance x.



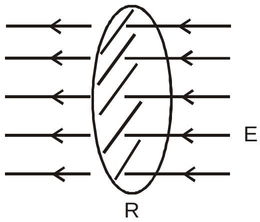

**Solution :** We can divide the circular region into small rings.




Lets take a ring of radius r and width dr. Flux through this small element $d\phi = E\, ds \cos\theta$

$$\therefore \quad \phi_{net} = \int E\, ds \cos\theta = \int_{r=o}^{r=R} \frac{KQ}{(x^2+r^2)} (2\pi\, r\, dr) \left(\frac{x}{\sqrt{x^2+r^2}}\right) = \frac{Q}{2\in_0}\left[1-\frac{x}{\sqrt{x^2+R^2}}\right]$$

**Case III :** Curved surface in uniform electric field. Suppose a circular surface of radius R is placed in a uniform electric field as shown.




Flux passing through the surface $\phi = E\,(\pi R^2)$

(ii) Now suppose, a hemispherical surface, is placed in the electric field. Flux through hemispherical surface:

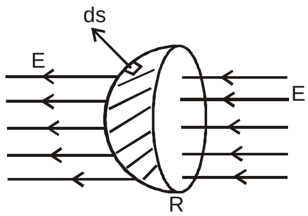
$\phi = \int E ds \cos\theta$

$\phi = E \int ds \cos\theta$

where, $\int ds \cos\theta$ is

projection of the spherical surface Area on base.

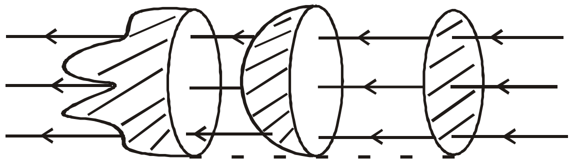
& $\int ds \cos\theta = \pi R^2$




So, $\phi = E(\pi R^2)$ = same Ans. as in previous case

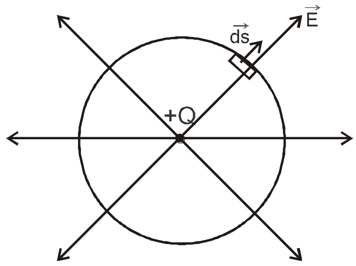
So, we can conclude that

**If the number of electric field lines passing through two surfaces are same, then flux passing through these surfaces will also be same, irrespective of the shape of surface**



$\phi_1 = \phi_2 = \phi_3 = E(\pi R^2)$

**Case IV : Flux through a closed surface :** Suppose there is a spherical surface and a charge 'q' is placed at centre.




$\therefore$ Flux through the spherical surface

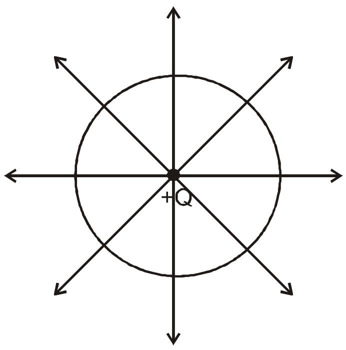
$$\phi = \int \vec{E}\cdot\vec{ds} = \int E\,ds \qquad (\text{as } \vec{E} \text{ is along } \vec{ds} \text{ (normal)})$$

$$\therefore \quad \phi = \frac{1}{4\pi\varepsilon_o}\frac{Q}{R^2}\int ds \quad \text{where, } \int ds = 4\pi R^2$$

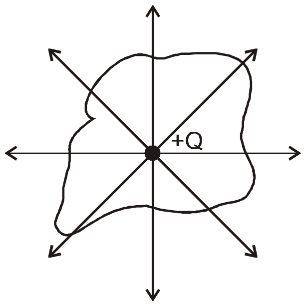
$$\phi = \left(\frac{1}{4\pi\varepsilon_0}\frac{Q}{R^2}\right)\left(4\pi R^2\right) \qquad \Rightarrow \qquad \phi = \frac{Q}{\varepsilon_o}$$

Now if the charge Q is enclosed by any other closed surface, still same no of lines of force will pass through the surface.

So, here also flux will be $\phi = \frac{Q}{\varepsilon_o}$. That's what Gauss Theorem is.




$\phi = \frac{Q}{\varepsilon_o}$




$\phi = \frac{Q}{\varepsilon_o}$

# 16. GAUSS'S LAW IN ELECTROSTATICS OR GAUSS'S THEOREM

This law was stated by a mathematician Karl F Gauss. This law gives the relation between the electric field at a point on a closed surface and the net charge enclosed by that surface. This surface is called Gaussian surface. It is a closed hypothetical surface. Its validity is shown by experiments. It is used to determine the electric field due to some symmetric charge distributions.

## 16.1 Statement and Details :

Gauss's law is stated as given below :

The surface integral of the electric field intensity over any closed hypothetical surface (called Gaussian surface) in free space is equal to $\frac{1}{\varepsilon_0}$ times the total charge enclosed within the surface. Here, $\varepsilon_0$ is the permittivity of free space.

If S is the Gaussian surface and $\sum_{i=1}^{n} q_i$ is the total charge enclosed by the Gaussian surface, then according to Gauss's law,

$$\phi_E = \oint \vec{E}\cdot\overrightarrow{dS} = \frac{1}{\varepsilon_0}\sum_{i=1}^{n} q_i .$$

The circle on the sign of integration indicates that the integration is to be carried out over the closed surface.

**Note :**

(i) Flux through Gaussian surface is independent of its shape.

(ii) Flux through Gaussian surface depends only on total charge present inside Gaussian surface.

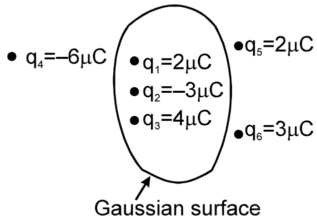
(iii) Flux through Gaussian surface is independent of position of charges inside Gaussian surface.

(iv) Electric field intensity at the Gaussian surface is due to all the charges present inside as well as outside the Gaussian surface.

(v) In a closed surface incoming flux is taken negative, while outgoing flux is taken positive, because $\hat{n}$ is taken positive in outward direction.

(vi) In a Gaussian surface, $\phi = 0$ does not imply E = 0 at every point of the surface but E = 0 at every point implies $\phi = 0$.

**Example 103.** Find out flux through the given Gaussian surface.




**Solution :** $\phi = \dfrac{Q_{in}}{\varepsilon_0} = \dfrac{2\mu C - 3\mu C + 4\mu C}{\varepsilon_0} = \dfrac{3 \times 10^{-6}}{\varepsilon_0}$ Nm²/C

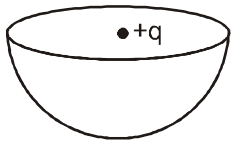
**Example 104.** If a point charge q is placed at the centre of a cube, then find out flux through any one face of cube.

**Solution :** Flux through all 6 faces = $\dfrac{q}{\varepsilon_0}$. Since, all the surfaces are symmetrical

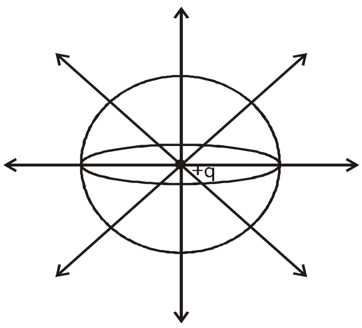
So, flux through one face = $\dfrac{1}{6}\dfrac{q}{\varepsilon_0}$

## 16.2 Flux through open surfaces using Gauss's Theorem :

**Example 105.** A point charge +q is placed at the centre of curvature of a hemisphere. Find flux through the hemispherical surface.




**Solution :** Lets put an upper half hemisphere. Now, flux passing through the entire sphere = $\dfrac{q}{\varepsilon_0}$

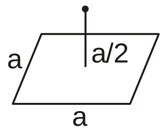


As the charge q is symmetrical to the upper half and lower half hemispheres, so half-half flux will emit from both the surfaces.

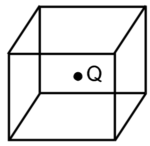
Flux emitting from lower half surface = $\frac{q}{2\varepsilon_0}$

Flux emitting from upper half surface = $\frac{q}{2\varepsilon_0}$

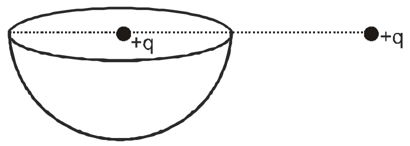
**Example 106.** A charge Q is placed at a distance a/2 above the centre of a horizontal, square surface of edge a as shown in figure. Find the flux of the electric field through the square surface.




**Solution :** We can consider imaginary faces of cube such that the charge lies at the centre of the cube. Due to symmetry, we can say that flux through the given area (which is one face of cube), $\phi = \frac{Q}{6\varepsilon_0}$

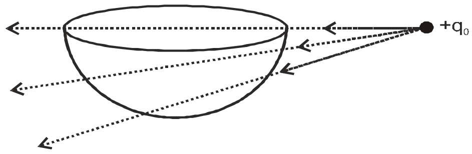



**Example 107.** Find flux through the hemispherical surface



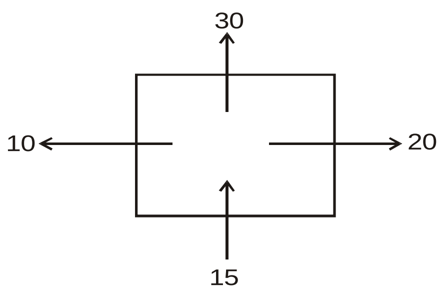

**Solution :** (i) Flux through the hemispherical surface due to +q = $\frac{q}{2\varepsilon_0}$ (we have seen in previous examples)

(ii) Flux through the hemispherical surface due to $+q_0$ is 0, because due to $+q_0$, field lines entering the surface = field lines coming out of the surface.




## 16.3 Finding $q_{in}$ from flux :

**Example 108.**




Flux (in S.I. units) coming out and entering a closed surface is shown in the figure. Find charge enclosed by the closed surface.

**Solution :** Net flux through the closed surface = + 20 + 30 + 10 -15 = 45 N.m$^2$/c

From Gauss`s theorem : $\phi_{net} = \frac{q_{in}}{\varepsilon_0}$

or $45 = \frac{q_{in}}{\varepsilon_0}$ $\quad\therefore\quad$ $q_{in} = (45)\varepsilon_0$

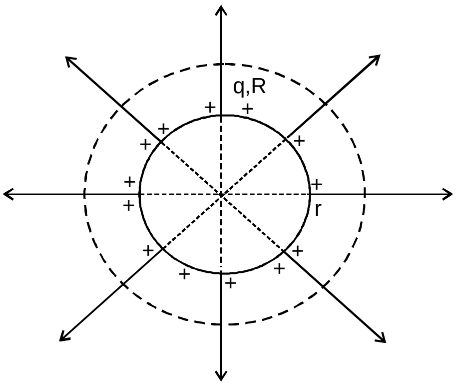
## 16.4 Finding electric field from Gauss`s Theorem :

From Gauss`s theorem, we can say $\int \vec{E}.\vec{ds} = \phi_{net} = \frac{q_{in}}{\varepsilon_0}$

### 16.4.1 Finding E due to a spherical shell :

**Electric field outside the Sphere :**

Since, electric field due to a shell will be radially outwards. So lets choose a spherical Gaussian surface Applying Gauss′s theorem for this spherical Gaussian surface,




$\int \vec{E}\vec{ds} = \phi_{net} = \frac{q_{in}}{\varepsilon_o} = \frac{q}{\varepsilon_o}$

$\downarrow$

$\int |\vec{E}||\vec{ds}|\cos 0$ (because $\vec{E}$ is normal to the surface)

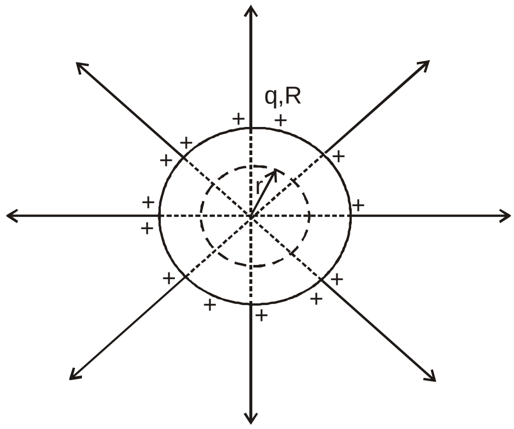
$\downarrow$

$E \int ds$ (because value of E is constant at the surface)

$E(4\pi r^2)$ ($\therefore \int ds$ total area of the spherical surface = $4\pi r^2$)

$\Rightarrow E(4\pi r^2) = \frac{q_{in}}{\varepsilon_0} \quad \therefore E_{out} = \frac{q}{4\pi\varepsilon_o r^2}$

**Electric field inside a spherical shell :** Lets choose a spherical Gaussian surface inside the shell. Applying Gauss`s theorem for this surface




$\int \vec{E}\vec{ds} = \phi_{net} = \frac{q_{in}}{\varepsilon_0} = 0$

$\downarrow$

$\int |\vec{E}||\vec{ds}|\cos 0$

$\downarrow$

$E \int ds$

$\downarrow$

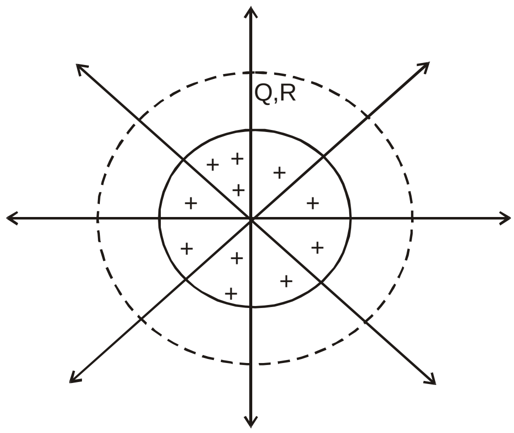
$E(4\pi r^2) \Rightarrow E(4\pi r^2) = 0$

$\therefore E_{in} = 0$

### 16.4.2 Electric field due to solid sphere (having uniformly distributed charge Q and radius R) :

**Electric field outside the sphere :**

Direction of electric field is radially outwards, so we will choose a spherical Gaussian surface Applying Gauss`s theorem




$\int \vec{E}\vec{ds} = \phi_{net} = \frac{q_{in}}{\varepsilon_o} = \frac{Q}{\varepsilon_o}$

$\downarrow$

$\int |\vec{E}||\vec{ds}|\cos 0$

$\downarrow$

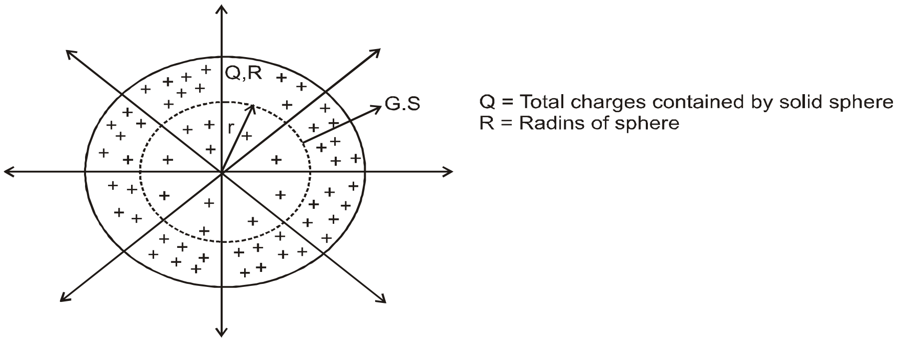
$E\int ds$

$\downarrow$

$E\,(4\pi r^2) \quad \Rightarrow \quad E\,(4\pi r^2) = \dfrac{Q}{\varepsilon_o}$

or $\quad E_{out} = \dfrac{Q}{4\pi\varepsilon_o r^2}$

**Electric field inside a solid sphere :**




For this choose a spherical Gaussian surface inside the solid sphere Applying Gauss`s theorem for this surface

$$\oint \vec{E}\,\vec{ds} = \phi_{net} = \frac{q_{in}}{\varepsilon_o} = \frac{\dfrac{Q}{\frac{4}{3}\pi R^3} \times \frac{4}{3}\pi r^3}{\varepsilon_o} = \frac{Qr^3}{\varepsilon_o R^3}$$

$\downarrow$

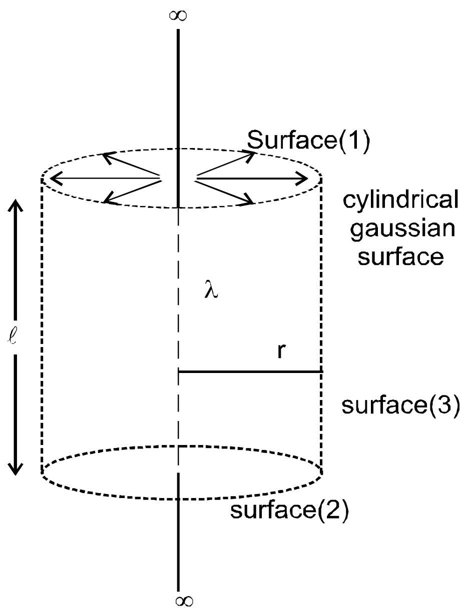
$\oint E\,ds$

$\downarrow$

$E\,(4\pi r^2) \qquad \Rightarrow E(4\pi r^2) = \dfrac{Qr^3}{\varepsilon_o R^3}$

$E = \dfrac{Qr}{4\pi\varepsilon_o R^3} \qquad \therefore \quad E_{in} = \dfrac{kQ}{R^3}\,r$

### 16.4.3 Electric field due to infinite line charge (having uniformly distributed charged of charge density $\lambda$) :



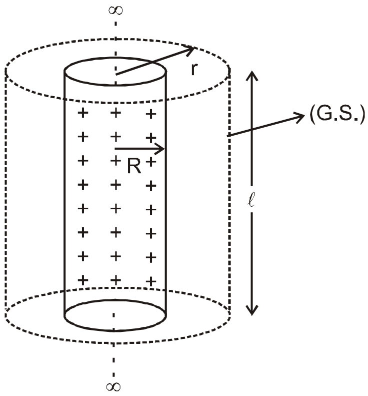
Electric field due to infinitely wire is radial so we will choose cylindrical Gaussian surface as shown is figure:

$$\phi_{net} \rightarrow \phi_1 = 0 \quad \phi_2 = 0 \quad \phi_3 \neq 0 = \frac{q_{in}}{\varepsilon_o} = \frac{\lambda\ell}{\varepsilon_o}$$

$$\phi_3 = \int \vec{E}.\vec{ds} = \int E\,ds = E\int ds = E\,(2\pi r\ell)$$

$$\because \quad E\,(2\pi r\ell) = \frac{\lambda\ell}{\varepsilon_o} \qquad \therefore \qquad E = \frac{\lambda}{2\pi\varepsilon_o r} = \frac{2k\lambda}{r}$$

### 16.4.4 Electric field due to infinitely long charged tube (having uniform surface charge density σ and radius R)) :




**(i) E outside the tube :** Lets choose a cylindrical Gaussian surface of length $\ell$ :

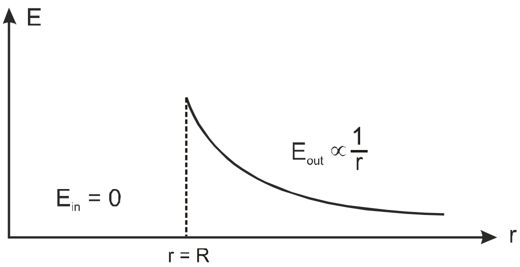
$$\therefore \phi_{net} = \frac{q_{in}}{\varepsilon_o} = \frac{\sigma 2\pi R\ell}{\varepsilon_o} \qquad \Rightarrow \qquad E_{out} \times 2\pi r\ell = \frac{\sigma 2\pi R\ell}{\varepsilon_o} \qquad \therefore \qquad E = \frac{\sigma R}{r\,\varepsilon_0}$$

**(ii) E inside the tube :**

Lets choose a cylindrical Gaussian surface inside the tube.

$$\phi_{net} = \frac{q_{in}}{\varepsilon_o} = 0$$

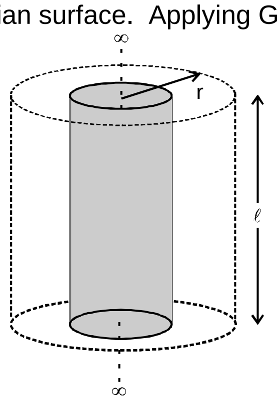
So $E_{in} = 0$




### 16.4.5 E due to infinitely long solid cylinder of radius R (having uniformly distributed charge in volume (volume charge density ρ)) :

**(i) E at outside point :-**

Lets choose a cylindrical Gaussian surface. Applying Gauss`s theorem :

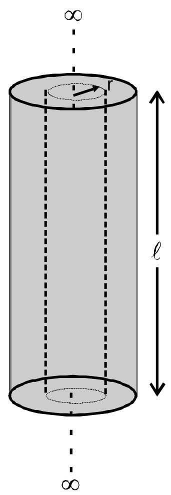


$$E \times 2\pi r\ell = \frac{q_{in}}{\varepsilon_o} = \frac{\rho \times \pi R^2 \ell}{\varepsilon_o}$$

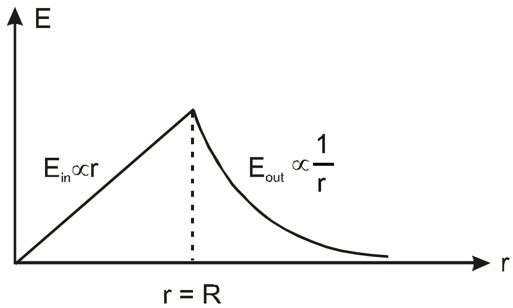
$$E_{out} = \frac{\rho R^2}{2r\,\varepsilon_0}$$

**(ii) E at inside point :**

Lets choose a cylindrical Gaussian surface inside the solid cylinder. Applying Gauss`s theorem




$$E \times 2\pi r\ell = \frac{q_{in}}{\varepsilon_o} = \frac{\rho \times \pi r^2 \ell}{\varepsilon_o} \Rightarrow E_{in} = \frac{\rho r}{2\varepsilon_0}$$




## 17. CONDUCTOR AND IT'S PROPERTIES [FOR ELECTROSTATIC CONDITION]

(i) Conductors are materials which contain large number of free electrons which can move freely inside the conductor.

(ii) In electrostatics conductors are always equipotential surfaces.

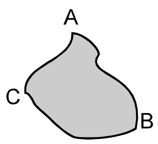
(iii) Charge always resides on outer surface of conductor.

(iv) If there is a cavity inside the conductor having no charge then charge will always reside only on outer surface of conductor.

(v) Electric field is always perpendicular to conducting surface.

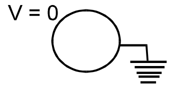
(vi) Electric lines of force never enter into conductors.

(vii) Electric field intensity near the conducting surface is given by formula $\vec{E} = \frac{\sigma}{\varepsilon_0}\hat{n}$




$$\overrightarrow{E_A} = \frac{\sigma_A}{\varepsilon_0}\hat{n} \;;\; \overrightarrow{E_B} = \frac{\sigma_B}{\varepsilon_0}\hat{n} \text{ and } \overrightarrow{E_C} = \frac{\sigma_C}{\varepsilon_0}\hat{n}$$

(viii) When a conductor is grounded its potential becomes zero.




(ix) When an isolated conductor is grounded then its charge becomes zero.

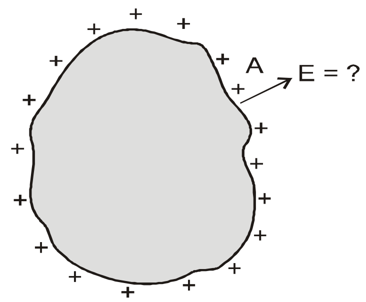
(x) When two conductors are connected there will be charge flow till their potentials become equal.

(xi) Electric pressure : Electric pressure at the surface of a conductor is given by formula

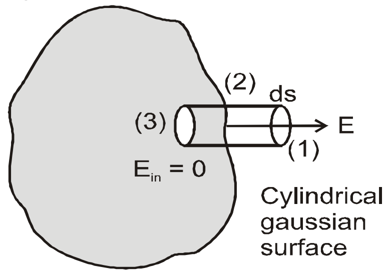
$P = \dfrac{\sigma^2}{2\varepsilon_0}$, where $\sigma$ is the local surface charge density.

## FINDING FIELD DUE TO A CONDUCTOR

Suppose we have a conductor and at any 'A', local surface charge density = $\sigma$. We have to find electric field just outside the conductor surface.




For this, let’s consider a small cylindrical Gaussian surface, which is partly inside and partly outside the conductor surface, as shown in figure. It has a small cross section area ds and negligible height.




Applying Gauss's theorem for this surface :

$$\phi_{net} = \frac{q_{in}}{\varepsilon_0} = \frac{\sigma\, ds}{\varepsilon_0}$$

flux through surface (1) $\phi_1 = Eds$ (because $\vec{E}$ is normal to the surface of conductor)

flux through surface (2) $\phi_2 = 0$ ($\vec{E}$ is normal to curved Gaussian surface )

flux through surface (3) $\phi_3 = 0$ (as E inside the conductor = 0)

So, $Eds = \dfrac{\sigma ds}{\varepsilon_0} \Rightarrow E = \dfrac{\sigma}{\varepsilon_0}$

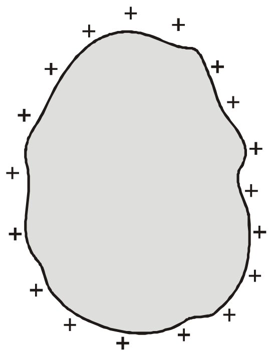
Electric field just outside the surface of conductor :

$E = \dfrac{\sigma}{\varepsilon_0}$ (direction will be normal to the surface)

in vector form: $\vec{E} = \dfrac{\sigma}{\varepsilon_0}\hat{n}$ (Here, $\hat{n}$ = unit vector normal to the conductor surface)

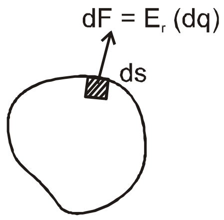
## ELECTROSTATIC PRESSURE AT THE SURFACE OF THE CONDUCTOR :

Suppose a conductor is given some charge. Due to repulsion, all the charges will reach the surface of the conductor. But the charges will still repel each other. So an outward force will be felt by each charge due to others. Due to this force, there will be some pressure at the surface, which is called electrostatic pressure.



To find the electrostatic pressure, lets take a small surface element having Area 'ds'.

Force on this element due to the remaining charges :




$$dF = \begin{pmatrix}\text{electric field at} \\ \text{that place due to} \\ \text{remainig charges}\end{pmatrix}\begin{pmatrix}\text{charge of} \\ \text{the small} \\ \text{element}\end{pmatrix}$$

Let electric field at that point due to the remaining charges = $E_r$

and charge of the small element = $dq = \sigma ds$

$\Rightarrow\ dF = (E_r)\ (dq) = (E_r)\ (\sigma ds)$

So, pressure on this small element

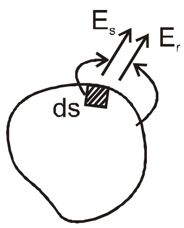
$$P = \frac{dF}{ds} = \frac{(E_r)(\sigma ds)}{ds} \Rightarrow P = (E_r)\ (\sigma) \qquad .....(1)$$

Now to find pressure, we have to find $E_r$ (electric field at that position due to the remaining charges)

Suppose,

Electric field due to the small element near the surface = $E_s$

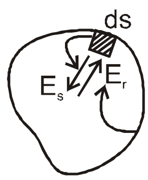
Electric field due to the remaining part near the surface = $E_r$




At a point just outside the surface, electric field due to the small element ($E_s$) will be normally out wards, and electric field due to the remaining part ($E_r$) will also be normally out wards.

So Net electric field just outside the surface $= E_s + E_r$ and we have proved that electric field just outside the conductor surface $= \frac{\sigma}{\varepsilon_o}$

$$\Rightarrow\ E_s + E_r = \frac{\sigma}{\varepsilon_o} \qquad .....(2)$$




Now, lets see the electric field just inside the metal surface. Here, electric field due to the remaining charges ($E_r$) will be in the same direction (normally outward), but the electric field due to the small element will be in opposite direction (normally inward)

So net electric field just inside the metal surface $= E_r - E_s$ and we know that electric field inside a conductor = 0

$$\text{So, } E_r - E_s = 0 \ \Rightarrow\ E_r = E_s \qquad .....(3)$$

from eqn. (2) and eqn. (3), we can say that :

$$2E_r = \frac{\sigma}{\varepsilon_o} \quad \Rightarrow \quad E_r = \frac{\sigma}{2\varepsilon_o}$$

Now, we can easily find the pressure from eqn.(1)

$$P = (E_r)\ (\sigma) = \frac{\sigma}{2\varepsilon_o}\ (\sigma) = \frac{\sigma^2}{2\varepsilon_o}$$

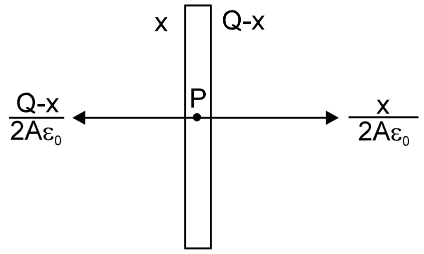
So, electrostatic pressure at the surface of the conductor $P = \frac{\sigma^2}{2\varepsilon_o}$

where, $\sigma$ = local surface charge density.

**Example 109.** Prove that if an isolated (isolated means no charges are near the sheet) large conducting sheet is given a charge then the charge distributes equally on its two surfaces.

**Solution :** Let there is x charge on left side of sheet and Q–x charge on right side of sheet.

Since, point P lies inside the conductor, so $E_P = O$



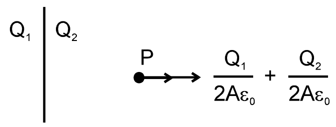

or $\frac{x}{2A\varepsilon_O} - \frac{Q-x}{2A\varepsilon_O} = 0 \quad \Rightarrow \quad \frac{2x}{2A\varepsilon_O} = \frac{Q}{2A\varepsilon_O}$

$\Rightarrow \; x = \frac{Q}{2} \;\; \& \;\; Q - x = \frac{Q}{2}$

So, charge is equally distributed on both sides.

**Example 110.** If an isolated infinite sheet contains charge $Q_1$ on its one surface and charge $Q_2$ on its other surface, then prove that electric field intensity at a point in front of sheet will be $\frac{Q}{2A\varepsilon_O}$, where $Q = Q_1 + Q_2$

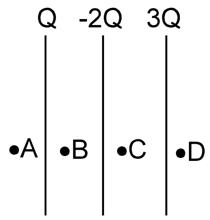
**Solution :** Electric field at point P :




$$\vec{E} = \vec{E}_{Q_1} + \vec{E}_{Q_2} = \frac{Q_1}{2A\varepsilon_0}\hat{n} + \frac{Q_2}{2A\varepsilon_0}\hat{n} = \frac{Q_1 + Q_2}{2A\varepsilon_0}\hat{n} = \frac{Q}{2A\varepsilon_0}\hat{n}$$

[This shows that the resultant field due to a sheet depends only on the total charge of the sheet and not on the distribution of charge on individual surfaces].

**Example 111.** Three large conducting sheets placed parallel to each other at finite distance contain charges Q, –2Q and 3Q respectively. Find electric field at points A, B, C and D




**Solution :** (i) $E_A = E_Q + E_{-2Q} + E_{3Q}$. Here $E_Q$ means electric field due to 'Q'.

$E_A = \frac{(Q - 2Q + 3Q)}{2A\varepsilon_0} = \frac{2Q}{2A\varepsilon_0} = \frac{Q}{A\varepsilon_0}$, towards left

(ii) $E_B = \frac{Q - (-2Q + 3Q)}{2A\varepsilon_0} = 0$

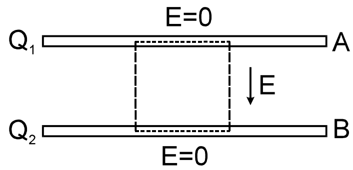
(iii) $E_C = \frac{(Q - 2Q) - (3Q)}{2A\varepsilon_0} = \frac{-4Q}{2A\varepsilon_0} = \frac{-2Q}{A\varepsilon_0}$, towards right

$\Rightarrow \dfrac{2Q}{A\varepsilon_0}$ towards left

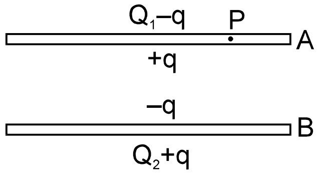
(iv) $E_D = \dfrac{(Q-2Q+3Q)}{2A\varepsilon_0} = \dfrac{2Q}{2A\varepsilon_0} = \dfrac{Q}{A\varepsilon_0}$, towards right

**Example 112.** Two conducting plates A and B are placed parallel to each other. A is given a charge $Q_1$ and B a charge $Q_2$. Prove that the charges on the inner facing surfaces are of equal magnitude and opposite sign.

**Solution :** Consider a Gaussian surface as shown in figure. Two faces of this closed surface lie completely inside the conductor where the electric field is zero. The flux through these faces is, therefore, zero. The other parts of the closed surface which are outside the conductor are parallel to the electric field and hence the flux on these parts is also zero. The total flux of the electric field through the closed surface is, therefore zero. From Gauss's law, the total charge inside this closed surface should be zero. The charge on the inner surface of A should be equal and opposite to that on the inner surface of B.




The distribution should be like the one shown in figure. To find the value of q, consider the field at a point P inside the plate A. Suppose, the surface area of the plate (one side) is A. Using the equation, $E = \sigma / (2\varepsilon_0)$, the electric field at P




Due to the charge $Q_1 - q = \dfrac{Q_1 - q}{2A\varepsilon_0}$ (downward)

Due to the charge $+ q = \dfrac{q}{2A\varepsilon_0}$ (upward),

Due to the charge $- q = \dfrac{q}{2A\varepsilon_0}$ (downward),

and due to the charge $Q_2 + q = \dfrac{Q_2 + q}{2A\varepsilon_0}$ (upward).

The net electric field at P due to all the four charged surfaces is (in the downward direction)

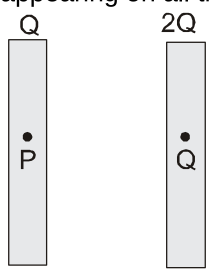
$$E_p = \frac{Q_1 - q}{2A\varepsilon_0} - \frac{q}{2A\varepsilon_0} + \frac{q}{2A\varepsilon_0} - \frac{Q_2 + q}{2A\varepsilon_0}$$

As the point P is inside the conductor, this field should be zero. Hence,

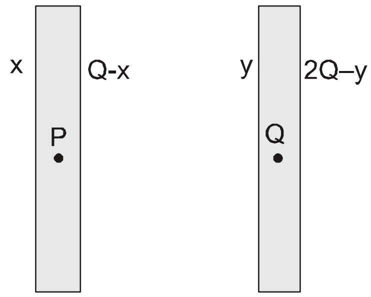
$Q_1 - q - q + q - Q_2 - q = 0$ or, $q = \dfrac{Q_1 - Q_2}{2}$

This result is a special case of the following result. When charged conducting plates are placed parallel to each other, the two outermost surfaces get equal charges and the facing surfaces get equal and opposite charges.

**Example 113.** Two large parallel conducting sheets (placed at finite distance) are given charges Q and 2Q respectively. Find out charges appearing on all the surfaces.



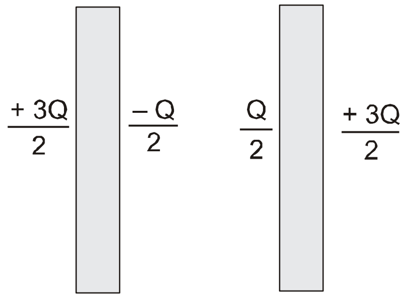
**Solution :** Let there is x amount of charge on left side of first plate. So, on its right side charge will be Q–x. Similarly, for second plate there is y charge on left side and 2Q – y charge is on right side,




$E_p = 0$ (By property of conductor)

$$\Rightarrow \quad \frac{x}{2A\varepsilon_o} - \left\{\frac{Q-x}{2A\varepsilon_o} + \frac{y}{2A\varepsilon_o} + \frac{2Q-y}{2A\varepsilon_o}\right\} = 0$$

We can also say that charge on left side of P = charge on right side of P



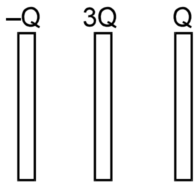

$$x = Q - x + y + 2Q - y \quad \Rightarrow \quad x = \frac{3Q}{2},\ Q - x = \frac{-Q}{2}$$

Similarly, for point Q : x + Q – x + y = 2Q – y

$\Rightarrow$ y = Q/2, 2Q – y = 3Q/2

So, final charge distribution of plates is

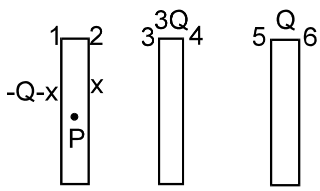
**Example 114.** Figure shows three large metallic plates with charges – Q, 3Q and Q respectively. Determine the final charges on all the surfaces.




**Solution :** We assume that charge on surface 2 is x. Following conservation of charge, we see that surfaces 1 has charge (– Q – x). The electric field inside the metal plate is zero. So, field at P is zero.

Resultant field at P

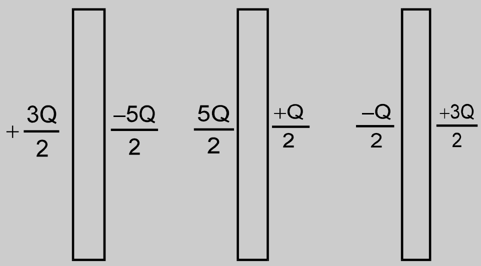
$E_P = 0$



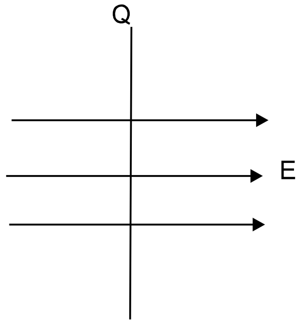

$$\Rightarrow \quad \frac{-Q-x}{2A\varepsilon_0} = \frac{x+3Q+Q}{2A\varepsilon_0} \quad \text{or} \quad -Q - x = x + 4Q$$

$$\therefore \quad x = \frac{-5Q}{2}$$

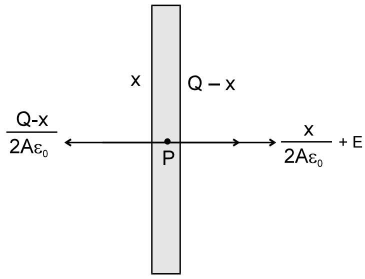
**Note :** We see that charges on the facing surfaces of the plates are of equal magnitude and opposite sign. This can be in general proved by Gauss theorem also. Remember this, it is important result. Thus the final charge distribution on all the surfaces is as shown in figure :



**Example 115.** An isolated conducting sheet of area A and carrying a charge Q is placed in a uniform electric field E, such that electric field is perpendicular to sheet and covers all the sheet. Find out charges appearing on its two surfaces.




**Solution :** Let there is x charge on left side of plate and Q – x charge on right side of plate




$\because \quad E_P = 0$

$$\therefore \quad \frac{x}{2A\varepsilon_0} + E = \frac{Q-x}{2A\varepsilon_0} \qquad \text{or} \qquad \frac{x}{A\varepsilon_0} = \frac{Q}{2A\varepsilon_0} - E$$

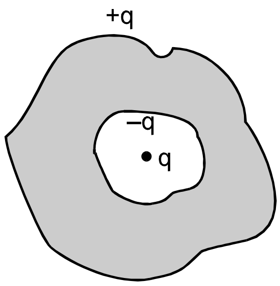
$$\therefore \quad x = \frac{Q}{2} - EA\varepsilon_0 \qquad \text{and} \qquad Q - x = \frac{Q}{2} + EA\varepsilon_0$$

So, charge on one side is $\frac{Q}{2}$ – $EA\varepsilon_0$ and other side $\frac{Q}{2}$ + $EA\varepsilon_0$

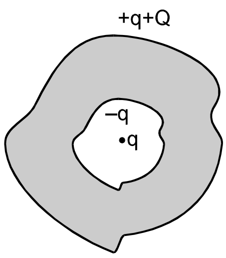
**Note :** Solve this question for Q = 0 without using the above answer and match that answer with the answer that you will get by putting Q = 0 in the above answer.

**17.1 Some other important results for a closed conductor:**

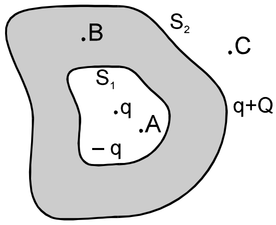
**(i)** If a charge q is kept in the cavity then –q will be induced on the inner surface and +q will be induced on the outer surface of the conductor (it can be proved using Gauss theorem)



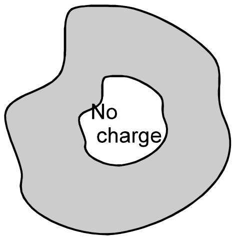

**(ii)** If a charge q is kept inside the cavity of a conductor and conductor is given a charge Q then –q charge will be induced on inner surface and total charge on the outer surface will be q + Q. (it can be proved using Gauss theorem)




**(iii)** Resultant field, due to q (which is inside the cavity) and induced charge on $S_1$, at any point outside $S_1$ (like B, C) is zero. Resultant field due to **q + Q on $S_2$ and any other charge outside $S_2$**, at any point inside of surface $S_2$ (like A, B) is zero

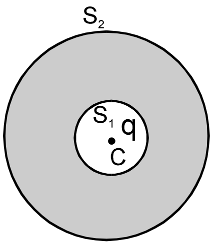



**(iv)** Resultant field in a charge free cavity in a closed conductor is zero. There can be charges outside the conductor and on the surface also. Then also, this result is true. No charge will be induced on the inner most surface of the conductor.



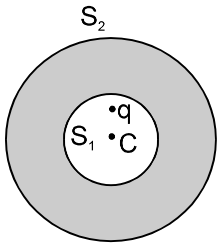

**(v)** Charge distribution for different types of cavities in conductors

(A)



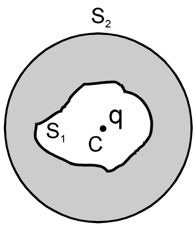

charge is at the common centre
($S_1$, $S_2 \rightarrow$ spherical)

(B)



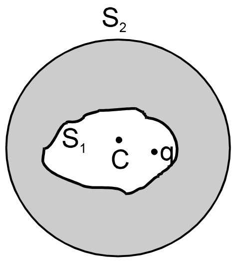

charge is not at the common centre
($S_1$, $S_2 \rightarrow$ spherical)

(C)



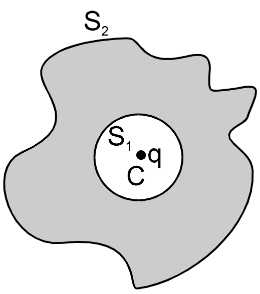

charge is at the centre of $S_2$
($S_2 \rightarrow$ spherical)

(D)



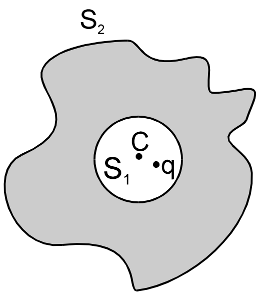

charge is not at the centre of $S_2$
($S_2 \rightarrow$ spherical)

(E)



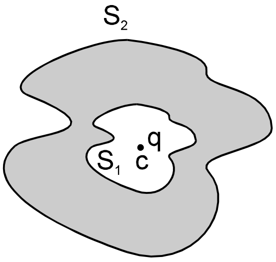

charge is at the centre of
$S_1$ (Spherical)

(F)



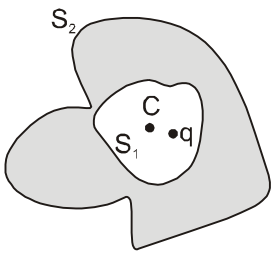

charge not at the centre
of $S_1$(Spherical)

(G)




charge is at the geometrical centre

(H)



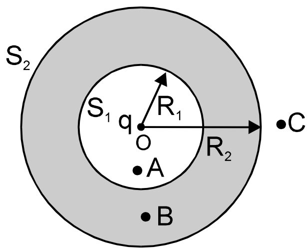

charge is not at the geometrical centre

Using the result that $\vec{E}_{res}$ in the conducting material should be zero and using result (iii) we can show that

| Case | A | B | C | D | E | F | G | H |
|---|---|---|---|---|---|---|---|---|
| $S_1$ | Uniform | Nonuniform | Nonuniform | Nonuniform | Uniform | Nonuniform | Nonuniform | Nonuniform |
| $S_2$ | Uniform | Uniform | Uniform | Uniform | Nonuniform | Nonuniform | Nonuniform | NonUniform |

**Note :** In all cases, charge on inner surface $S_1$ = –q and on outer surface $S_2$ = q. The distribution of charge on '$S_1$' will not change even if some charges are kept outside the conductor (i.e. outside the surface $S_2$). But the charge distribution on '$S_2$' may change if some charges(s) is/are kept outside the conductor.

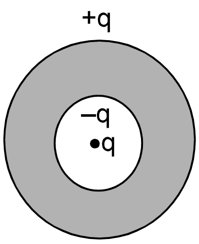
**Example 116.** An uncharged conductor of inner radius $R_1$ and outer radius $R_2$ contains a point charge q at the centre as shown in figure




(i) Find $\vec{E}$ and V at points A,B and C

(ii) If a point charge Q is kept outside the sphere at a distance 'r' (>>$R_2$) from centre, then find out resultant force on charge Q and charge q.

**Solution :** At point A :




$V_A = \frac{Kq}{OA} + \frac{Kq}{R_2} + \frac{K(-q)}{R_1}$ , $\vec{E}_A = \frac{Kq}{OA^3}\overrightarrow{OA}$

**Note :** Electric field at 'A' due to –q of $S_1$ and +q of $S_2$ is zero individually because they are uniformly distributed

At point B : $V_B = \frac{Kq}{OB} + \frac{K(-q)}{OB} + \frac{Kq}{R_2} = \frac{Kq}{R_2}$ , $E_B = 0$

At point C : $V_C = \frac{Kq}{OC}$ , $\vec{E}_C = \frac{Kq}{OC^3}\overrightarrow{OC}$

(ii) Force on point charge Q :

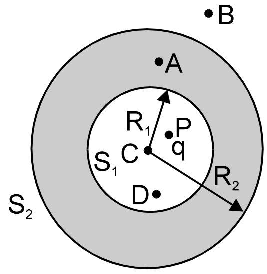
(**Note :** Here, force on 'Q' will be only due to 'q' of $S_2$ (see result (iii) )

$\vec{F}_Q = \frac{KqQ}{r^2}\hat{r}$ (r = distance of 'Q' from centre 'O')

Force on point charge q:

$\vec{F}_q = 0$ (using result (iii) & charge on $S_1$ uniform)

**Example 117.** An uncharged conductor of inner radius $R_1$ and outer radius $R_2$ contains a point charge q placed at point P (not at the centre) as shown in figure. Find out the following :




(i) $V_C$ (ii) $V_A$ (iii) $V_B$ (iv) $E_A$
(v) $E_B$ (vi) Force on charge Q, if it is placed at B.

**Solution :** (i) $V_C = \dfrac{Kq}{CP} + \dfrac{K(-q)}{R_1} + \dfrac{Kq}{R_2}$

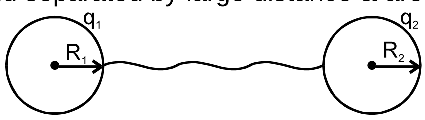
**Note :** –q on $S_1$ is non-uniformly distributed. Still it produces potential $\dfrac{K(-q)}{R_1}$ at 'C' because 'C' is at distance '$R_1$' from each point of '$S_1$'.

(ii) $V_A = \dfrac{Kq}{R_2}$ (iii) $V_B = \dfrac{Kq}{CB}$

(iv) $E_A$ = O (point is inside metallic conductor)

(v) $E_B = \dfrac{Kq}{CB^2}\,\hat{CB}$ (vi) $F_Q = \dfrac{KQq}{CB^2}\,\hat{CB}$

**(vi) Sharing of charges :** Two conducting hollow spherical shells of radii $R_1$ and $R_2$ having charges $Q_1$ and $Q_2$ respectively and separated by large distance & are joined by a conducting wire




Let final charges on spheres are $q_1$ and $q_2$ respectively. Potential on both spherical shell becomes equal after joining. Therefore,

$$\frac{Kq_1}{R_1} = \frac{Kq_2}{R_2} \Rightarrow \frac{q_1}{q_2} = \frac{R_1}{R_2} \qquad .....(i)$$

and, $q_1 + q_2 = Q_1 + Q_2$ .....(ii)

from (i) and (ii) : $q_1 = \dfrac{(Q_1+Q_2)R_1}{R_1+R_2}$ $q_2 = \dfrac{(Q_1+Q_2)R_2}{R_1+R_2}$

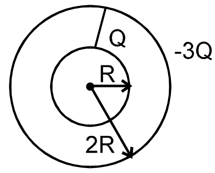
ratio of charges : $\dfrac{q_1}{q_2} = \dfrac{R_1}{R_2} \Rightarrow \dfrac{\sigma_1 4\pi R_1^2}{\sigma_2 4\pi R_2^2} = \dfrac{R_1}{R_2}$

$\therefore$ ratio of surface charge densities : $\dfrac{\sigma_1}{\sigma_2} = \dfrac{R_2}{R_1}$

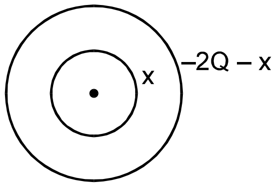
**Ratio of final charges :** $\dfrac{q_1}{q_2} = \dfrac{R_1}{R_2}$

**Ratio of final surface charge densities :** $\dfrac{\sigma_1}{\sigma_2} = \dfrac{R_2}{R_1}$

**Example 118.** The two conducting spherical shells are joined by a conducting wire which is cut after some time when charge stops flowing. Find out the charge on each sphere after that.




**Solution :** After cutting the wire, the potential of both the shells is equal




Thus, potential of inner shell, $V_{in} = \dfrac{Kx}{R} + \dfrac{K(-2Q-x)}{2R} = \dfrac{k(x-2Q)}{2R}$

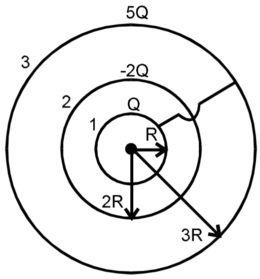
and potential of outer shell, $V_{out} = \dfrac{Kx}{2R} + \dfrac{K(-2Q-x)}{2R} = \dfrac{-KQ}{R}$

As, $V_{out} = V_{in}$

$\Rightarrow \quad \dfrac{-KQ}{R} = \dfrac{K(x-2Q)}{2R} \quad \Rightarrow \quad -2Q = x - 2Q \quad \Rightarrow \quad x = 0$

So, charge on inner spherical shell = 0 and outer spherical shell = – 2Q.

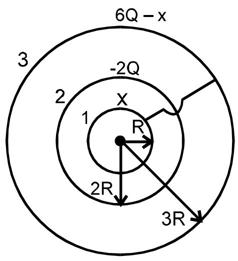
**Example 119.** Find charge on each spherical shell after joining the inner most shell and outer most shell by a conducting wire. Also find charges on each surface.




**Solution :** Let the charge on the innermost sphere be x.

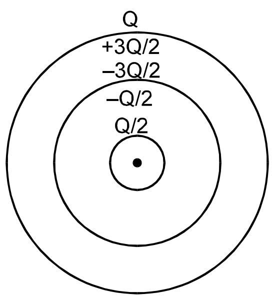
Finally potential of shell 1 = Potential of shell 3

$$\therefore \quad \frac{Kx}{R} + \frac{K(-2Q)}{2R} + \frac{K(6Q-x)}{3R} = \frac{Kx}{3R} + \frac{K(-2Q)}{3R} + \frac{K(6Q-x)}{3R}$$




$3x - 3Q + 6Q - x = 4Q \; ; \; 2x = Q \; ; \; x = \dfrac{Q}{2}$

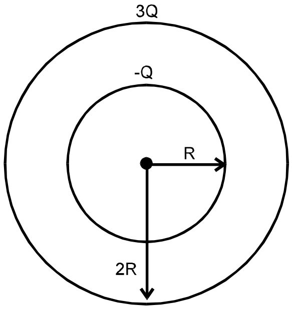
$\therefore$ Charge on innermost shell = $\dfrac{Q}{2}$ & Charge on outermost shell = $\dfrac{5Q}{2}$




Charge on middle shell = –2Q

$\therefore$ Final charge distribution is as shown in figure.

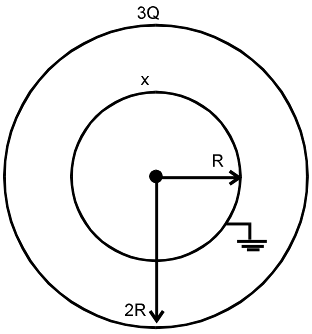
**Example 120.** Two conducting hollow spherical shells of radii R and 2R carry charges – Q and 3Q respectively. How much charge will flow into the earth if inner shell is grounded ?




**Solution :** When inner shell is grounded to the Earth then the potential of inner shell will become zero because potential of the Earth is taken to be zero.

$$\frac{Kx}{R} + \frac{K3Q}{2R} = 0$$

or $x = \frac{-3Q}{2}$, (the charge that has appeared on inner shells after grounding)




$\Rightarrow \frac{-3Q}{2} - (-Q) = \frac{-Q}{2}$ [hence, charge flown into the Earth $= \frac{Q}{2}$]

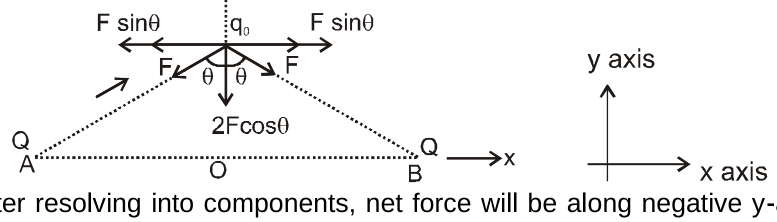
**Example 121.** An isolated conducting sphere of charge Q and radius R is connected to a similar uncharged sphere (kept at a large distance) by using a high resistance wire. After a long time, what is the amount of heat loss ?

**Solution :** When two conducting spheres of equal radii are connected, charge is equally distributed on them (Result VI). So, we can say that heat loss of system :

$$\Delta H = U_i - U_f = \left(\frac{Q^2}{8\pi\varepsilon_0 R} - 0\right) - \left(\frac{Q^2/4}{8\pi\varepsilon_0 R} + \frac{Q^2/4}{8\pi\varepsilon_0 R}\right) = \frac{Q^2}{16\pi\varepsilon_0 R}$$

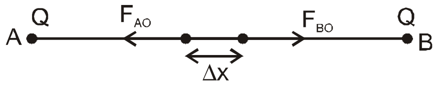
**Problem 1.** Two equal positive point charges 'Q' each are fixed at points B(a, 0) and A(–a, 0). Another negative point charge $q_0$ is also placed at O(0, 0) then prove that the equilibrium at 'O' is
(i) Stable for displacement in Y-direction.
(ii) Unstable for displacement in X-direction.

**Solution :** (i) When charge is shifted along y-axis:
Let x-y direction as :-




After resolving into components, net force will be along negative y-axis so the particle will return to its original position. So, it is stable equilibrium

(ii) When negative charge $q_0$ is shifted along x-axis.



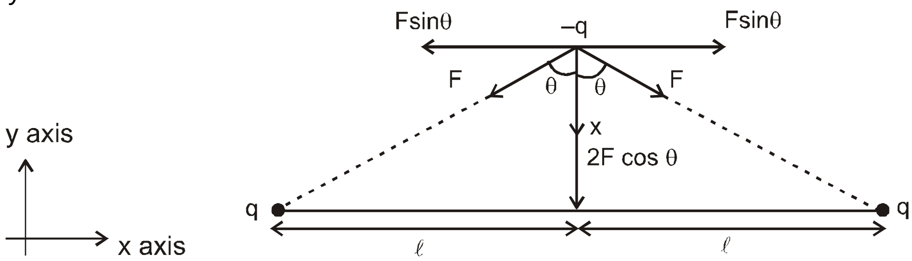

Initially, $\vec{F}_{AO} + \vec{F}_{BO} = \vec{0}$ $\Rightarrow$ $\left|\vec{F}_{AO}\right| = \left|\vec{F}_{BO}\right| = \dfrac{KQq_0}{d^2}$

When charge $q_0$ is slightly shifted towards + x axis by small distance $\Delta x$ then $\left|\vec{F}_{BO}\right| > \left|\vec{F}_{AO}\right|$

Also, these forces are attractive forces (due to negative charge)
Therefore, the particle will move towards positive x-axis and will not return to its original position So, it is unstable equilibrium for negative charge.

**Problem 2.** A particle of mass m and charge –q is located midway between two fixed charged particles each having a charge q and a distance $2\ell$ apart. Prove that the motion of the particle will be SHM if it is displaced slightly along perpendicular bisector and released. Also find its time period.

**Solution :** Let x-y direction is taken as :




Particle is shifted along y-axis by a small displacement x.
After resolving component of forces between q and –q charges :
By figure. $F_{net}$ in x-axis = 0 [$F_{net}$ = net force on –q charge]

Net force on – q charge in y direction = $-2F\cos\theta = -2.\ \dfrac{kqq}{(x^2+\ell^2)} . \dfrac{x}{(x^2+\ell^2)^{1/2}}$

$|\vec{F}| = \dfrac{2Kq^2x}{(x^2+\ell^2)^{3/2}}$ $\Rightarrow$ $ma = \dfrac{2Kq^2x}{\ell^3}$ (for $x << \ell$) (a = acceleration of – q charge)

$\Rightarrow$ $a = \dfrac{2Kq^2}{m\ell^3} . x$ (downwards)

This is equation of S.H.M. ($a = -\omega^2 x$)
So, time period of this charge (–q) :

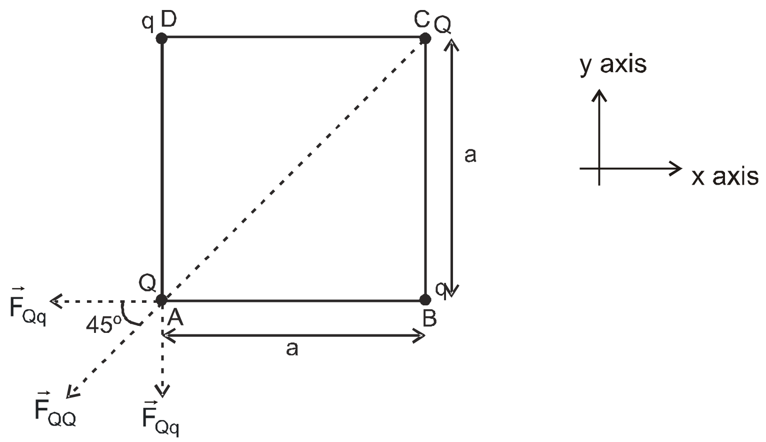
$T = 2\pi\sqrt{\dfrac{m\ell^3}{2Kq^2}}$ **Ans.**

**Problem 3.** Two charges Q each, are placed at two opposite corners of a square. A charge q is placed at each of the other two corners.

(a) If the resultant force on Q is zero, how are Q and q related ?

(b) Could q be chosen to make the resultant force on each charge zero ?

**Solution :** (a) Let on a square ABCD, charges are placed as shown




Now, forces on charge Q (at point A) due to other charge are $\vec{F}_{QQ}$, $\vec{F}_{Qq}$ and $\overrightarrow{F}_{Qq}$ respectively as shown in figure.

$F_{net}$ on Q $= \vec{F}_{Q.Q} + \vec{F}_{Qq} + \vec{F}_{Qq}$ (at point A)

But $F_{net} = 0$

So, $\Sigma F_x = 0$

$\Sigma F_x = -F_{QQ}\cos 45° - F_{Qq}$

$\Rightarrow \frac{KQ^2}{(\sqrt{2}a)^2} \cdot \frac{1}{\sqrt{2}} + \frac{KQq}{a^2} = 0 \Rightarrow q = -\frac{Q}{2\sqrt{2}}$ **Ans.**

(b) For resultant force on each charge to be zero :

From previous data, force on charge Q is zero when $q = -\frac{Q}{2\sqrt{2}}$ If for this value of charge q, force on q is zero, then and only then the value of q exists for which the resultant force on each charge is zero.

**Force on q :**

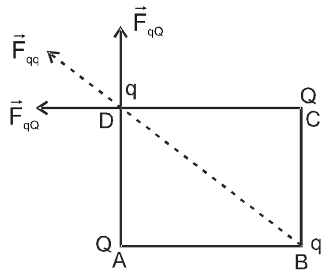
Forces on charge q (at point D) due to other three charges are $\vec{F}_{qQ}$, $\vec{F}_{qq}$ and $\vec{F}_{qQ}$ respectively as shown in figure.

Net force on charge q :

$\vec{F}_{net} = \vec{F}_{qq} + \vec{F}_{qQ} + \vec{F}_{qQ}$ But $\vec{F}_{net} = 0$

So, $\Sigma F_x = 0$

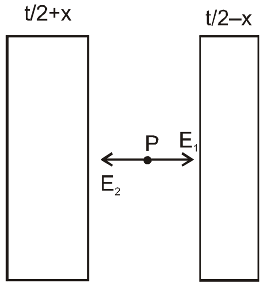
$\Sigma F_x = -\frac{Kq^2}{(\sqrt{2}a)^2} \cdot \frac{1}{\sqrt{2}} - \frac{KQq}{(a)^2} \Rightarrow q = 2\sqrt{2} - Q$




But from previous condition, $q = -\frac{Q}{2\sqrt{2}}$

So, no value of q makes the resultant force on each charge zero.

**Problem 4.** An infinitely large non-conducting sheet of thickness t and uniform volume charge density $\rho$ is given in which left half of the sheet contains charge density $\rho$ and right half contains charge density. Find the electric field at the symmetry plane of this sheet




**Solution :** We can consider two sheets of thickness $\left(\frac{t}{2}-x\right)$ and $\left(\frac{t}{2}+x\right)$

When a point lies inside the sheet.

Net electric field at point P : E = $E_1 - E_2 = \frac{Q_1}{2A\varepsilon_0} - \frac{Q_2}{2A\varepsilon_0}$

[$Q_1$ : charge on left sheet; $Q_2$ = charge of right sheet]

$$= \frac{A\rho\left(\frac{t}{2}+x\right) - 2\rho A\left(\frac{t}{2}-x\right)}{2A\varepsilon_0} = \frac{\rho\left[3x-\frac{t}{2}\right]}{2\varepsilon_0}$$

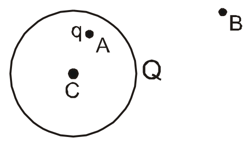
At the symmetry plane, x = 0 So, $E = -\frac{\rho t}{4\varepsilon_0}$ **Ans.**

**Problem 5.** Figure shows a uniformly charged thin non-conducting sphere of total charge Q and radius R. If point charge q is situated at point 'A' which is at a distance r < R from the centre of the sphere, then find out following:

(i) Force acting on charge q.

(ii) Electric field at centre of sphere.

(iii) Electric field at point B.

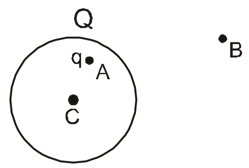



**Solution :** (i) Electric field inside a hollow sphere = 0

$\therefore$ Force on charge q.

F = qE = q × 0 = 0

(ii) Net electric field at centre of sphere




$\vec{E}_{net} = \vec{E}_1 + \vec{E}_2$

$E_1$ = field due to sphere = 0

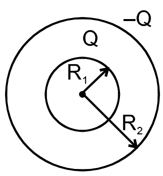
$E_2$ = field due to this charge = $\frac{Kq}{r^2}$

$\therefore$ $E_{net} = \frac{Kq}{r^2}$

(iii) Electric field at B due to charge on sphere, $\vec{E}_1 = \frac{KQ}{r_1^2}\hat{r}_1$ and due to charge q at A, $\vec{E}_2 = \frac{Kq}{r_2^2}\hat{r}_2$

So, $\vec{E}_{net} = \vec{E}_1 + \vec{E}_2 = \frac{KQ}{r_1^2}\hat{r}_1 + \frac{Kq}{r_2^2}\hat{r}_2$ where $r_1$ = CB and $r_2$ = AB

**Problem 6.** Figure shows two concentric spheres of radii $R_1$ and $R_2$ ($R_2 > R_1$) which contain uniformly distributed charges Q and –Q respectively. Find out electric field intensities at the following positions :

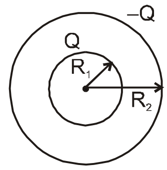



(i) $r < R_1$ (ii) $R_1 \le r < R_2$ (iii) $r \ge R_2$

**Solution :** Net electric field = $E_1 + E_2$

$E_1$ = field due to sphere of radius $R_1$

$E_2$ = field due of sphere of radius $R_2$




(i) $E_1 = 0, E_2 = 0$ $\therefore$ $E_{net} = 0$

(ii) $E_1 = \frac{KQ}{r^2}$, $E_2 = 0$ $\Rightarrow$ $\vec{E} = \frac{Kq}{r^2}\hat{r}$

(iii) $\vec{E}_1 = \frac{Kq}{r^2}\hat{r}$ $\vec{E}_2 = \frac{Kq}{r^2}(-\hat{r})$ $\Rightarrow$ $\vec{E}_{net} = \vec{E}_1 + \vec{E}_2 = 0$

**Problem 7.** A solid non conducting sphere of radius R and uniform volume charge density $\rho$ has centre at origin. Find out electric field intensity in vector form at following positions :

(i) (R, 0, 0) (ii) (0, 0, R/2) (iii) (R, R, R)

**Solution :** For uniformly charged non-conducting sphere, electric field inside the sphere :

$$\vec{E} = k\frac{Q\vec{r}}{R^3} = \frac{\rho\vec{r}}{3\varepsilon_o} \text{ (for } r < R)$$

and electric field outside the sphere

$$\vec{E}_o = \frac{KQ}{r^2}.\hat{r} = \frac{1}{4\pi\varepsilon_o}.\frac{\rho\frac{4}{3}\pi R^3}{r^2}\hat{r} = \frac{\rho R^3}{3\varepsilon_o r^2}.\hat{r} \text{ (for } r \geq R)$$

(i) (R, 0, 0) means it is at the surface $\vec{r} = R\hat{r}$ and $\hat{r} = \hat{i}$

$$\therefore \vec{E}_o = \frac{\rho R^3}{3\varepsilon_o R^2}(\hat{i}) = \frac{\rho R}{3\varepsilon_o}.\hat{i}$$

(ii) $(0, 0, \frac{R}{2})$

means point is inside the sphere

$$\vec{r} = \frac{R}{2}\hat{k} \Rightarrow \vec{E} = \frac{\rho R}{6\varepsilon_o}\hat{k}$$

(iii) For position (R, R, R)

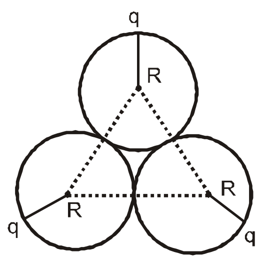
$$\vec{r} = R(\hat{i} + \hat{j} + \hat{k}) \Rightarrow \hat{r} = \frac{\hat{i}+\hat{j}+\hat{k}}{\sqrt{3}}, \ r = R\sqrt{3}$$

means point (R,R,R) is outside the sphere

$$\therefore \vec{E} = \frac{\rho R^3}{3\varepsilon_o(3R^2)}.\frac{(\hat{i}+\hat{j}+\hat{k})}{\sqrt{3}} = \frac{\rho R}{9\sqrt{3}\varepsilon_o}(\hat{i}+\hat{j}+\hat{k})$$ **Ans.**

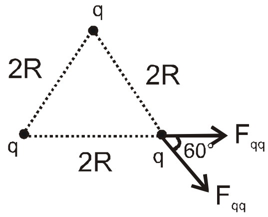
**Problem 8.** Three identical spheres each having a charge q (uniformly distributed) and radius R, are kept in such a way that each touches the other two. Find the magnitude of the electric force on any one sphere due to other two.

**Solution :** Given three identical spheres each having a charge q and radius R are kept as shown.




For any external point, sphere behaves like a point charge. So it becomes a triangle having point charges at its corners.

$$|\vec{F}_{qq}| = \frac{kq^2}{4R^2}$$




So, net force (F) = $2.\ \frac{kq^2}{4R^2}.\cos\frac{60}{2} = 2.\ \frac{kq^2}{4R^2}\ \frac{\sqrt{3}}{2} = \frac{\sqrt{3}}{4}.\ \frac{kq^2}{R^2}$ **Ans.**

**Problem 9.** A uniform electric field of 20 N/C exists in the vertically downward direction. Find the increase in the electric potential as one goes up through a height of 40cm.

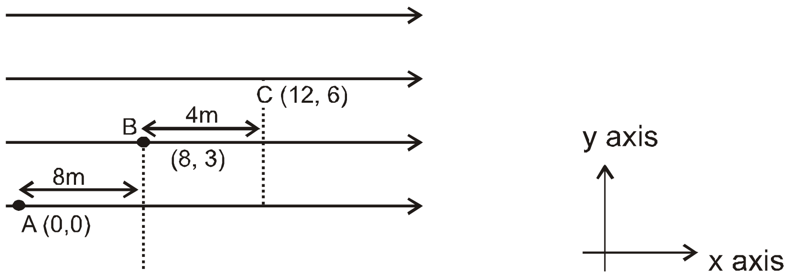
**Solution :** $E = -\frac{dv}{dr} \Rightarrow dv = -\vec{E}.d\vec{r}$

for $\vec{E}$ = constant $\Rightarrow$ $\Delta V = -\vec{E}.\vec{\Delta r}$

$\Delta V = -20(-\hat{j}).(40 \times 10^{-2})\hat{j} = 8$ volts.

**Problem 10.** An electric field of 10 N/C exists along the x-axis in space. Calculate the potential difference $V_B - V_A$, where the points A and B are given by –
(a) A = (0,0) ; B = (8m , 3m) (b) A = (8m, 3m) ; B = (12m , 6m)
(c) A = (0,0) ; B = (12m , 6m)

**Solution :** Electric field in x - axis means $\vec{E} = 10\hat{i}$

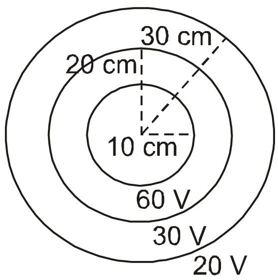



(a) $|\Delta V_{AB}| = \vec{E}.\vec{d} = 10\hat{i}\ .\ 8\ \hat{i} = 80$ V $\Rightarrow$ $V_B - V_A = -80V$

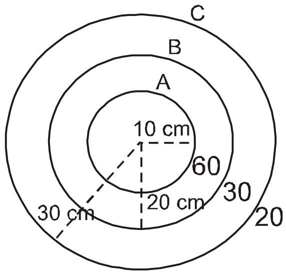
(b) $|\Delta V_{BC}| = \vec{E}.\vec{d} = 10\hat{i}\ .\ 4\hat{i} = 40$ volt $\Rightarrow$ $V_C - V_B = -40V$

(c) $|\Delta V_{AC}| = \vec{E}.\vec{d} = 10\hat{i}\ .\ 12\hat{i} = 120$ volt $\Rightarrow$ $V_C - V_A = -120$ V

**Problem 11.** Some equi-potential surfaces are shown in figure. What can you say about the magnitude and the direction of the electric field ?




**Solution :** We know, that the electric field is always perpendicular to equipotential surface. So, making electric field lines perpendicular to the surface, we find that these lines are originating from the centre. So, the field is similar to that due to a point charge placed at the centre. So, comparing the given potentials with that due to point charge, we have,



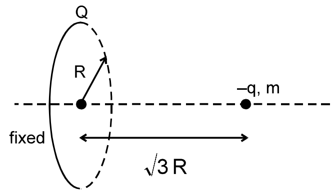

$V = \frac{KQ}{r} \Rightarrow KQ = V_A r_A = V_B r_B = V_C r_C = 6$ V-m

Hence, electric field at distance r can be given by

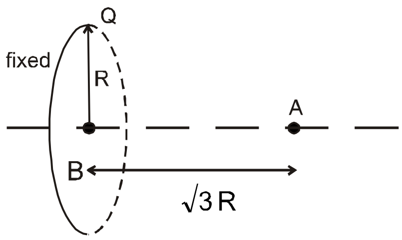
$E = \frac{KQ}{r^2} = \frac{6}{r^2}$ V/m

As the electric field lines are directed towards the decreasing potential. So, electric field is along radially outward direction.

**Problem 12.** A point charge of charge –q and mass m is released with negligible speed from a distance $\sqrt{3}R$ on the axis of a fixed uniformly charged ring of charge Q and radius R. Find out its velocity when it reaches at the centre of the ring.




**Solution :** As potential due to uniformly charged ring at its axis (at x distance) is :




$V = \frac{kQ}{\sqrt{R^2 + x^2}}$ ;

So, potential at point A due to ring

$V_1 = \frac{kQ}{\sqrt{R^2 + 3R^2}} = \frac{kQ}{2R}$

So potential energy of charge –q at point A

$P.E._1 = \frac{-kQq}{2R}$ and potential at point B, $V_2 = \frac{kQ}{R}$

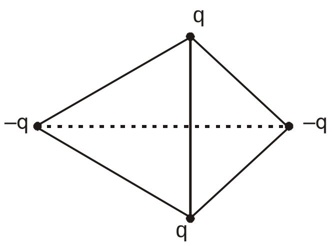
So, potential energy of charge –q at point B : $P.E._2 = \frac{-kQq}{R}$

Now by energy conservation : $P.E._1 + K.E._1 = P.E_2 + K.E_2$

$\frac{-kQq}{2R} + 0 = \frac{-kQq}{R} + \frac{1}{2}mv^2 \Rightarrow v^2 = \frac{kQq}{mR}$

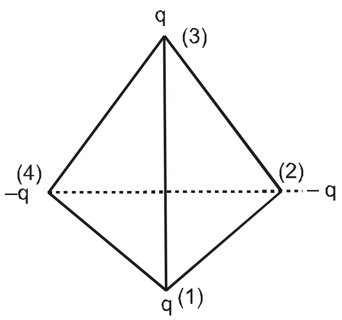
So velocity of charge – q at point B $v = \sqrt{\frac{kQq}{mR}}$ **Ans.**

**Problem 13.** Four small point charges (each of equal magnitude q) are placed at four corners of a regular tetrahedron of side a. Find out potential energy of charge system




**Solution :** Potential energy of system : $U = U_{12} + U_{13} + U_{14} + U_{23} + U_{24} + U_{34}$

$\therefore U = \frac{-kq^2}{a} + \frac{kq^2}{a} + \frac{-kq^2}{a} + \frac{-kq^2}{a} + \frac{kq^2}{a} + \frac{-kq^2}{a}$




Total potential energy of this charge system $U = \frac{-2kq^2}{a}$

**Problem 14.** If $V = x^2y + y^2z$ then find $\vec{E}$ at (x, y, z)

**Solution :** Given $V = x^2y + y^2z$ and $\vec{E} = -\frac{\partial v}{\partial r}$

$$\vec{E} = -\left[\frac{\partial V}{\partial x}\hat{i} + \frac{\partial V}{\partial y}\hat{j} + \frac{\partial V}{\partial z}\hat{k}\right] \Rightarrow \vec{E} = -[2xy\,\hat{i} + (x^2 + 2yz)\,\hat{j} + y^2\hat{k}]$$

**Problem 15.** Magnitude of electric field depends only on the x – coordinate as $\vec{E} = \frac{20}{x^2}\hat{i}$ V/m. Find

(i) The potential difference between two points A (5m, 0) and B(10m, 0).

(ii) Potential at x = 5 if V at $\infty$ is 10 volt.

(iii) In part (i), does the potential difference between A and B depend on whether the potential at $\infty$ is 10 volt or something else.

**Solution :** Given, $\vec{E} = \frac{20}{x^2}\hat{i}$ V/m

We know that : $\int dV = -\int \vec{E}.\vec{dr} \Rightarrow \int_{V_1}^{V_2} dV = -\int_{x_1}^{x_2} E_x dx = -\int_{x_1}^{x_2}\frac{20}{x^2}dx$

$\therefore$ Potential difference, $\Delta V = \frac{20}{x}\Big|_{x_1}^{x_2} \Rightarrow V_2 - V_1 = \frac{20}{x_2} - \frac{20}{x_1}$

(i) Potential difference between point A and B ($\Delta V$ for A to B)

$V_B - V_A = \frac{20}{10} - \frac{20}{5} = -2$ volt

(ii) $\Delta V$ for $x = \infty$ to $x = 5$

$V_5 - V_\infty = \frac{20}{5} - \frac{20}{\infty}$ $\therefore$ $V_5 = 10 + 4 = 14$ volt

(iii) Potential difference between two points does not depend on reference value of potential. So, the potential difference between A and B does not depend on whether the potential at $\infty$ is 10 volt or something else.

**Problem 16.** If $E = 2r^2$ then find V(r)

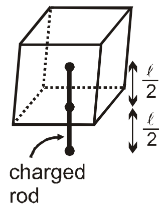
**Solution :** Given : $E = 2r^2$

we know that : $\int dv = -\int \vec{E}.\vec{dr} = -\int 2r^2\,dr \Rightarrow V(r) = \frac{-2r^3}{3} + c$ **Ans.**

**Problem 17.** A charge Q is uniformly distributed over a rod of length $\ell$. Consider a hypothetical cube of edge $\ell$ with the centre of the cube at one end of the rod. Find the minimum possible flux of the electric field through the entire surface of the cube.

**Solution :** According to Gauss law : Flux depends upon charge inside the closed hypothetical surface. So, for minimum possible flux through the entire surface of the cube, charge inside it should be minimum.

Linear charge density of rod = $\frac{Q}{\ell}$



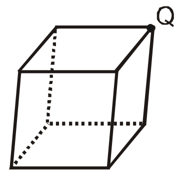

and minimum length of rod inside the cube = $\frac{\ell}{2}$

So, charge inside the cube = $\frac{\ell}{2}.\frac{Q}{\ell} = \frac{Q}{2}$

So, flux through the entire surface of the cube = $\frac{\Sigma q}{\epsilon_0} = \frac{Q}{2\epsilon_0}$

**Problem 18.** A charge Q is placed at a corner of a cube. Find the flux of the electric field through the six surfaces of the cube.

**Solution :** By Gauss's law, $\phi = \frac{q_{in}}{\varepsilon_0}$. Here, since Q is kept at the corner, so only $\frac{q}{8}$ charge is inside the cube. (Since, complete charge can be enclosed by 8 such cubes)

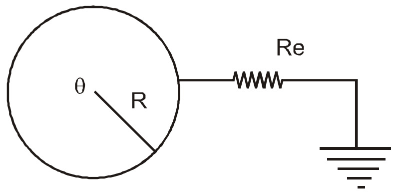



$\therefore \quad q_{in} = \frac{Q}{8}$ So, $\phi = \frac{q_{in}}{\varepsilon_0} = \frac{Q}{8\varepsilon_0}$ **Ans.**

**Problem 19.** An isolated conducting sphere of charge Q and radius R is grounded by using a high resistance wire. What is the amount of heat loss ?

**Solution :** When sphere is grounded, it's potential become zero which means all charge goes to earth (since sphere is conducting and isolated)
So, all energy in sphere is converted into heat

So, total heat loss = $\frac{kQ^2}{2R}$




**Problem 20.** An isolated conducting sheet of area A and carrying a charge Q is placed in a uniform electric field E, such that electric field is perpendicular to sheet and covers all the sheet. Find charges appearing on left and right surfaces of the conducting sheet. Also find the resultant electric field on the left and right side of the plate.

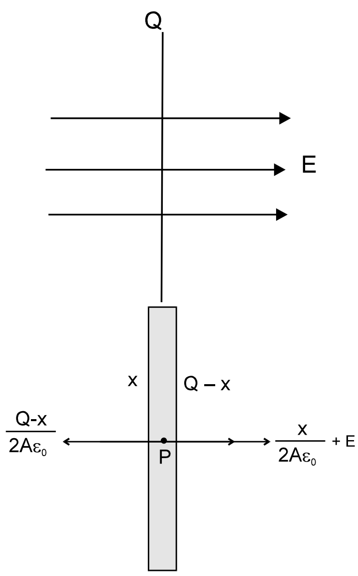
**Solution :** Let there is x charge on left side of plate and Q – x charge on right side of plate

$\because \quad E_P = 0$

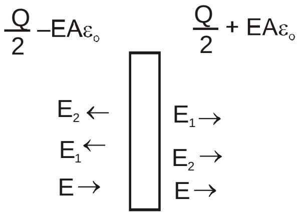
$\therefore \quad \frac{x}{2A\varepsilon_0} + E = \frac{Q-x}{2A\varepsilon_0}$ or $\frac{x}{A\varepsilon_0} = \frac{Q}{2A\varepsilon_0} - E$

$\therefore \quad x = \frac{Q}{2} - EA\varepsilon_0$ and $Q - x = \frac{Q}{2} + EA\varepsilon_0$

So, charge on one side is $\frac{Q}{2} - EA\varepsilon_0$ and other side $\frac{Q}{2} + EA\varepsilon_0$



Electric field due to conducting sheet at a point outside the sheet $= \dfrac{Q_{net}}{2A\varepsilon_o} = \dfrac{Q}{2A\varepsilon_o}$ and having direction away from plate $E_1$ and $E_2$ are electric field due to charges on left and right side of the plate respectively.

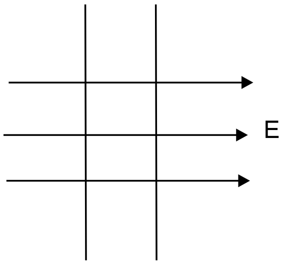



Now,

$$E_{left} = E - E_1 - E_2 = E - \frac{\left(\frac{Q}{2} - EA\varepsilon_o\right) + \left(\frac{Q}{2} + EA\varepsilon_o\right)}{2A\varepsilon_o} = E - \frac{Q}{2A\varepsilon_o} \text{ (towards right)}$$

$$E_{right} = E + E_1 + E_2 = E + \frac{Q}{2A\varepsilon_o} \text{ (towards right)}$$

**Problem 21.** Two uncharged and parallel conducting sheets, each of area A are placed in a uniform electric field E at a finite distance from each other, such that electric field is perpendicular to sheets and covers all the sheets. Find out charges appearing on its two surfaces.




**Solution :** Plates are conducting so net electric field inside these plates should be zero. So, electric field due to induced charges (on the surface of the plate) balance the outside electric field.

Here $\vec{E}_i$ = induced electric field

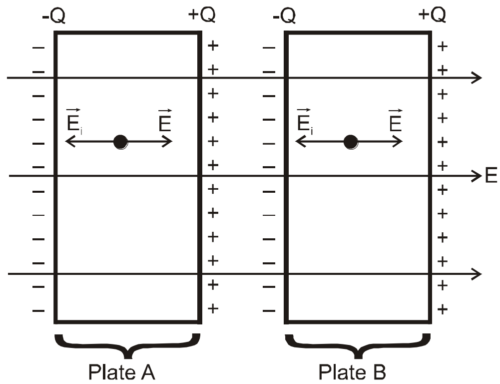
For both plates, $\vec{E}_i + \vec{E} = 0$

$\Rightarrow \quad \vec{E}_i = -\vec{E}$ .....(1)

Let charge induced on surfaces are +Q and – Q, then $|\vec{E}_i| = \dfrac{Q}{A\varepsilon_0}$

By equation (1)

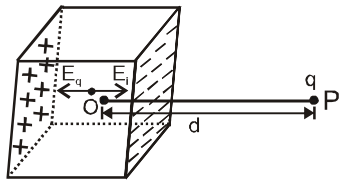
$\dfrac{Q}{A\varepsilon_0} = E \quad \Rightarrow \quad Q = AE\varepsilon_0$ **Ans.**




**Problem 22.** A positive charge q is placed in front of a conducting solid cube at a distance d from its centre. Find the electric field at the centre of the cube due to the charges appearing on its surface.

**Solution :** Here $\vec{E}_i = \vec{E}$ electric field due to induced charges and $E_q$ = electric field due to charge q

We know that net electric field in a conducting cavity is equal to zero.




i.e. $\vec{E} = \vec{0}$ at the centre of the cube.

$\Rightarrow \quad \vec{E}_i + \vec{E}_q = \vec{0}$

$\Rightarrow \quad \vec{E}_i = -\vec{E}_q \quad \Rightarrow \quad \vec{E}_i = -\dfrac{kq}{d^2}\overrightarrow{PO}$ **Ans.**